

高強度テラヘルツパルスで誘起する 非線形光学現象

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物質—細胞統合システム拠点

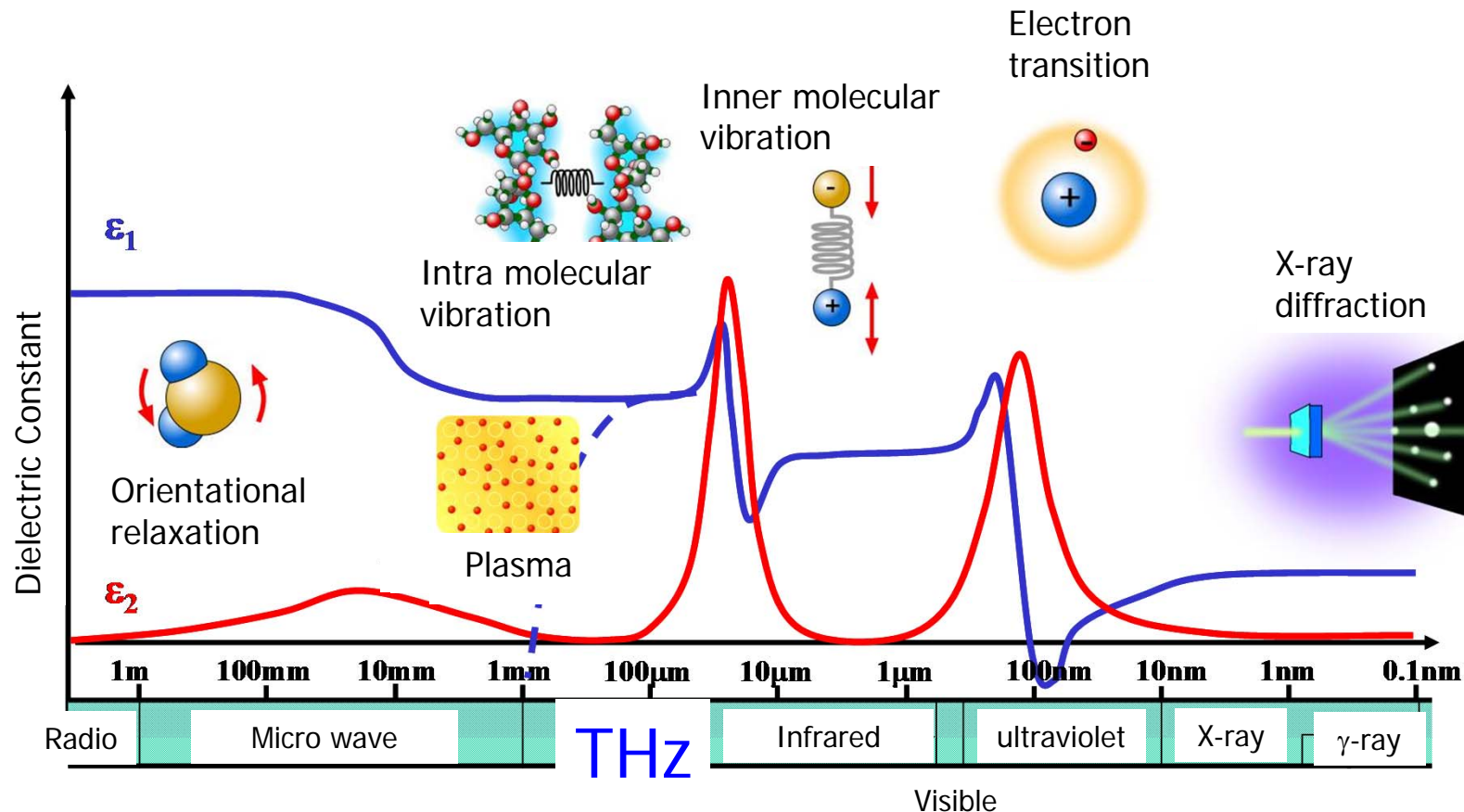
Contents



- THz time-domain spectroscopy
- Intense THz pulse generation
 - Tilted-pulse front scheme with a LiNbO_3 crystal
H. Hirori, et al., Appl. Phys. Lett. 98, 091106 (2011)
- Nonlinear THz phenomena
 - Carrier multiplication
(GaAs/AlGaAs multiple quantum wells)
H. Hirori, et al., Nature Comms. 2, 594(2011)

THz region (0.1-10 THz)

(1THz \triangleq 300 μ m \triangleq 33cm⁻¹ \triangleq 4.1meV)

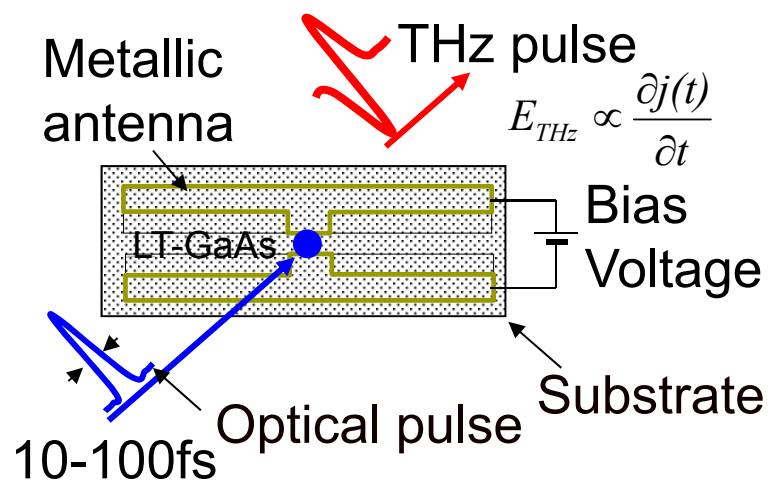


Superconducting gap, Soft phonon modes in ferroelectrics, Excitonic resonance, subband transition, Intra molecular vibration of bio-molecules, Rotational mode of gases, etc.

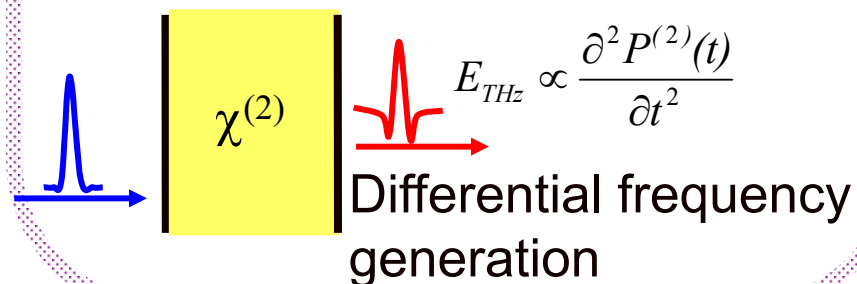
THz generation and detection

Generation

Photoconductive antenna



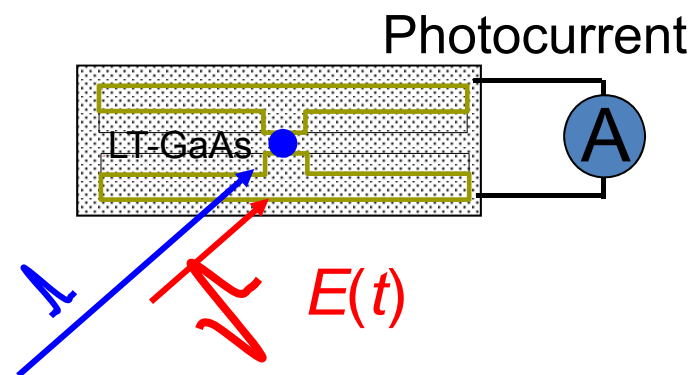
Nonlinear optical process



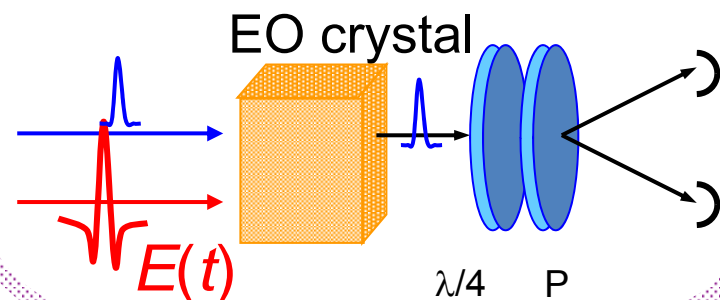
✓ 0.1-100 THz (3 mm-3 μ m, 0.4-400 meV) is available.

Detection

Photoconductive antenna



Electro optic sampling



THz Time-domain spectroscopy

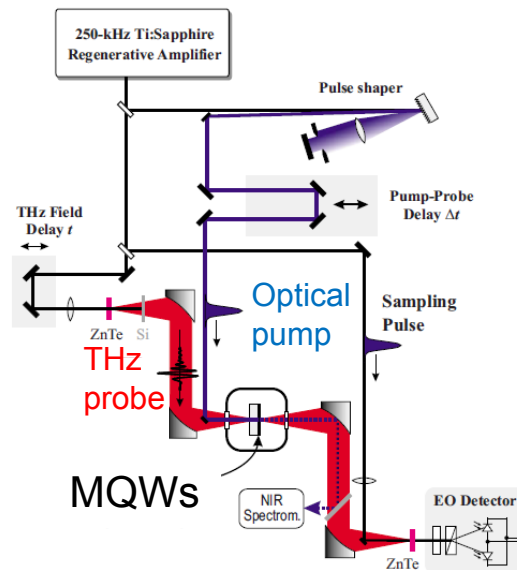
✓ Electric field measurement \Rightarrow Complex dielectric constants.

$$E(\omega), \phi(\omega)$$

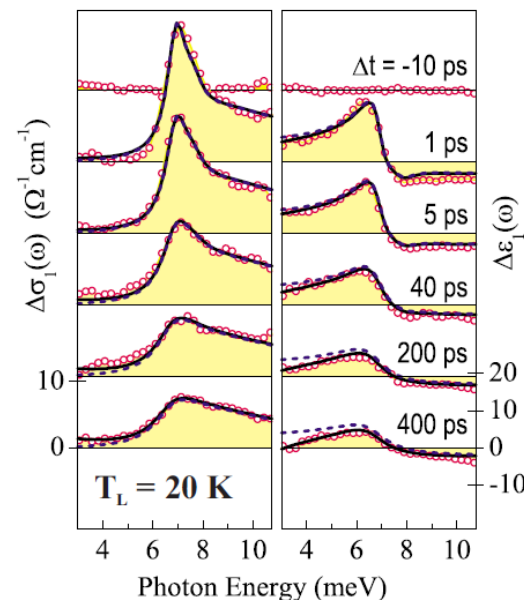
✓ Allowing for ultrafast time-resolved measurement.
 Ultrafast phenomena in semiconductors, phase-transition and superconducting materials.

R. Huber et al., Nature (2001). C. Kaindl et al., Nature (2003).

T. Suzuki and R. Shimano, PRL(2009).



C. Kaindl et al., PRB(2009).



Contents

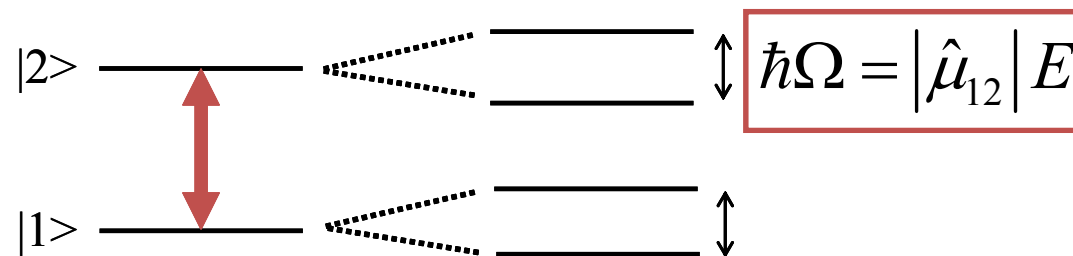


- THz time-domain spectroscopy
- **Intense THz pulse generation**
 - **Tilted-pulse front scheme with a LiNbO₃ crystal**
H. Hirori, et al., Appl. Phys. Lett. 98, 091106 (2011)
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Nonlinear phenomena in solids with intense THz pulse

- Resonantly THz-Driven Systems**

Rabi oscillations, dressed states, the AC (or optical) Stark effect, the Autler-Townes effect, electromagnetically-induced transparency, gain without inversion, ..., etc.



- Nonlinear transport phenomena**

Bloch oscillation, Inter-valley scattering (Gunn's effect), Impact ionizations (carrier multiplication)

- Available for inducing phase transition phenomena**

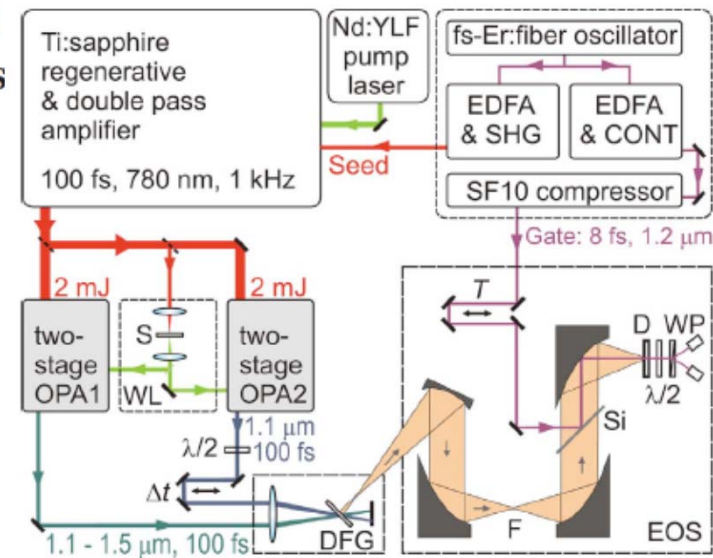
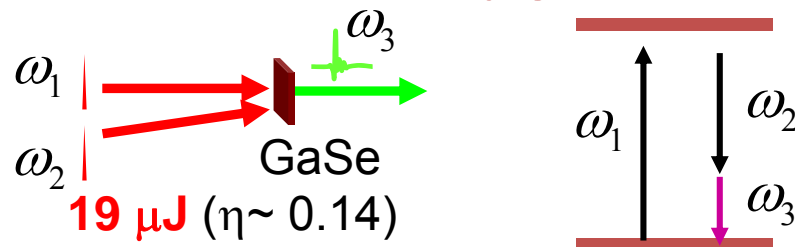
Intense THz pulse generation (10-80 THz)

December 1, 2008 / Vol. 33, No. 23 / OPTICS LETTERS 2767

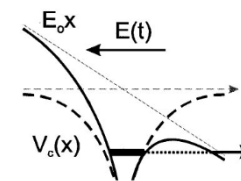
Phase-locked generation and field-resolved detection of widely tunable terahertz pulses with amplitudes exceeding 100 MV/cm

Alexander Sell, Alfred Leitenstorfer, and Rupert Huber*

Differential frequency generation



- Field strengths that outer-shell electrons in atoms experiences.
- Higher frequency region (10-80 THz) has a GOOD source! (19 μJ, $\eta \sim 0.14$).
- Lower frequency ???

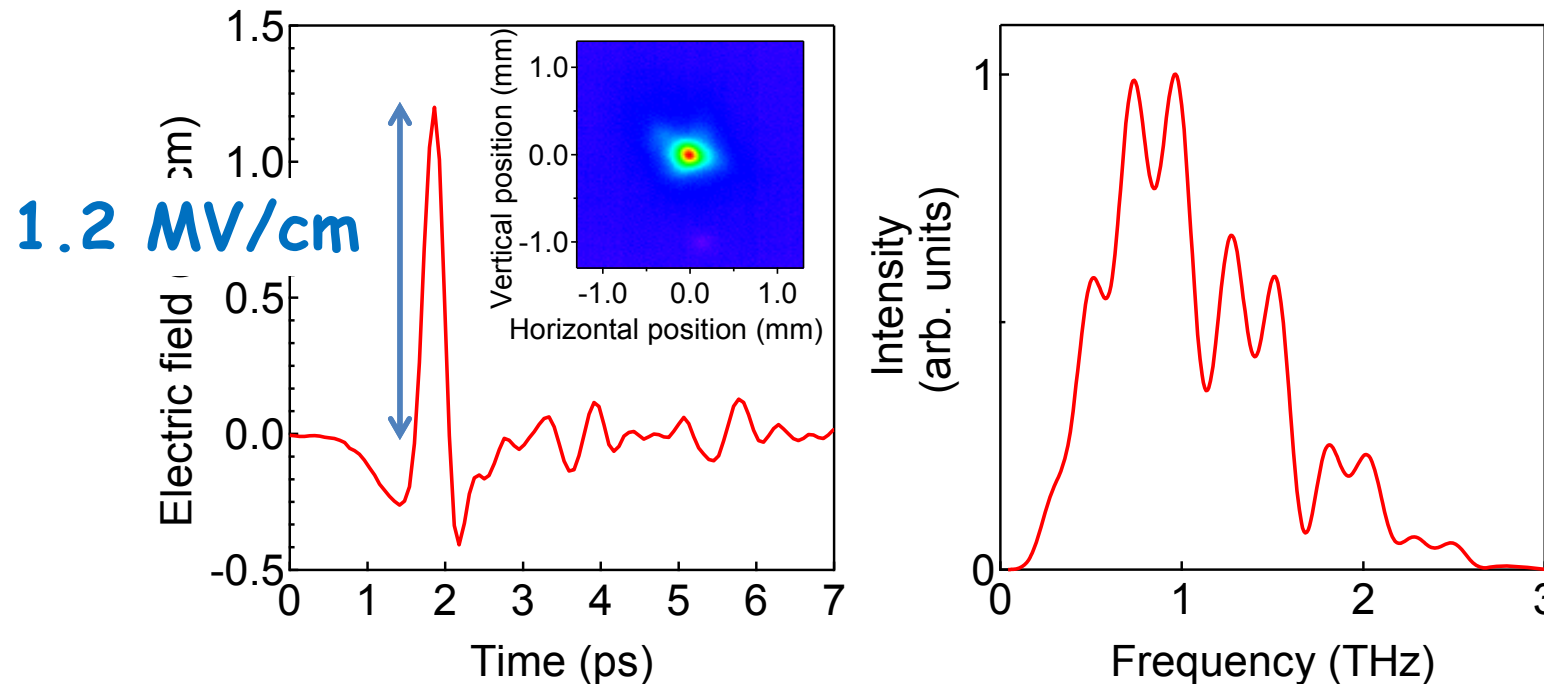


$$e/4\pi\epsilon_0 a_B^2 \sim 100 \text{ MV/cm} - 1 \text{ GV/cm}$$

Tilted-pulse front scheme with LiNbO_3 crystal (below 3 THz)



Temporal profile and spectra H. Hirori, et al., APL (2011)



- Nearly half-cycle sub-picosecond pulse with a maximum peak field of 1.2 MV/cm ($3 \mu\text{J}$, $\eta \sim 10^{-3}$, spot diameter $\sim 300 \mu\text{m}$).
- Around 1 THz of center frequency.
- Optical rectification of femtosecond lasers.

Characteristic of LiNbO_3 for THz generation

Good point

- High second order nonlinear susceptibility
- Large band-gap energy (3.7 eV) (less multi-photon absorption)

Appl. Phys. B 78, 593 (2004)

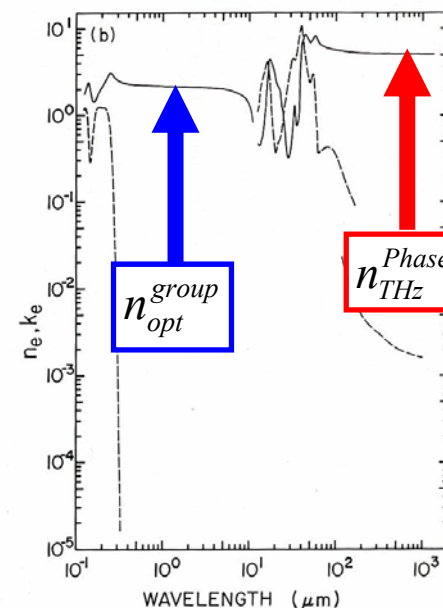
Material	d (pm/V)
CdTe	81.8
GaAs	65.6
GaP	24.8
ZnTe	68.5
GaSe	28.0
<u>LiTaO₃</u>	<u>161</u>
<u>LiNbO₃</u>	<u>168</u>

Bad point

- Large difference between $n_{\text{THz}}^{\text{phase}}$ (= 5) and $n_{\text{opt}}^{\text{group}}$ (= 2.2)

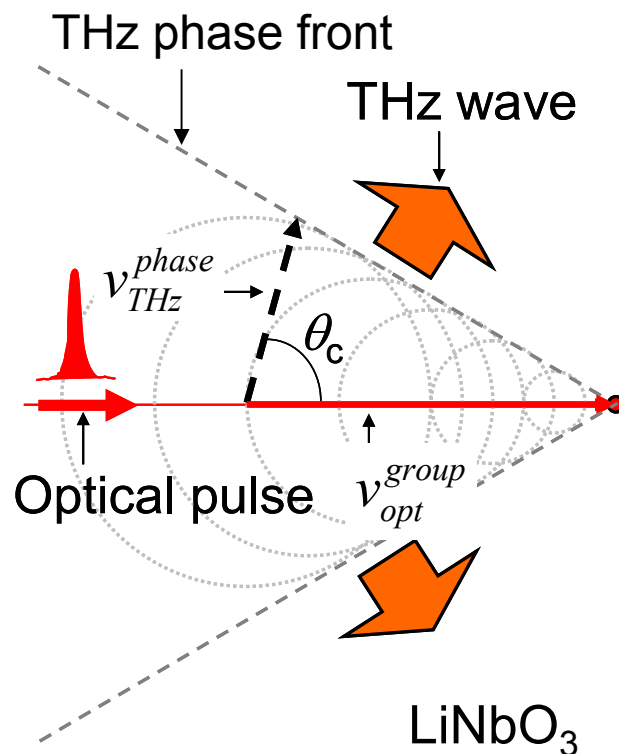
$$\Rightarrow v_{\text{THz}}^{\text{phase}} \ll v_{\text{opt}}^{\text{group}}$$

Velocity (phase) mismatching



D. Palik, Academic Press (1985).

THz Cherenkov wave



D. H. Auston et al.,
PRL **53**, 1555 (1984).

$$v_{THz}^{phase} \ll v_{opt}^{group}$$

➡ The condition necessary for the shock (or Cherenkov) wave radiation of a supersonic aircraft or bullet.

Direction of THz wave

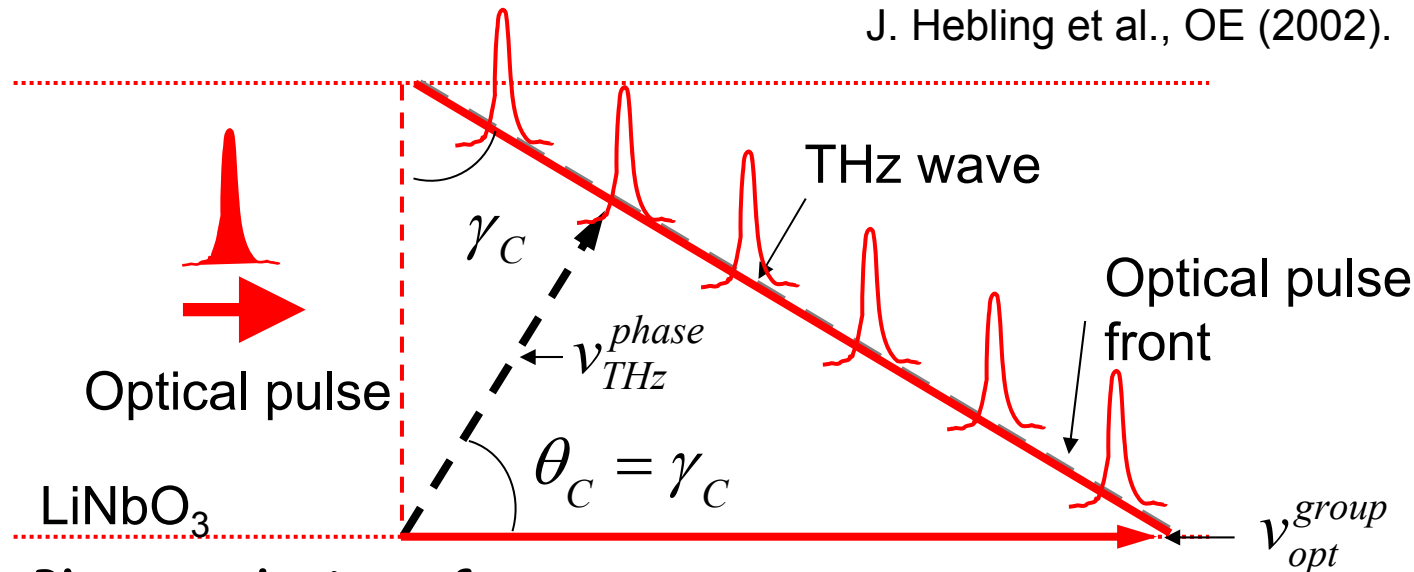
$$v_{THz}^{phase} = v_{opt}^{group} \cos \theta_C$$

$$\cos \theta_C = n_{opt}^{group} / n_{THz}^{phase}$$

➡ $\theta_C = 62 \text{ deg}$

Velocity matching by the tilted-pulse-front scheme

J. Hebling et al., OE (2002).



Phase velocity of THz wave

$$v_{THz}^{phase} = v_{opt}^{group} \cos \theta_C$$

Noncollinear velocity of tilted optical pulse

$$v_{opt}^{\gamma_C} = v_{opt}^{group} \cos \gamma_C$$

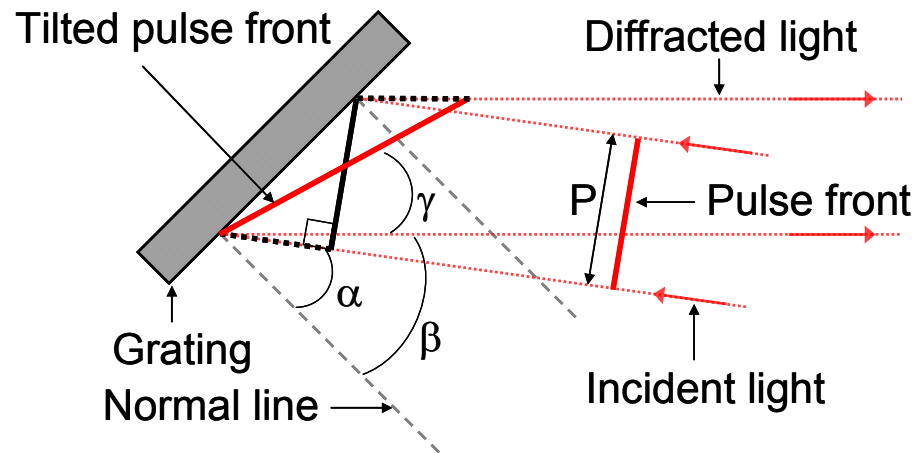
When $\gamma_C = \theta_C$,

$$\Rightarrow v_{THz}^{phase} = v_{opt}^{\gamma_C}$$

Phase matching condition can be satisfied.

How to tilt the pump pulse front?

Grating



$$\tan \gamma = \frac{\cos \beta}{\sin \alpha + \sin \beta}$$

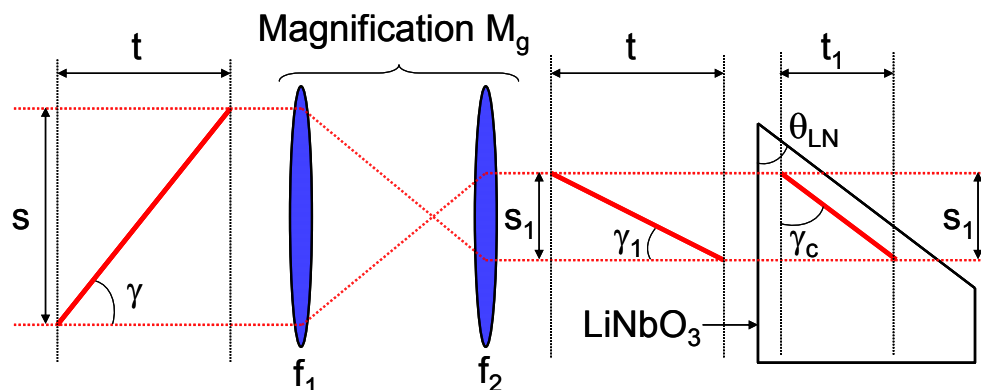
$$\sin \alpha + \sin \beta = mp\lambda_0,$$

m : diffraction order

p : groove density of a grating

λ_0 : the central wavelength

Lens pair

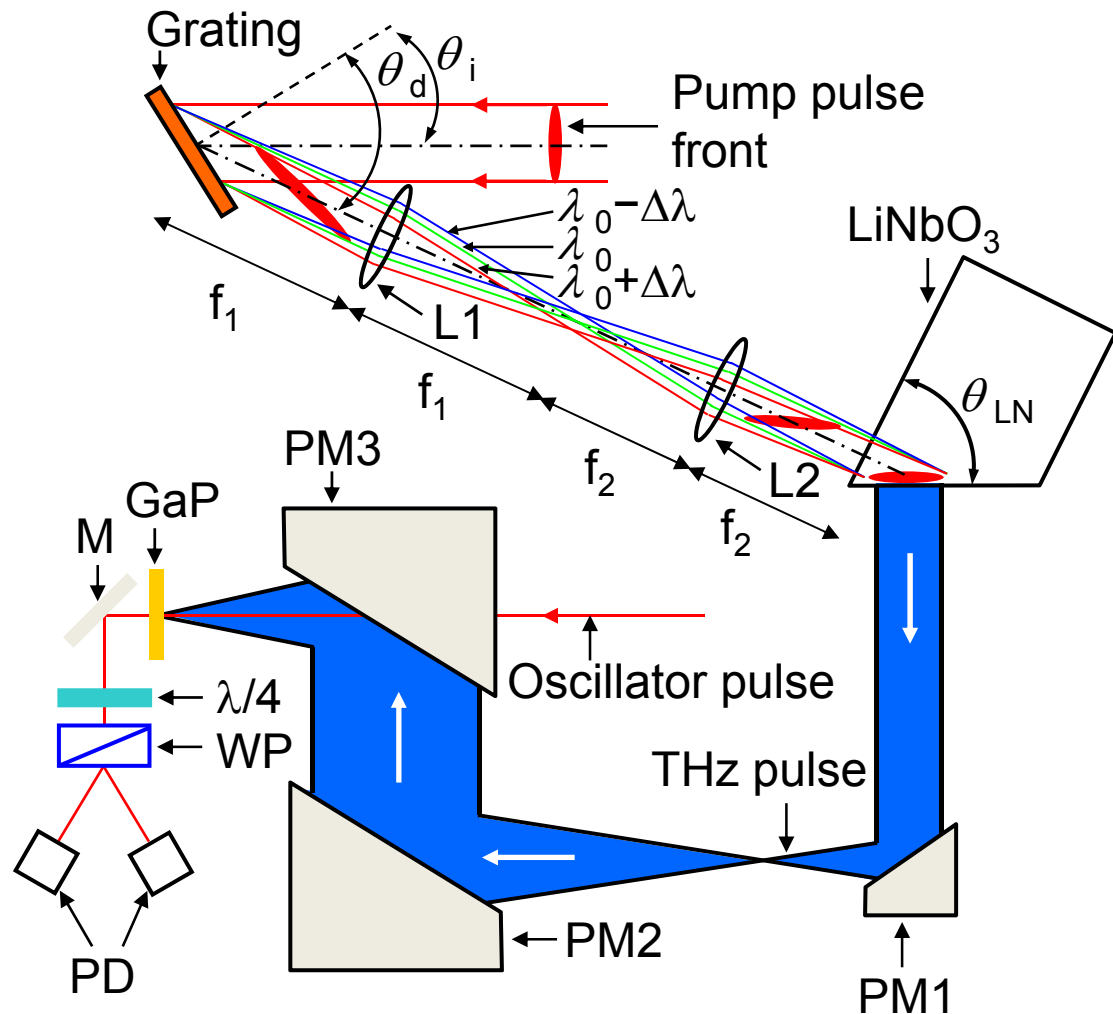


$$\tan \gamma_c = \frac{m\lambda_0 p}{n_g M_g \cos \beta}$$

n_g : group refractive index of LN crystal

$$\gamma_c = \theta_c = 62^\circ$$

THz pulse generation setup



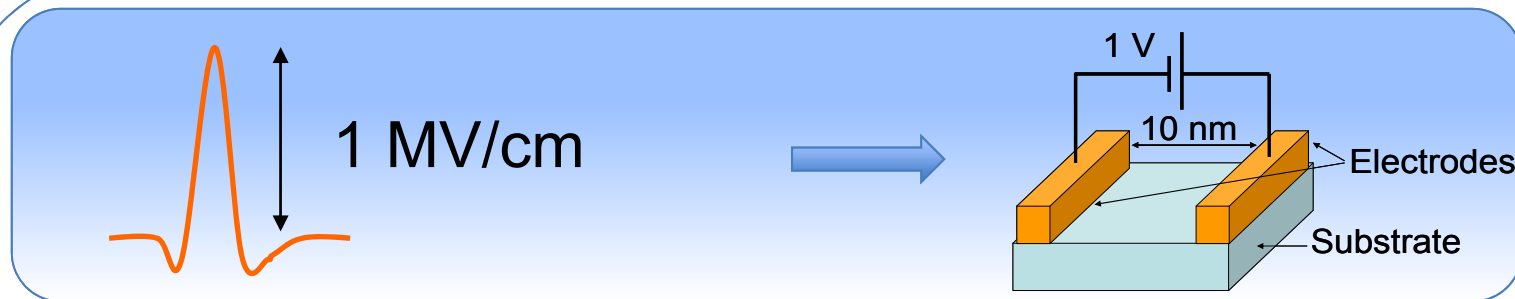
Laser source

- Ti: sapphire
- 4 mJ/pulse
- 1 kHz
- 780 nm
- 85 fs (FWHM)

Key factors

- Good phase matching condition
- High nonlinear susceptibility
- High pumping power
- Tight focusing the collimated THz beam

Strong picosecond DC electric field



- Achieved in actual electric devices
- Useful for characterization of electric devices

Nonlinear transport phenomena

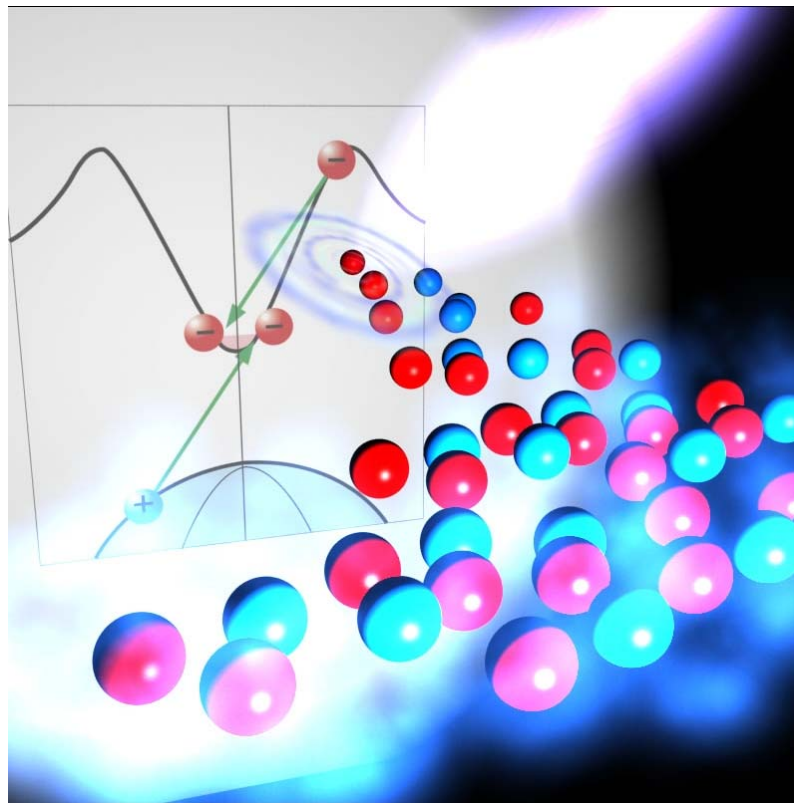
- Bloch oscillations
- Inter-valley scattering
- Impact ionization
- Zener tunneling

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H. Hirori, et al., Nature Comms. 2, 594(2011)

Extraordinary carrier multiplication in GaAs MQWs gated by intense terahertz pulse



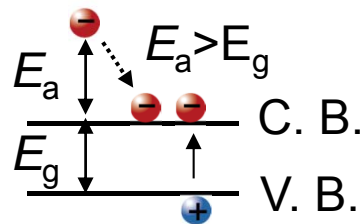
H. Hirori, et al., Nature Commun.
2, 594(2011)

Nonlinear Transport Phenomena in Semiconductors

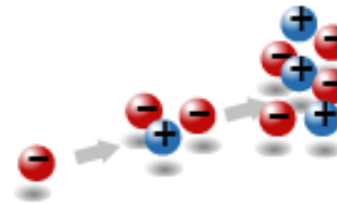
Carrier multiplication in high electric field

- Electron accelerated by electric field can gain a kinetic energy and excite other electrons.

Impact ionization

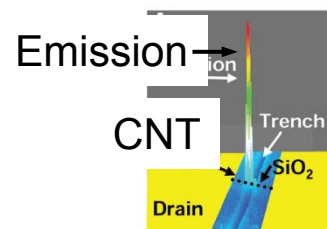


Carrier multiplication



Various applications

- Avalanche photodiodes
- Electroluminescent and photovoltaic nano-devices



J. Chen et al., Science (2005).

Fundamental

- Important for nonlinear transport phenomena.

➡ The elementary scattering process has been unclear.

Purpose



Clarifying carrier multiplication process of GaAs under high electric field

- THz pulse with an amplitude exceeding 1 MV/cm

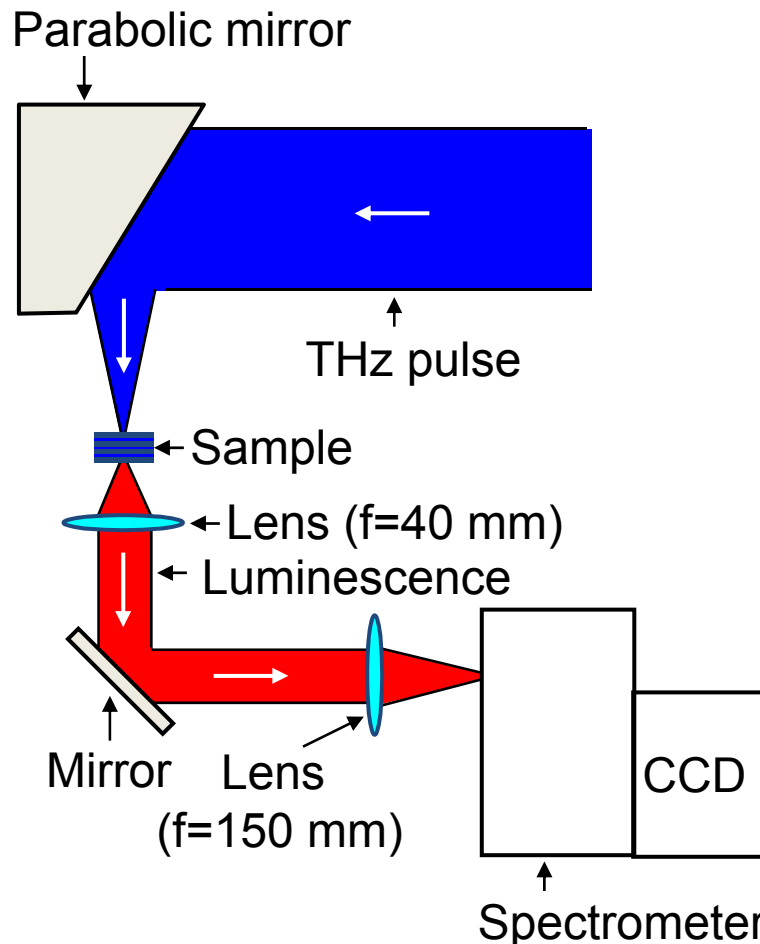
⇒ Sufficient carrier multiplication.

- Luminescence measurement

⇒ Direct evidence of carrier generation.

Experimental setup for THz induced luminescence

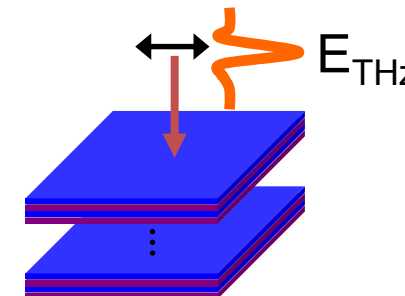
Setup



GaAs quantum wells

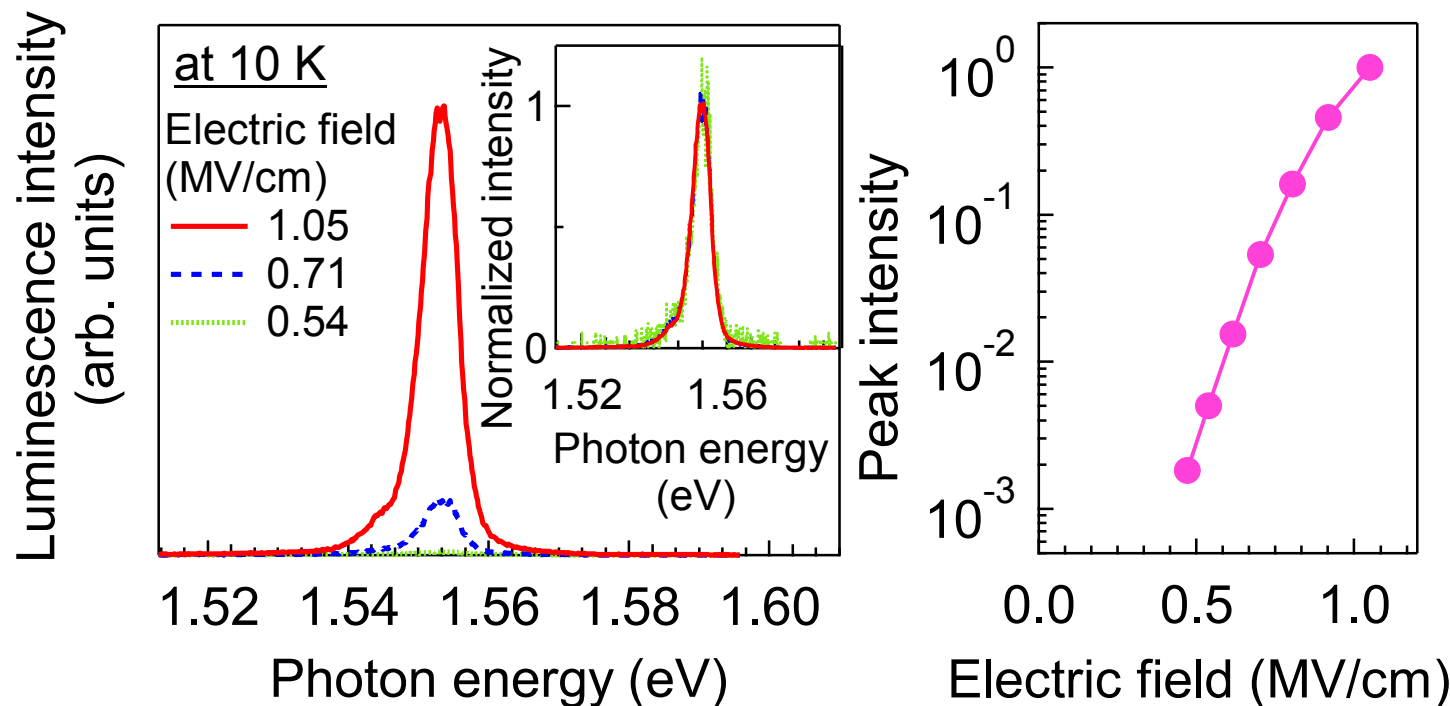


- Nominally non-doped (Residual dominant impurity donors are Sulfur.)
- Low temperature measurement (10 -150 K)
- The polarization of electric field is in plane and along the (100) direction of the sample



THz induced luminescence without photoexcitation

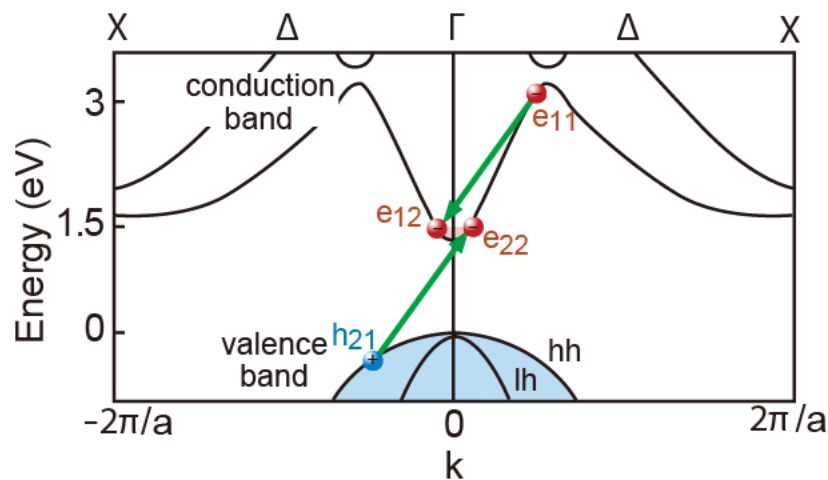
Electric field dependence



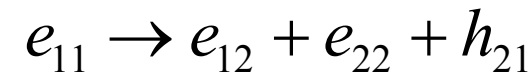
- Luminescence centered around 1.55 eV.
- Electric field dependence shows extremely nonlinear ($\propto \mathcal{E}^8$).
- The number of carriers increases by about three orders of magnitude.

Carrier multiplication process in GaAs

Single impact ionization



- Electrons seeded from residual impurity donors by the field ionization
- Doubling the number of electrons.



e_{ij}, h_{ij} : electron and hole

Carrier multiplication

Carrier density $N(\varepsilon)$ after $\langle n_I \rangle$ times impact ionization events:

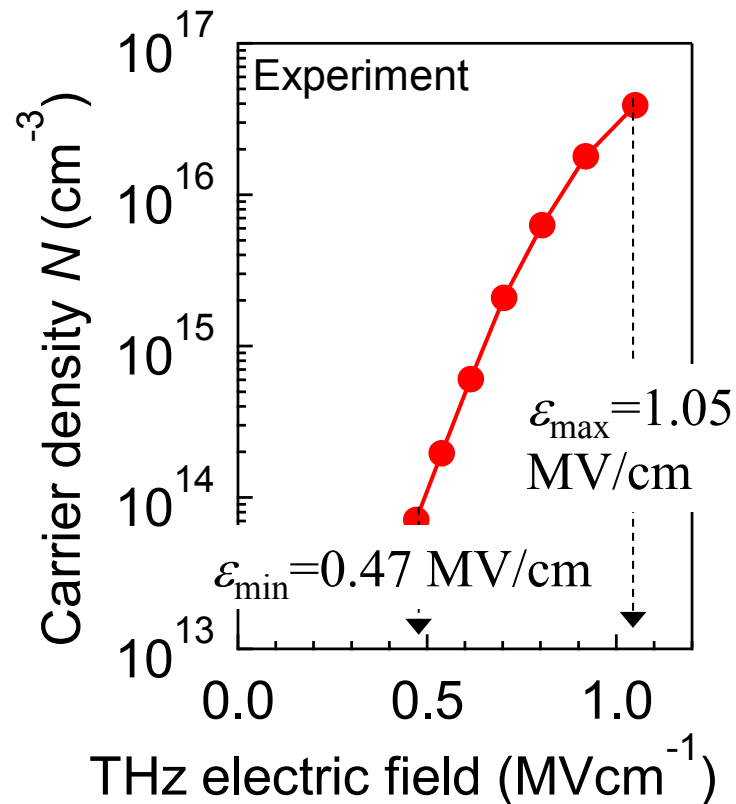
$$N(\varepsilon) = N_0 \times 2^{\langle n_I \rangle}$$

N_0 : initial electron density

$\langle n_I \rangle$: number of impact ionization events

Impact ionization number derived from experimental result

Carrier density

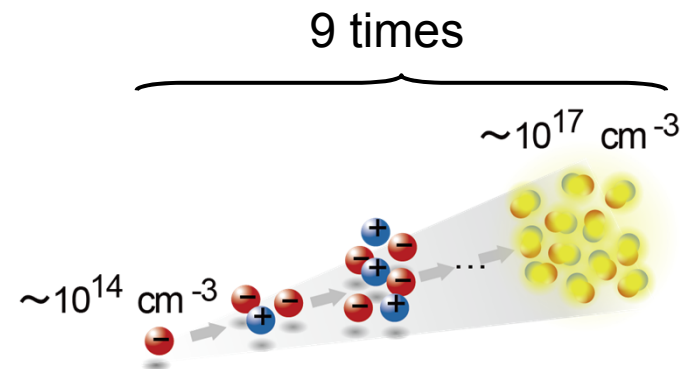


Increment of impact ionization number $\langle \Delta n_I \rangle$

$$N(\varepsilon) = N_0 \times 2^{\langle n_I \rangle}$$

$$\langle \Delta n_I \rangle = \log_2(N(\varepsilon_{\max})/N(\varepsilon_{\min})) \sim 9$$

- 9 impact ionizations induce numerous $e-h$ pairs.



How many times do impact ionization occur?

Changing k with electric field $\varepsilon(t)$

$$\hbar \frac{dk(t)}{dt} = -e\varepsilon(t)$$

$k(t)$: average electron wavenumber

$\varepsilon(t)$: electric field

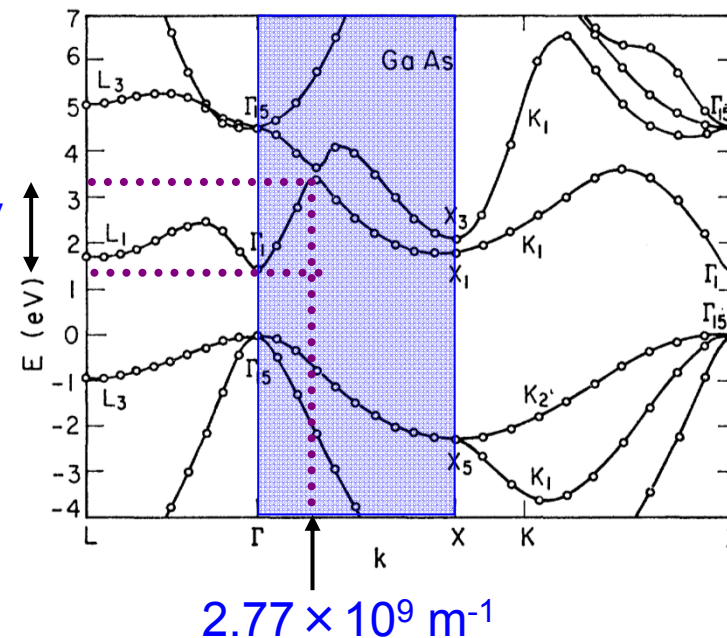
e : elementary charge

\hbar : Planck's constant

$E_{th} = 1.7 \text{ eV}$

The GaAs dispersion relation yields the E_{th} of 1.7 eV at a wavenumber k of $2.77 \times 10^9 \text{ m}^{-1}$.

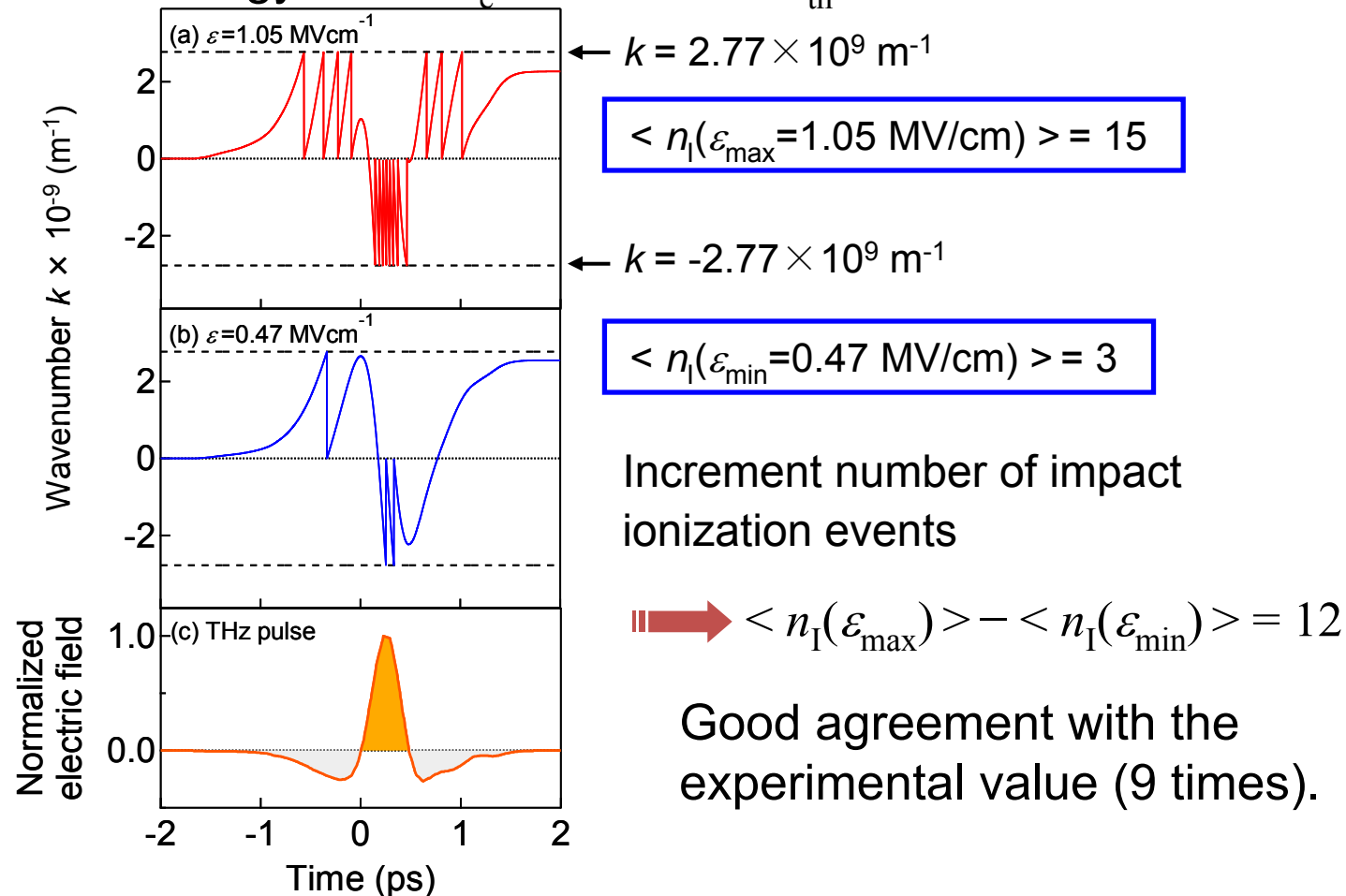
M. L. Cohen et al. *Phys. Rev.* **141**, 789 (1966).



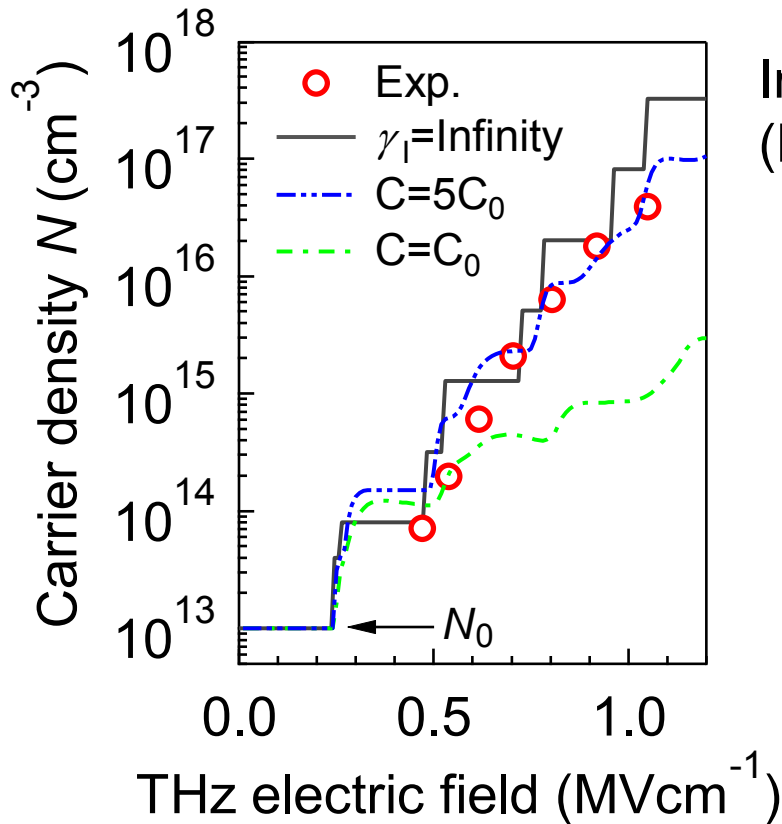
➡ The number of times the k exceeds $2.77 \times 10^9 \text{ m}^{-1}$ in electric field of THz pulse should be calculated.

Calculated number of impact ionization events

- Assumption that the accelerated electrons lose all kinetic energy when K_e achieve the E_{th} .



Experiment v.s. calculation



Impact ionization rate
(Keldysh formula):

$$\gamma_I = C(E - E_{th})^2, \quad E > E_{th}.$$

$$C = C_0 = 870 \text{ ps}^{-1}\text{eV}^{-2} \text{ and } E_{th} = 1.7 \text{ eV}.$$

M. V. Fischetti and S. E. Laux,
Phys. Rev. B **38**, 9721 (1988).

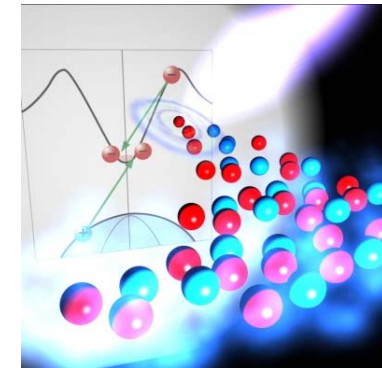
Good agreement of impact ionization numbers and carrier densities between the experimental and calculated results.

Summary



Result

- ✓ Generating the world's strongest THz pulse.
(Achieving at electric field of 1 MV/cm.)
- ✓ Observation of bright luminescence by irradiating 1-MV/cm THz pulse to GaAs MQWs ($E_g \sim 400 h\nu_{\text{THz}}$).



Outlook

- ✓ Further increasing up the THz generation efficiency.
- ✓ Studying Zener tunneling and Bloch oscillations with higher electric field.
- ✓ Applying new materials showing phase transitions.