

## Two-Dimensional Electronic Structure of La and P Co-Doped $\text{CaFe}_2\text{As}_2$ Studied by ARPES

The relation between the occurrence of high-temperature superconductivity and the underlying electronic structure in iron-based superconductors is a major unresolved problem. We revealed the electronic structure of La and P co-doped  $\text{CaFe}_2\text{As}_2$  superconductor with the critical temperature ( $T_c$ ) = 45 K by angle-resolved photoemission spectroscopy. Its Fermi surface topology consists of nearly two-dimensional hole- and electron-like Fermi surfaces. Compared to other high- $T_c$  iron-based superconductors, we discuss the relationship between the dimensionality of Fermi surface topology and the value of  $T_c$ .

The discovery of high critical temperature (high- $T_c$ ) superconductivity in  $\text{LaFeAsO}_{1-x}\text{F}_x$  [1] has triggered intensive research on iron-based superconductors [2-4]. Density functional theory calculations for iron-based superconductors predict that Fermi Surfaces (FSs) in these materials are composed of nearly cylindrical (two-dimensional) hole- and electron-like FSs at the Brillouin zone center and corner, respectively, as shown in Fig. 1(a) [5, 6]. For the emergence of iron-based superconductivity, the FSs nesting that can enhance spin and/or orbital fluctuations has been believed to be important for realizing exotic Cooper pairing [5-8].

Experimentally, FS topology and its nesting condition for several iron-based superconductors have been directly studied by angle-resolved photoemission spectroscopy (ARPES) [9-22]. The FS topology for  $\text{AFe}_2\text{As}_2$  (so-called 122-type bulk superconductors, where A represents alkali or alkali-earth metals) has revealed that high- $T_c$  122-type bulk superconductors have both hole- and electron-like FSs with the electron-like FSs in common [19-22]. However, the dimensionality of the FSs, especially hole-like FSs, shows marked material dependence: while in hole-doped  $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$  ( $T_c = 38$  K), all hole-like FSs have a cylindrical shape [Fig. 1(b)] [19, 20], in electron-doped  $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$  ( $T_c = 25$  K), two- and three-dimensional hole-like FSs are present [Fig. 1(c)] [21, 22]. Since the presence of such a three-dimensional FS is not favorable for superconductivity

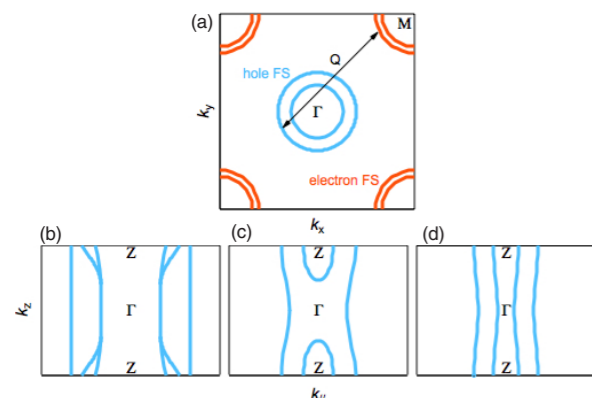


Figure 1: (a) Fermi surface (FS) and the nesting with nesting vector  $Q$  in the  $k_x$ - $k_y$  plane for iron-based superconductors. The sky-blue and orange curved lines represent hole- and electron-like FSs, respectively. (b)-(d) Hole-like FSs in the  $k_x$ - $k_y$  plane for  $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$ ,  $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$ , and  $\text{Ca}_{0.82}\text{La}_{0.18}\text{Fe}_2(\text{As}_{0.94}\text{P}_{0.06})_2$  obtained from refs. 19, 21 and 24, respectively.

resulting from the FS nesting, the relationship between the FS nesting and superconductivity is still unclear in iron-based superconductors. In this regard, the recent discovery of electron-doped  $\text{CaFe}_2\text{As}_2$  with  $T_c = 45$  K, namely La and P co-doped  $\text{Ca}_{1-x}\text{La}_x\text{Fe}_2(\text{As}_{1-x}\text{P}_x)_2$  [23], may be of importance. This value of  $T_c$  is the highest yet reported among iron-based bulk superconductors that have been studied by ARPES. In order to check the characteristic FS topology for the emergence of high- $T_c$  in iron-based superconductors, we performed ARPES measurements for  $\text{Ca}_{0.82}\text{La}_{0.18}\text{Fe}_2(\text{As}_{0.94}\text{P}_{0.06})_2$  superconductor [24].

High-quality single crystals of  $\text{Ca}_{0.82}\text{La}_{0.18}\text{Fe}_2(\text{As}_{0.94}\text{P}_{0.06})_2$  were grown as described elsewhere [23]. ARPES measurements were carried out at BL-28A of the Photon Factory using circularly polarized light ( $h\nu = 40$ – $86$  eV) and at BL-9A of the Hiroshima Synchrotron Radiation Center using circularly and linear polarized light ( $h\nu = 19$ – $31$  eV). The total energy resolution was set to 10–30 meV. Clean surfaces were obtained by *in situ* cleaving of crystal in a working vacuum better than  $3 \times 10^{-9}$  Pa and measured at 60 K (above  $T_c$ ).

In Figs. 2(a) and (b), we show ARPES intensity plots taken along the  $\Gamma$ -X direction with *s*- and *p*-polarized light, respectively ( $h\nu = 31$  eV). We observed one hole-like band whose top is located around 30 meV below  $E_F$  ( $\alpha_1$ ) and two hole-like bands with crossing  $E_F$  ( $\alpha_2$  and  $\beta$ ). In the *s*-pol data along the Z-X [Fig. 2(c)], we found that the top of the  $\alpha_1$  band is located around 10 meV below  $E_F$ , indicating that the  $\alpha_1$  band cannot form FS between the  $\Gamma$  and Z points. On the other hand, the  $\alpha_2$  and  $\beta$  bands cross  $E_F$  at Z, as seen in the *p*-pol data [Fig. 2(d)]. Since the position of the Fermi wave vector ( $k_F$ ) of the  $\alpha_2$  and  $\beta$  bands has almost no  $k_z$  dependence [Fig. 2(h)], we find these bands form nearly two-dimensional hole-like FSs. Around the M and A points, we observed two electron-like bands ( $\epsilon$  and  $\delta$ ) as seen in Figs. 2(e) and (f). These form elliptical electron-like FSs [Fig. 2(g)], and their  $k_F$  position shows sizeable undulation along the  $k_z$  direction [Fig. 2(i)] reflecting the elliptical shape of these FSs and the shape of the boundary in the body-centered tetragonal Brillouin zone. Figs. 2(g-i) show the shape of the four FSs ( $\alpha_2$ ,  $\beta$ ,  $\epsilon$ ,  $\delta$ ) observed by our ARPES measurements [24]. The total carrier number deduced from the volume of FSs of

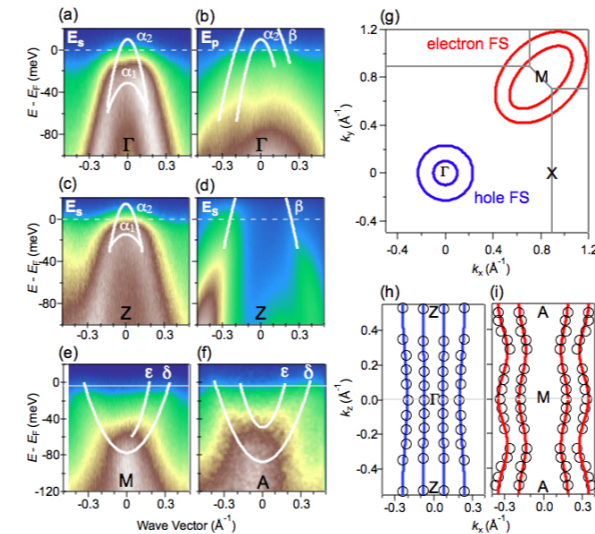


Figure 2: (a), (b) ARPES intensity plots taken along the  $\Gamma$ -X direction with *s*- and *p*-polarized light, respectively ( $h\nu = 31$  eV). (c), (d) are the same as (a), (b) but measured along the Z-X direction ( $h\nu = 19$  eV). (e), (f) ARPES intensity plots taken at M and A points, respectively ( $h\nu = 69$  eV and  $h\nu = 86$  eV). In (a)-(f), the white curves represent the band dispersion. (g) The FS topology in the  $k_x$ - $k_y$  plane. The blue and red curved lines represent hole- and electron-like FSs, respectively. (h) and (i) The FS topology in the  $k_x$ - $k_z$  plane. In (h) and (i), open circles show the positions of  $k_F$  deduced from the momentum distribution curve analysis of photon-energy dependent ARPES data [24].

$0.12 \pm 0.07$  e/Fe agrees with the value expected from the chemical composition of 0.09 e/Fe.

Next we discuss the implications of the present ARPES results for iron-based superconductivity. As shown in Figs. 1(b) and (d), all hole-like FSs in hole-doped  $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$  and electron-doped  $\text{Ca}_{0.82}\text{La}_{0.18}\text{Fe}_2(\text{As}_{0.94}\text{P}_{0.06})_2$  have nearly a two-dimensional shape in common [19, 20]. From these results, we suggest that the two-dimensional FS topology, leading to the good FS nesting condition, is universal for high- $T_c$  superconductivity regardless of the type of doped carrier. This supports the unconventional mechanisms for superconductivity in iron-based superconductors.

In summary, we investigated the three-dimensional electronic structure near  $E_F$  in electron-doped  $\text{Ca}_{0.82}\text{La}_{0.18}\text{Fe}_2(\text{As}_{0.94}\text{P}_{0.06})_2$  ( $T_c = 45$  K). We observed a nearly two-dimensional FS topology similar to that of  $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$ , demonstrating the common identity of the electronic structure to induce high- $T_c$  in 122-type iron-based superconductors.

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## BEAMLINE

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