

Observation of spin structure in CrNb<sub>3</sub>S<sub>6</sub> by means of resonant soft x-ray scatteringTakashi Honda<sup>1,\*</sup>, Yuichi Yamasaki<sup>2,3</sup>, Hironori Nakao<sup>1</sup>, Yuichi Murakami<sup>1</sup>, Takahiro Ogura<sup>4</sup>, Yusuke Kousaka<sup>5,6</sup>, and Jun Akimitsu<sup>7</sup><sup>1</sup>CMRC and PF, Institute of Materials Structure Science, KEK, Tsukuba 305-0801, Japan<sup>2</sup>Department of Applied Physics and QPEC, University of Tokyo, Tokyo 113-8656, Japan<sup>3</sup>RIKEN CEMS, Wako 351-0198, Japan<sup>4</sup>Department of Physics and Mathematics, Aoyama-Gakuin University, Kanagawa 252-5258, Japan.<sup>5</sup>Graduate School of Science, Hiroshima University, Hiroshima, 739-8526, Japan<sup>6</sup>Center for Chiral Science, Hiroshima, 739-8526, Japan<sup>7</sup>Faculty of Science, Okayama University, Okayama 700-8530, Japan

We clarified a formation of chiral magnetic soliton lattice in mono-axial chiral magnet CrNb<sub>3</sub>S<sub>6</sub> as seen via small-angle resonant soft x-ray scattering (RSXS) near the Cr *L*-absorption edge. In our results, the magnetic-field dependence of the higher harmonic magnetic diffraction and that of chiral magnetic soliton lattice constant are found to agree well with a theoretical magnetic structure predicted based on the chiral XY-spin model. We also observe deformations of the spin structure from the predicted chiral magnetic soliton lattice near the critical temperature.

## 1 Introduction

CrNb<sub>3</sub>S<sub>6</sub> is one of the chiral magnets, and forms chiral soliton lattice (CSL) by applying small magnetic fields perpendicular magnetic modulation vector along *c*-axis. It has been investigated by theory and by magnetic, electric transport properties measurements on single crystals.[1,2] In order to study the magnetic-field dependence of the magnetic structure for the CSL in more detail, we performed a small-angle RSXS.

## 2 Experiment

A single crystal of CrNb<sub>3</sub>S<sub>6</sub> with a volume of ~0.01 mm<sup>3</sup> was grown by the chemical vapor transport method.[3] A thin plate with a thickness of ~120 nm for small-angle RSXS observation was prepared by the focused ion beam (FIB) thinning method (SMI3200; Seiko Instruments Inc., Japan). The sample was affixed with carbon contacts on a Si substrate with a square hole of 10 × 10 μm<sup>2</sup>. Small-angle RSXS measurements were carried out at BL-16A. An in-vacuum CCD camera (2024 × 2024 pixels, Roper Industrial Inc.), positioned downstream of the sample, was used to record the RSXS intensity.

## 3 Results and Discussion

Figure 1(a) shows the experimental setup for the transmitted small-angle RSXS measurements. Figure 1(b) displays the CCD image measured at 195 mT with the circularly polarized soft x-ray of 577 eV at 80 K. An application of magnetic fields induces higher harmonic magnetic peaks. Finally, 7 spots are discerned with  $q = \pm 0.052, \pm 0.107, \pm 0.154$  and  $0.207 \text{ nm}^{-1}$  at 195 mT [see Figs. 1(b,c)]. Magnetic-field dependence of  $q$  is consistent with a predicted theoretical curve based on the chiral XY-spin model including the correction of the demagnetizing field, which strongly depend on the shape of the specimen and is proportion to the magnetization. Near the critical temperature ( $T_C \sim 119.6 \text{ K}$ ), the

dependence deviates from the theoretical curve with the correction. Additionally, according to the temperature dependence of FWHM of  $q$ , the magnetic correlation of helix becomes short-range order (SRO) above  $T_C$ . We are able to define two kinds of phases; Helix <  $T_C \sim 119.6 \text{ K}$  < SRO helix <  $T_{SRO} \sim 121.6 \text{ K}$ . In the SRO phase, it indicates that FM order and short-range modulation due to DM interaction emerges.

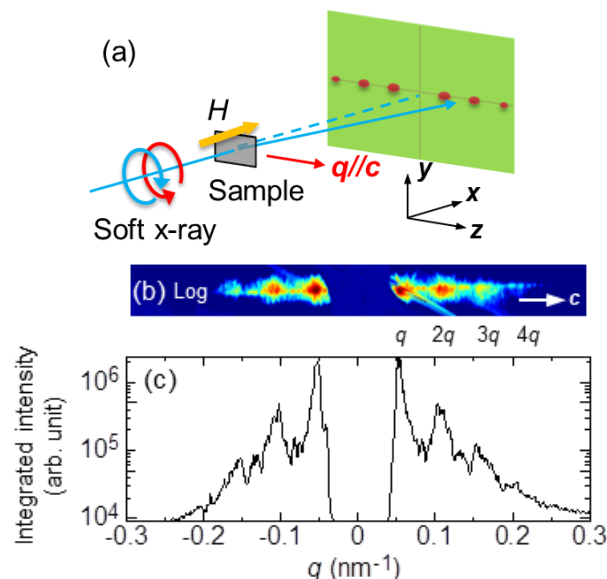


Fig. 1: (a) A transmission-type setup for small-angle RSXS. (b,c) CCD images of small-angle RSXS, and the profiles of it for CSL phase, respectively.

## References

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