

## ***In situ* synchrotron-radiation photoemission spectroscopic study of strain-controlled $\text{La}_{0.6}\text{Sr}_{0.4}\text{MnO}_3$ thin films**

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### **Introduction**

Hole-doped perovskite manganese oxides  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$  (LSMO) have attracted much attention because of their interesting magnetic and electronic properties such as colossal magnetoresistance, half-metallicity, and metal-insulator transition. The extensive studies have demonstrated that the physical properties of these manganese oxides can be controlled by epitaxial strain.<sup>1</sup> However, little is known on the effects of epitaxial strain on the electronic structure of these films. In this study, we report *in situ* photoemission (PES) study of LSMO ( $x = 0.4$ ) thin films deposited on  $\text{LaAlO}_3$  (LAO) with lattice mismatch of  $-2\%$ ,  $(\text{LaAlO}_3)_{0.3}(\text{SrAl}_{0.5}\text{Ta}_{0.5}\text{O}_3)_{0.7}$  (LSAT) with that of  $\pm 0\%$ , and  $\text{SrTiO}_3$  (STO) with that of  $+1\%$  substrates to investigate the changes in the electronic structure of LSMO thin films.

### **Experimental**

The strain-controlled LSMO ( $x = 0.4$ ) has been epitaxially grown on LAO, LSAT, and STO substrates in a laser MBE chamber connected to a synchrotron radiation PES system at BL2C of the Photon Factory.<sup>2</sup> The film thickness was estimated to be 40 nm by monitoring the intensity of the specular spot in reflection high-energy electron diffraction (RHEED) patterns. The coherent growth of these films was confirmed by the four-circle X-ray diffraction measurements. The PES spectra were taken *in situ* with a total energy resolution of 150 meV in the photon energy range of 600–700 eV.

### **Results and Discussion**

Figure 1 shows the valence band photoemission spectra of LSMO thin films under different epitaxial strain from different substrates, where LSMO/STO and LSMO/LSAT are ferromagnetic metal, while LSMO/LAO is antiferromagnetic insulator. The valence band spectra consist of four main structures labelled as A, B, C, and D. From the band structure calculation<sup>3</sup>, the two prominent structures A and B are assigned to O  $2p$  states, while a shoulder structure C and a broad hump structure D closest to the Fermi level ( $E_F$ ) are assigned to Mn  $3d t_{2g}$  and  $e_g$  states, respectively. At the first glance, the overall valence band structure does not seem to change with epitaxial strain. However, closer look reveals that the peak position

of the  $e_g$  state shifts to the higher binding energy side by 100 meV from LSMO/LSAT (strain free) to LSMO/LAO (compressive strain). The shift of  $e_g$  state is more clearly seen in the expansion near  $E_F$  (inset). The energy shift of  $e_g$  state in LSMO/LAO is explained by the energy splitting of degenerated  $e_g$  state to  $d(x^2-y^2)$  and  $d(3z^2-r^2)$  states due to the Jahn-Teller distortion induced by the compressive epitaxial strain.<sup>3</sup> The results suggest that the resultant formation of energy gap at  $E_F$  is responsible for the strain-induced metal-insulator transition in the strain-controlled LSMO.

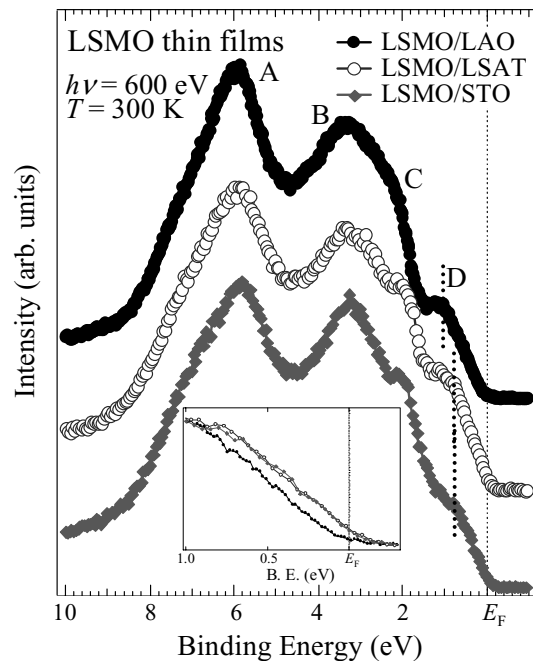


Fig. 1: The valence band spectra of LSMO thin films on various substrates. Inset shows the spectra near  $E_F$ .

### **References**

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