

Critical phenomena of lattice and charge systems in $\beta\text{-Na}_{1/3}\text{V}_2\text{O}_5$

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Introduction

A series of beta-vanadium bronzes ($\beta\text{-A}_{1/3}\text{V}_2\text{O}_5$) exhibits some exotic physical properties under high pressure (HP).

A basic structure is common to all $\beta\text{-A}_{1/3}\text{V}_2\text{O}_5$ series, $P2_1/a$ space group. In spite of the common structure, physical properties are quite different and depend on a carrier number. When the A site is occupied by an ion having 1+ valence ($\text{Na}^+\dots$), the system shows superconductivity under HP [1]. On the other hand, when the A site is occupied by an ion having 2+ valences ($\text{Sr}^{2+}\dots$), the system shows devil's flower-like behaviours under HP [2]. However, most of the samples ($A = \text{Na}, \text{Sr}$ and so on) commonly show a charge ordering with $b' = 6b$ unit cell at low temperature in spite of a different modulation of the charge instability above metal-insulator (M-I) transition at T_c [2].

To study an intrinsic difference among $\text{A}_{1/3}\text{V}_2\text{O}_5$ system, we have particularly focused upon critical phenomena of lattice and charge systems. A critical property of a sample well reflect a universal class where the sample belonging to. In this time, we will report results of a measurement of the critical phenomena of $\text{Na}_{1/3}\text{V}_2\text{O}_5$.

Experimental

Experiment was carried out at BL-4C of Photon Factory. For separating information of lattice (atomic shift) and charge, we used a resonant x-ray scattering (RXS) technique. An RXS intensity at the main edge of V-K absorption edge (around 5.47 keV) well reflects the charge disproportionation at V sites.

Results

Figure 1 shows energy spectra of 0 10/6 0 reflection measured at 100 K and 128.4 K. A fluorescence spectrum is also shown by a solid line. On-resonant intensities around $E = 5.485$ keV rapidly decreases rather than off-resonant intensities around $E = 5.44$ keV as the temperature increases near the critical temperature, $T_c = 128.96$ K. It means that the atomic shift is a primary order parameter, while the charge order is a secondary order parameter. Because the on-resonant intensity well corresponds to the charge ordering, while the off-resonant intensity corresponds to the atomic shifts.

Figure 2 shows a temperature dependence of the intensity measured at $E = 5.482$ keV (on resonant) and $E = 5.44$ keV (off resonant). Dotted curves represent the

fitting results using $f(T) = ((T_c - T)/T_c)^{2\beta(\beta')}$. The critical region was set as $126.5 < T < 128.5$ for eliminating a effect of a diffuse scattering. Finally, we obtained the critical exponent $\beta = 0.282(47)$ for the off-resonant data and $\beta' = 0.559(41)$ for the on-resonant data. The relationship between β and β' is $\beta' \sim 2\beta$, where it obeys a naive Landau prediction for the secondary order parameter.

It is suggested that the M-I transition at T_c of $\text{A}_{1/3}\text{V}_2\text{O}_5$ system is driven by a $b^* = 1/6$ lattice instability. The charge modulation tends to adjust the lattice modulation.

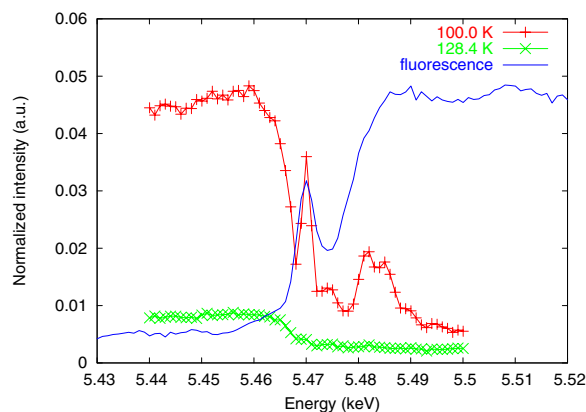


Fig. 1. The RXS and the fluorescence spectrum of $\beta\text{-Na}_{1/3}\text{V}_2\text{O}_5$.

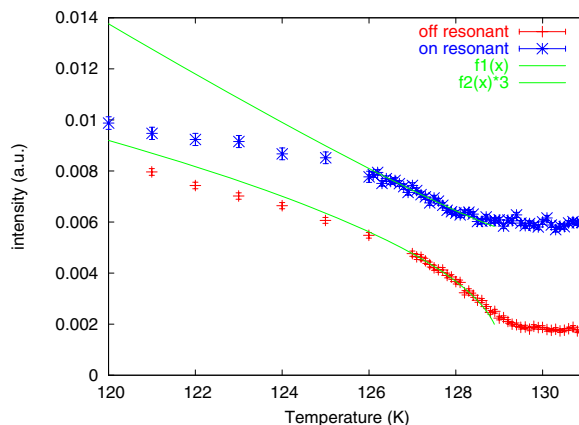


Fig. 2. Temperature dependence of the intensity measured at the on-resonant ($E = 5.482$ keV) and the off-resonant ($E = 5.44$ keV) energies.

References

- [1] T. Yamauchi *et al.*, Phys Rev. Lett. **89**, 057002 (2002).
- [2] T. Yamauchi *et al.*, private communications.