# Rotation of orbital stripes in $\operatorname{Pr}\left(\mathbf{S r}_{0.1} \mathbf{C a}_{0.9}\right)_{2} \mathbf{M n}_{2} \mathbf{O}_{7}$ 

Yusuke. TOKUNAGA ${ }^{1,2^{*}}$, Reiji. KUMAI ${ }^{3}$, Takahisa. ARIMA ${ }^{2,4}$, Yoshinori. TOKURA ${ }^{1,2,3,5}$<br>${ }^{1}$ ERATO-MF, Japan Science and Technology Agency, Tsukuba 305-0801, Japan<br>${ }^{2}$ ERATO-SSS, Japan Science and Technology Agency, Tsukuba 305-0801, Japan<br>${ }^{3}$ CERC, AIST, Tsukuba 305-0801, Japan<br>${ }^{4}$ IMRAM, Tohoku University, Sendai 980-8577, Japan<br>${ }^{5}$ Department of Applied Physics, University of Tokyo, Bunkyo-ku, Tokyo 113-8656, Japan

## Introduction

In half-doped manganites, charge- and orbital- ordered (CO-OO) state is widely realized. Below the CO-OO transition temperature $T_{\mathrm{c} 0}$, charges are ordered in a checkerboard type, forming alternating $\mathrm{Mn}^{3+} / \mathrm{Mn}^{4+}$ sites. Simultaneously, the $e_{\mathrm{g}}$ orbitals in $\mathrm{Mn}^{3+}$ sites are aligned in zigzag chain type. Because formations of orbital stripes and zigzag chains dominates the anisotropic charge dynamics in the CO-OO plane, the change in their direction connects to the control of the in-plane anisotropy of the optical and electrical conductivity. We have investigated the CO-OO states in a half-doped bilayered manganite $\operatorname{Pr}\left(\mathrm{Sr}_{0.1} \mathrm{Ca}_{0.9}\right)_{2} \mathrm{Mn}_{2} \mathrm{O}_{7}$.

## Results

Through single-crystal X-ray diffraction measurements we have determined the lattice structures of $\operatorname{Pr}\left(\mathrm{Sr}_{0.1} \mathrm{Ca}_{0.9}\right)_{2} \mathrm{Mn}_{2} \mathrm{O}_{7}$ including superstructures affected by CO-OO: There exist two CO-OO phases in this material; the higher temperature phase (CO1) and the lower temperature one (CO2). Space groups are Amam (with $a=5.410, b=5.462, c=19.277 \AA$ at 405 K ) for charge disordered phase $\left(T>T_{\mathrm{cO}}\right.$ ), Pbnm (with $a=5.412$, $\mathrm{b}=10.921, c=19.234 \AA$ at 330 K$)$ for $\mathrm{CO} 1\left(T_{\mathrm{CO} 1}>T>T_{\mathrm{CO} 2}\right)$, and $A m 2 m$ (with $a=10.812, b=5.475, c=19.203 \AA$ at 295 K) for $\mathrm{CO} 2\left(T<T_{\mathrm{CO} 2}\right)$, respectively. Note that the lattice structure is orthorhombic even above $T_{\text {coI }}$ because of tilt of $\mathrm{MnO}_{6}$ octahedra around the $a$-axis. The orthorhombic distortion couples with the CO-OO to determine the direction of orbital stripes in the bilayeres. This is distinct from the case of the wide-bandwidth analogue, $\mathrm{LaSr}_{2} \mathrm{Mn}_{2} \mathrm{O}_{7}$, which retains the tetragonal (I4/mmm) symmetry above the CO-OO transition temperature.

Figures 1(a) and (b) show the synchrotron X-ray oscillation photograph around the reflection 413 taken by using an imaging plate. Above $T_{\text {col, }}$ there is no additional spot except for fundamental reflections anticipated for orthorhombic Amam unit cell. In the higher temperature CO1 phase, superlattice spots appear at $u v \pm 1 / 2 w$ ( $u, v, w$ are integers) (Fig. 1(a)), indicating orbital zigzag-chains run along the $b$-axis, as depicted in Fig. 1(c). When the crystal is cooled below $T_{\text {co2 }}$, the superlattice spots along the $b^{*}$-axis disappear and instead emerge along the $a^{*}$-axis, i.e. at $u \pm 1 / 2 v w$ (Fig. 1(b)). This means that directions of orbital stripes and zigzag chains rotate by 90 degrees at $T_{\mathrm{C} 22}$ as depicted in Fig. 1(d),
while retaining the underlying lattice orthorhombic cell essentially intact.

It is also revealed that the rotation of orbital stripes on the orthorhombically distorted lattice breaks the symmetry of the space inversion and gives rise to an electrically polarized state. The charge polarized state is also confirmed by the large second harmonic generation susceptibility in the low-temperature $\mathrm{CO}-\mathrm{OO}$ state [2].

(c)



Fig. 1 X-ray oscillation photographs near the 413 reflection in $\operatorname{Pr}\left(\mathrm{Sr}_{0.1} \mathrm{Ca}_{0.9}\right)_{2} \mathrm{Mn}_{2} \mathrm{O}_{7}$ obtained by an imaging plate at 340 K (a), and 280 K (b), respectively. Red arrows indicate superlattice reflections. (c)(d) Schematics of charge and orbital ordering patterns at (c) $T_{\mathrm{CO} 1}>T>T_{\mathrm{CO} 2}$ (orbital stripes run along the $a$-axis) and (d) $T<T_{\text {c02 }}$ (orbital stripes run along the $b$-axis), respectively.

## References

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* yusuke-tokunaga@aist.go.jp

