

Electronic Structure of Delafossite Oxides  $\text{CuCr}_{1-x}\text{Mg}_x\text{O}_2$ 

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## 1 Introduction

Delafossite oxides  $\text{CuMO}_2$  ( $M$  = metal element) have various interesting physical properties both in fundamental and applicational terms;  $\text{CuAlO}_2$  is the first  $p$ -type transparent oxide semiconductor [1], for example, and  $\text{CuFeO}_2$  is a typical multiferroic compound [2]. The delafossite oxides have also considerable potential for thermoelectric materials [3] because of its layered structure of edge-shared  $\text{MO}_6$  octahedrons, which is the same as thermoelectric  $\text{NaCoO}_2$  [4]. Hole-doped  $\text{CuCr}_{1-x}\text{Mg}_x\text{O}_2$  is one of such candidates; in  $\text{CuCrO}_2$ ,  $3d^3$  electrons of the  $\text{Cr}^{3+}$  ions under the  $O_h$  local symmetry fill up the narrow  $\text{Cr } 3d t_{2g}$  band and thus steep density of states (DOS) at the Fermi level ( $E_F$ ) may be realized near the  $t_{2g}$  band edge in the hole-doped system  $\text{CuCr}_{1-x}\text{Mg}_x\text{O}_2$  with one of the highest conductivity among delafossite oxides [3]. In this paper, we briefly report our research on the valence-band electronic structure of this system.

## 2 Experiment

Polycrystalline samples of  $\text{CuCr}_{1-x}\text{Mg}_x\text{O}_2$  ( $x=0, 0.02, 0.03$ ) were prepared by the standard solid-state reaction [3]. Photoemission measurements were performed at BL-28A of Photon Factory in KEK. The samples were fractured *in situ* in the main chamber right before measurements in ultrahigh vacuum (better than  $1.2 \times 10^{-7}$  Pa) at 300 K. The intensity of the spectra was normalized using photon current of the exit mirror. The energy resolution of about 30 meV and the  $E_F$  location were both determined by Gold Fermi edge.

## 3 Results and Discussion

Figure 1 shows the valence-band spectra of the  $x=0.02$  sample taken with the photon energies around the  $\text{Cr } 3p$ - $3d$  (Panel (a)) and the  $\text{Cu } 3p$ - $3d$  (Panel (b)) resonance energies. The intensity of the near- $E_F$  leading structures (the 1.4-eV shoulder and the 2.3-eV peak) systematically changes with increasing photon energy, as displayed in Fig. 1 (c) and (d). In Panel (c), the 1.4-eV shoulder shows a clear resonance-type behavior with the maximum at 50.0 eV, the  $\text{Cr } 3p$ - $3d$  resonance energy, whereas the 2.3-eV peak shows a typical weak anti-resonance-type behavior with a dip at the  $\text{Cu } 3p$ - $3d$  resonance energy of

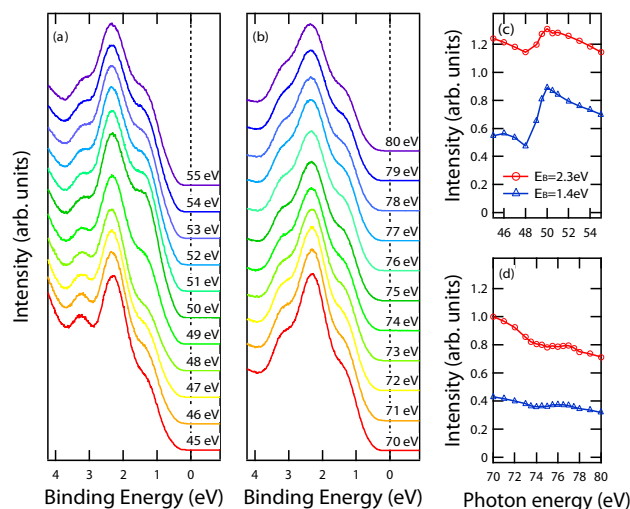


Fig. 1: Valence-band spectra of  $\text{CuCr}_{0.98}\text{Mg}_{0.02}\text{O}_2$  taken with the photon energy around (a) the  $\text{Cr } 3p$ - $3d$  resonance region and (b) the  $\text{Cu } 3p$ - $3d$  resonance region. (c) and (d) show CIS spectra of the 1.4-eV and 2.3-eV structures.

74.0 eV in Panel (d). Therefore, the  $\text{Cr } 3d$  character is dominant in the 1.4-eV shoulder and the  $\text{Cu } 3d$  character is dominant in the 2.3-eV peak. However, one can also notice that a weak resonance of the 2.3-eV peak exists at 50.0 eV and a tiny anti resonance of the 1.4 eV at 74.0 eV does. This is demonstrating that there exists sizeable hybridization between the  $\text{Cr } 3d$  and  $\text{Cu } 3d$  states via O  $2p$  states. Although the obtained energy diagram is somewhat different from previous works [5,6], it is in agreement with another previous work [7] and our band-structure calculations. It is also reasonable in terms of the difference between the O  $2p$ -to- $\text{Cr } 3d$  charge-transfer energy and the O  $2p$ -to- $\text{Cu } 3d$  charge-transfer energy [8].

## Acknowledgement

We would like to thank Prof. T. Mizokawa for useful discussions.

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