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Angle-resolved photoemission study of topological crystalline insulator Pb_{1,x}Sn_xTe

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1 Introduction

The discovery of topological insulators (TIs) triggered the search for new types of topological materials, and a recent theory predicted the existence of "topological crystalline insulators" (TCIs) in which metallic surface states are protected by point-group symmetry of the crystal structure. Such a TCI phase has been experimentally verified by angle-resolved photoemission spectroscopy (ARPES) experiments for narrow-gap IV-VI semiconductors SnTe [1], Pb_{0.6}Sn_{0.4}Te [2], and Pb_{0.77}Sn_{0.23}Se [3]. In those materials, the topological surface states measured on the (001) surface consist of double Dirac cones located at momenta slightly away from $\bar{\mathbf{X}}$ point in the (110) mirror plane of the crystal. Since there are four $\bar{\mathbf{X}}$ points on the boundary of the first surface Brillouin zone (BZ), there are a total of four Dirac cones within it. This is distinct from the three-dimensional (3D) TIs whose surface states are characterized by an odd number of Dirac cones. In contrast to the double Diraccone signature observed in the TCI phase, the ARPES measurements for isostructual PbTe [1] and Pb₀₈Sn₀₂Te [2] revealed the absence of any surface states, which strongly suggests a trivial-to-nontrivial topological quantum phase transition (QPT) in the solid-solution system Pb_{1-x}Sn_xTe. We have performed comprehensive ARPES studies of $Pb_{1-x}Sn_xTe$ with various Sn compositions (x = 0.0, 0.2, 0.3, 0.5, 0.6, 0.8, and 1.0) to clarify how the surface and bulk electronic states evolve as a function of Sn composition x, which would be useful for establishing practical understanding of TCIs [4].

2 Experiment

High-quality single crystals of the two end-member samples SnTe (x = 1.0) and PbTe (x = 0.0) were grown by the modified Bridgeman method, and the solid-solution samples (x = 0.2, 0.3, 0.5, 0.6, and 0.8) were prepared by a vapor transport method. ARPES measurements were performed with a VG-Scienta SES2002 electron analyzer with a tunable synchrotron light at the beamline BL28A at Photon Factory (KEK). We used circularly polarized light of 75-100 eV. The energy and angular resolutions were set at 10-30 meV and 0.2° , respectively.

3 Results and Discussion

Figures 3(a) and (b) show the near- $E_{\rm F}$ ARPES intensity and corresponding EDCs, respectively, taken at hv = 78.5 eV. From photon energy dependence (75-100 eV), we have elucidated that hv ~78.5 eV hits to the right k_z value to probe the L point of the bulk BZ projected to the \bar{X}

point of the surface BZ. To elucidate the evolution of the surface states and bulk VB as a function of x, we have measured all the samples with hv = 78.5 eV along the k cut to cross the bulk VB top. At x = 1.0, we observe a highly dispersive holelike band which is attributed to an admixture of bulk and surface bands with dominant contribution from the surface state near $E_{\rm F}$. As clearly seen, the holelike band still touches to the $E_{\rm E}$ at x = 0.3, suggesting the existence of metallic surface state. However, at x = 0.2 the holelike band is obviously away from $E_{\rm F}$, indicating the surface state vanishes at $x_{\rm c} \approx 0.25$. From the x dependence of the energy location of VB maximum (not shown), we also found that VB maximum approaches $E_{\rm E}$ as x is reduced from 1.0 toward x, and then it moves away from $E_{\rm F}$ when x passes $x_{\rm c}$, suggesting the band inversion at $x_{\rm c}$. With this band inversion, we conclude that the phase transition at $x_c \approx 0.25$ is indeed of topological origin, separating the TCI and trivial phase.



Fig. 1: (a), (b) x dependence of near- $E_{\rm F}$ ARPES intensity and corresponding EDCs, respectively, measured along the $k_{\rm y}$ cut across the bulk VB top that can be accessed with hv = 78.5 eV.

References

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