

An x-ray fluorescence holographic study on a  $\text{Bi}_2\text{Te}_3\text{Mn}_{0.1}$  topological insulatorShinya Hosokawa<sup>1,\*</sup>, Naohisa Happo<sup>2</sup>, Kouichi Hayashi<sup>3</sup>, A Ohnishi<sup>4</sup>, M Kitaura<sup>4</sup>,  
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## 1 Introduction

In the past,  $\text{Bi}_2\text{Te}_3$  was well-known as a thermoelectric material having a large value of thermopower. Nowadays, this material has also achieved much attention as a topological insulator (TI), where the so-called Dirac electrons with an almost zero mass can conduct on the surface of this material. From this reason,  $\text{Bi}_2\text{Te}_3$  is promising as a raw material for future computer systems of very fast and energy-saving.

Recently, Sasaki and coworkers [1] found that Mn- and some other impurity(M)-doped  $\text{Bi}_2\text{Te}_3\text{Mn}_x$  alloys have an excellent TI behavior. Moreover, they found some interesting features as TI materials. Firstly, the lifetime of the TI properties is extremely long, which is very important for applications. Secondly, the Dirac electrons are not scattered by impurities (clean surface Dirac electrons), which was detected by a magnetoresistance measurement at low temperature.

Then, an important question arises from the structural point of view: Where do the impurity atoms locate in the  $\text{Bi}_2\text{Te}_3$  crystal for showing such excellent TI properties?  $\text{Bi}_2\text{Te}_3$  has a layer structure with a hexagonal form. The interlayer positions would be most plausible for the impurities, but there is no reason to exclude the possibilities of substitutions with Bi or Te atoms or interstitial positions in the layer.

X-ray fluorescence holography (XFH) is a newly developed technique that enables one to draw three-dimensional (3D) atomic images around a specific element emitting fluorescent x-rays [2]. Due to the interference between the direct incident x-rays and those scattered by the surrounding atoms, the fluorescent x-ray intensity from the emitter slightly modulates with the incident x-ray angles by about 0.1%, from which 3D images can be obtained by simple Fourier transforms without any special models. We have performed the Mn  $K\alpha$  XFH measurements on  $\text{Bi}_2\text{Te}_3\text{Mn}_{0.1}$  single crystal at room temperature. In this paper, we report atomic images around the Mn atoms. Details of the results are given elsewhere [3].

## 2 Experiment

A single crystal of  $\text{Bi}_2\text{Te}_3\text{Mn}_{0.1}$  was grown by a modified Bridgman method, where  $\text{Bi}_2\text{Te}_3$  and Mn powders were melted and crystallized in an evacuated quartz ampoule several times by slow cooling. The sample was cooled from 850°C to 550°C with a cooling rate of 10°C/h. The crystal was cut and polished so as to

have a (001) flat surface larger than  $5 \times 5 \text{ mm}^2$ . The crystallinity of the sample was examined by taking a Laue photograph, and the concentrations was checked by EPMA analysis to be the nominal value over the sample.

XFH measurements were carried out at the beamline BL6C. The sample was placed on a two-axes table of a diffractometer. The measurements were performed in inverse mode by rotating the two axes, the exit angle of  $0^\circ \leq \theta \leq 75^\circ$  in steps of  $1.00^\circ$ , and the azimuthal angle of  $0^\circ \leq \phi \leq 360^\circ$  in steps of about  $0.35^\circ$ , of the sample stage. Incident x-rays were focused onto the (001) surface of the samples. Mn  $K\alpha$  fluorescent x-rays were collected using an avalanche photodiode detector with a cylindrical graphite crystal energy-analyzer. The XFH signals were recorded at eight different incident x-ray energies from 7.0 to 10.5 keV in steps of 0.5 keV. Each scan took about 12 h. Details of the experimental setup are given elsewhere [2].

Holographic oscillation data were obtained by subtracting the background from the fluorescent x-ray intensities and normalizing them to the incident x-ray intensities. An extension of the hologram data was carried out using the crystal symmetry of the hexagonal structure and the measured x-ray standing wave (XSW) lines. From the hologram patterns, 3D atomic configuration images were reconstructed using Barton's algorithm [4] by superimposing the holograms with eight different incident x-ray energies, which can highly suppress the appearance of twin images.

## 3 Results and Discussion

Figure 1 shows examples of the Mn  $K\alpha$  hologram pattern of the  $\text{Bi}_2\text{Te}_3\text{Mn}_{0.1}$  single crystal measured at an incident x-ray energy of 8.0 keV. The radial and angle directions indicate  $\theta$  and  $\phi$ , respectively, and the magnitudes are given as the color bars beside the figure. A roughly threefold symmetry including XSW signals was observed in the holographic pattern, indicating a good quality of the sample crystal.

Figure 2 shows the reconstructed atomic image of  $\text{Bi}_2\text{Te}_3\text{Mn}_{0.1}$  single crystal on the (001) plane around the central Mn atom marked by the circle in the figure. The image intensities were normalized to that of the strongest atoms, and are shown as the color bar beside the figure. To avoid confusions on the noise level artifacts and to emphasize the middle intensity range, the color variation was chosen to range from 0.25 to 0.75. As a guide for

eyes, cross-sections of dashed lines indicate ideal positions of the neighboring atoms obtained from x-ray diffraction provided that the Mn atom is located at a Bi or Te position. In the crystal structure, all of atoms on the (001) plane are the same element, Bi or Te, depending on the  $z$  value along the  $c$  axis.

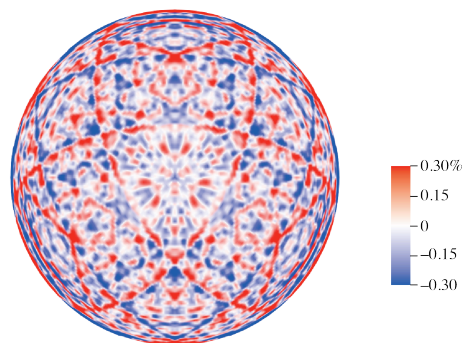


Fig. 1: Examples of Mn  $K\alpha$  hologram pattern obtained from (001) surfaces of  $\text{Bi}_2\text{Te}_3\text{Mn}_{0.1}$  single crystal measured at incident x-ray energy of 8.0 keV. The figure is taken from ref. [3].

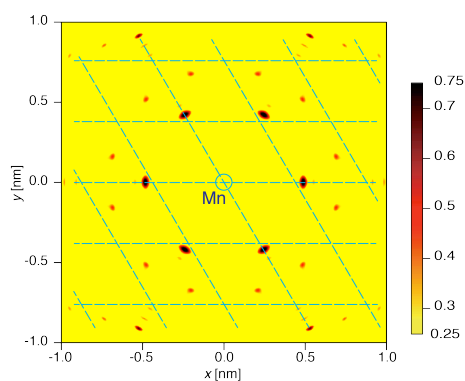


Fig. 1: Atomic images of  $\text{Bi}_2\text{Te}_3\text{Mn}_{0.1}$  single crystal on (001) plane around central Mn atoms marked by the circle. Cross-sections of dashed lines indicate ideal positions of neighboring atoms obtained from x-ray diffraction provided that the Mn atom is located at a Bi or Te position. The figure is taken from ref. [3].

As seen in Fig. 2, clear six atomic images are observed near the ideal positions of the hexagonal lattice. Thus, the Mn impurity atom does not enter an interlayer position, and would be located at a substitutional position of Bi or Te atoms in the original  $\text{Bi}_2\text{Te}_3$  crystal. Since the nearest-neighboring atoms are hardly visible (not shown in the figure) unlike usual covalent crystals, the Mn impurity atom would locate at a flexible site, and thus, it is most plausible that the Mn impurity atoms are substituted with an Bi atom. However, there is no reason to exclude the possibilities of substitutions with the Te atoms inside and surface of the layer.

The local lattice constant around the Mn impurity is slightly larger than the original  $\text{Bi}_2\text{Te}_3$  crystal. Note that the present XFH shows only images of the second neighboring Te atoms, and the other images are hardly visible (not shown). As discussed in a previous paper on mixed crystals, an impurity induces strong fluctuations in nearest-neighboring atomic positions, which results in very weak atomic images of nearest-neighboring atoms. In this material, however, further distant neighboring atoms are not observed unlike previous mixed crystals. From these findings, it would be reasonable to conclude that a strong flat clusters are formed, and the positional fluctuations between the clusters are very large.

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