

軟X線磁気円二色性による磁性ナノ構造 の研究

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SPring8 BL23-SU: 竹田幸治, 岡根哲夫, 斎藤祐児(原子力機構)

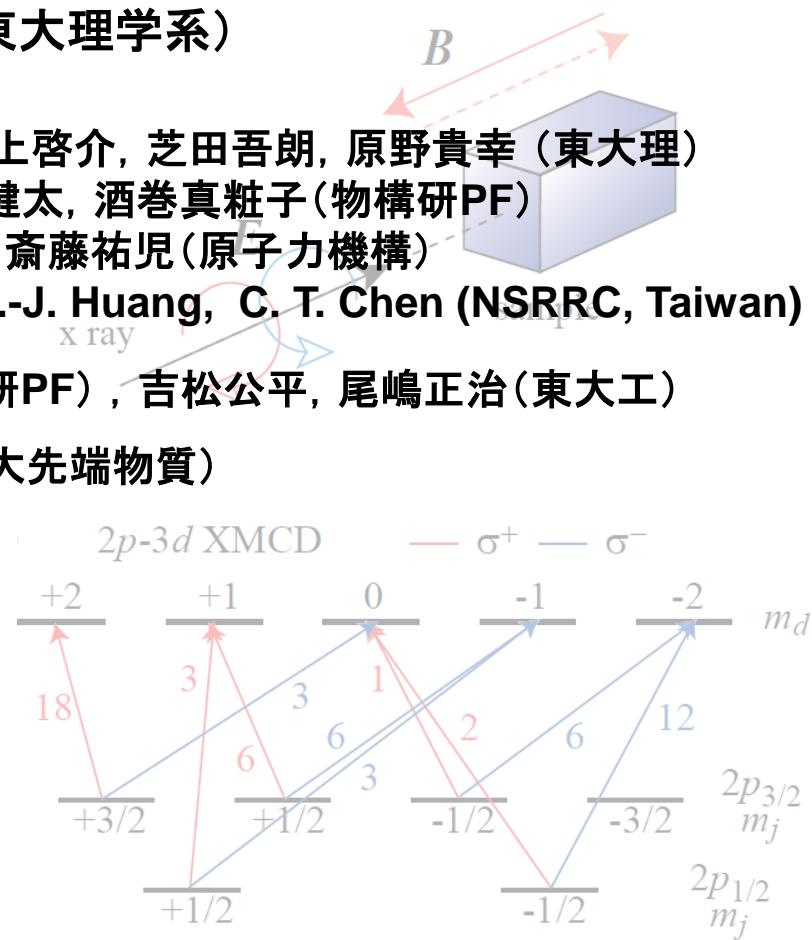
TLS BL-11A1: F. H. Chang, H.-J. Lin, D.-J. Huang, C. T. Chen (NSRRC, Taiwan)

SrRuO_3 , $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$: 組頭広志(物構研PF), 吉松公平, 尾嶋正治(東大工)

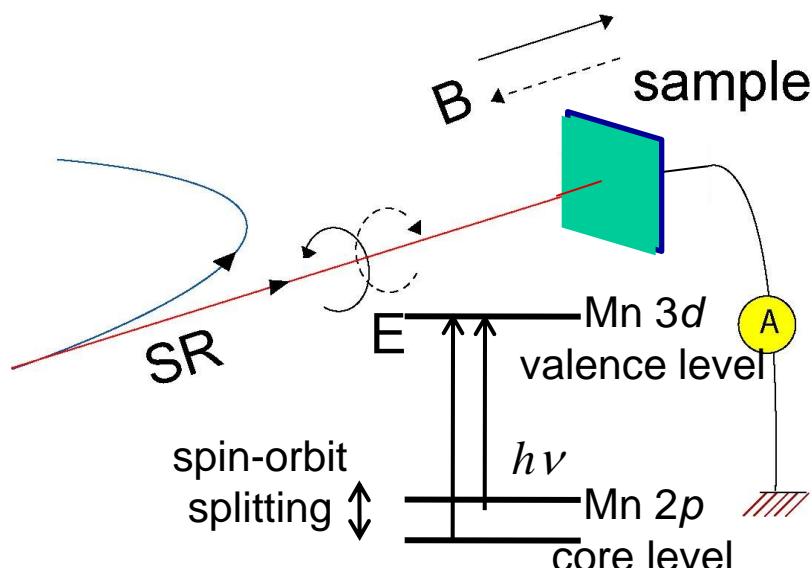
多重項計算, クラスター計算: 田中 新(広大先端物質)

概要

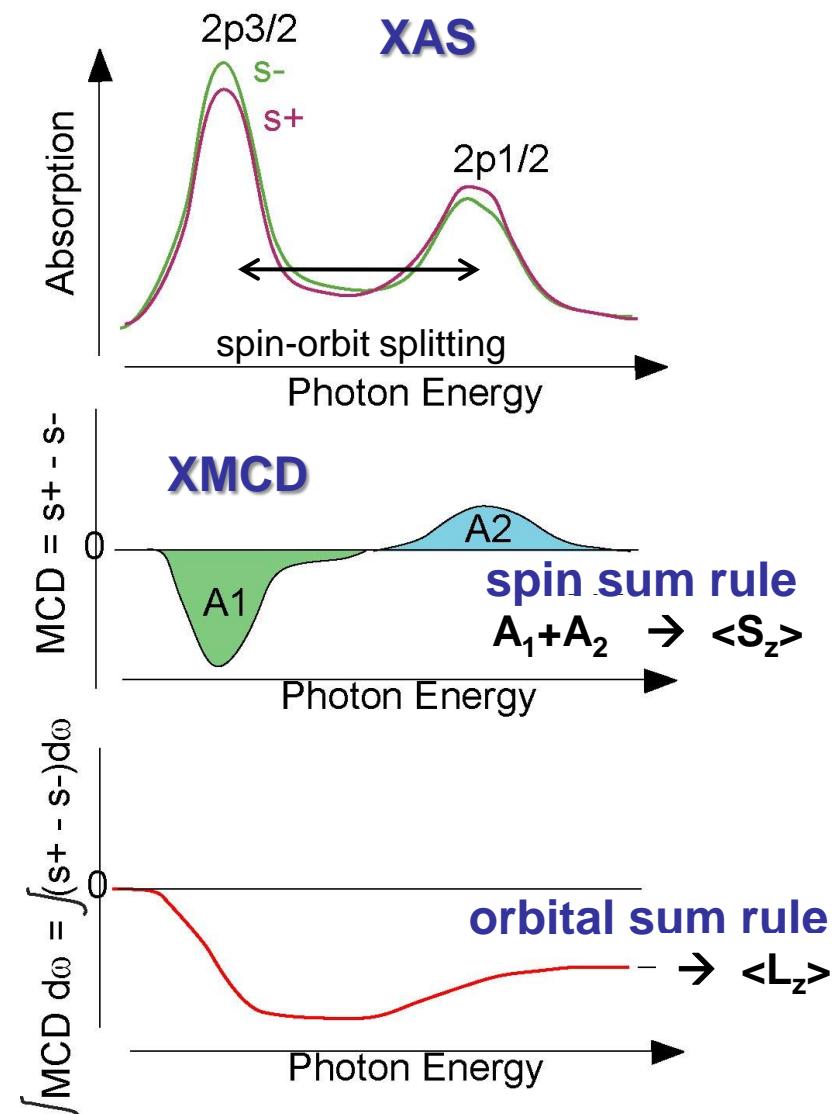
- ・ 測定原理, 特徴
- ・ これまでの研究例
 - 希薄磁性半導体
 - 酸化物薄膜
- ・ 今後の展望



X-ray magnetic circular dichroism (XMCD) in core-level x-ray absorption spectroscopy (XAS)

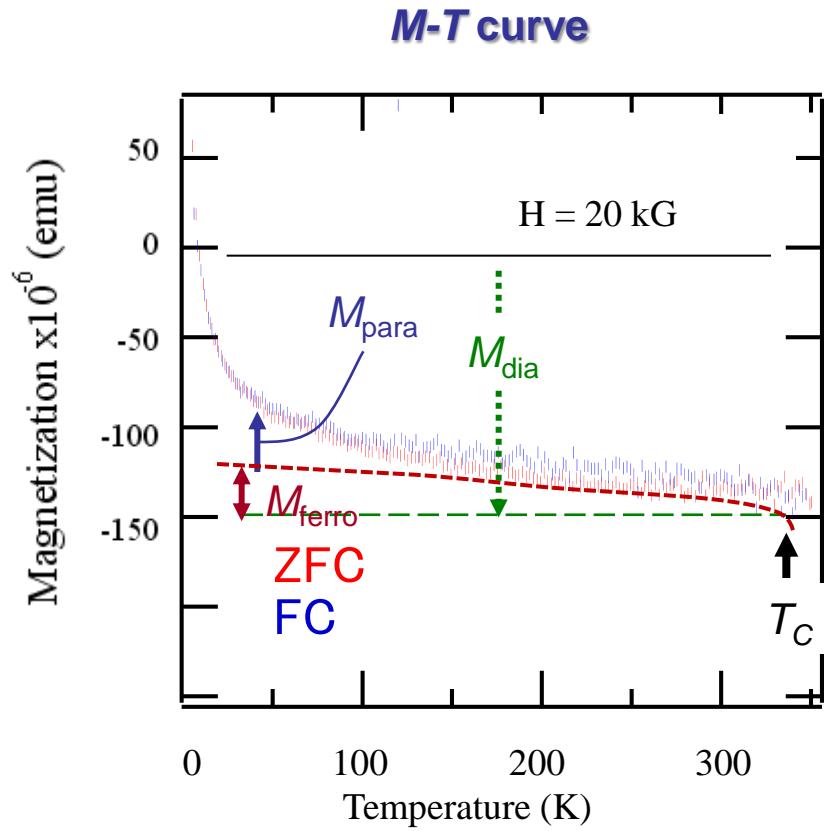
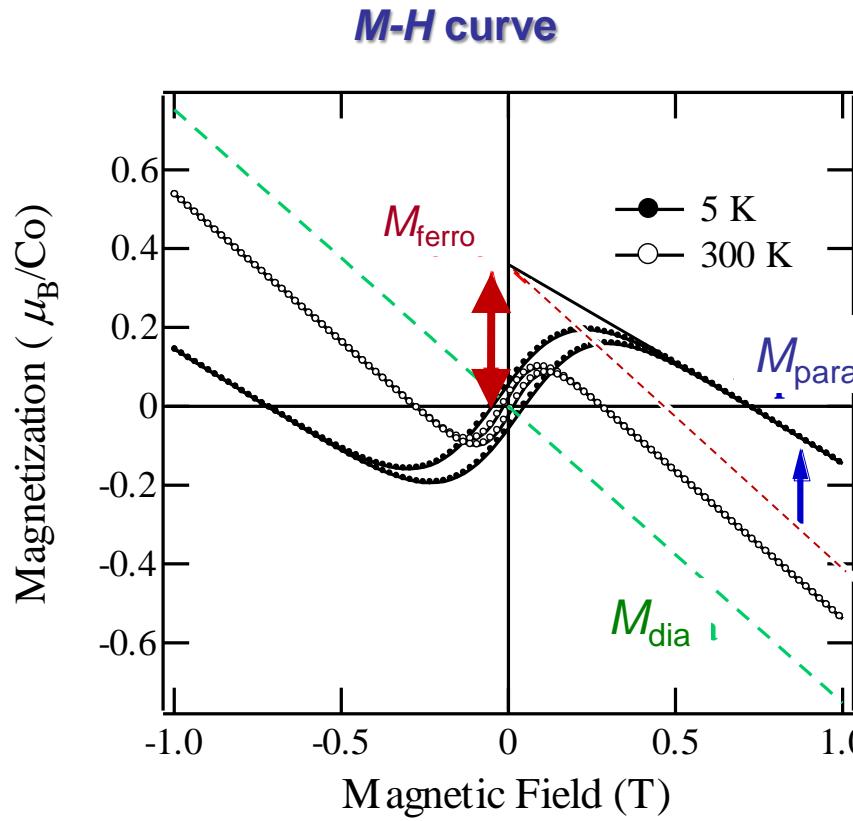


- Element-specific
- chemical state-selective
- Magnetism-selective



Ferromagnetic and paramagnetic components in magnetization and XMCD signals

SQUID data of thin film sample

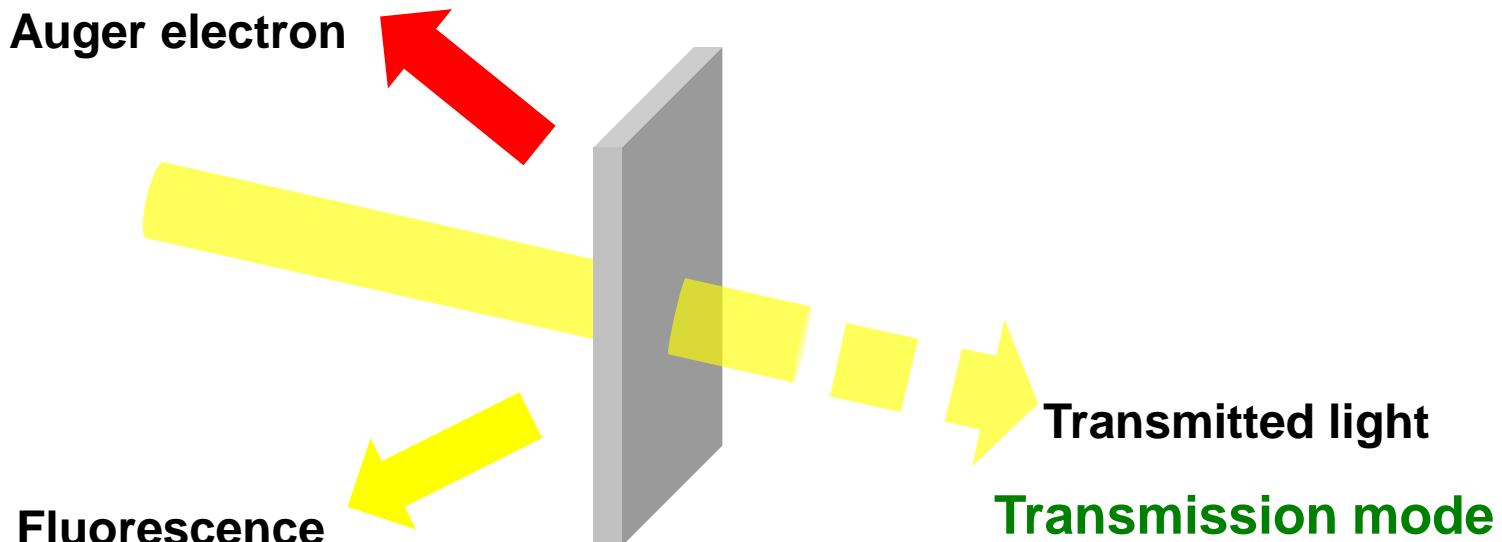


$$M = M_{\text{dia}} + M_{\text{ferro}} + M_{\text{para}} \rightarrow \text{XMCD signals}$$

Surface- and bulk-sensitive detection modes of XAS and XMCD measurements

Total electron yield (TEY) mode

probing depth ~ 3-5nm

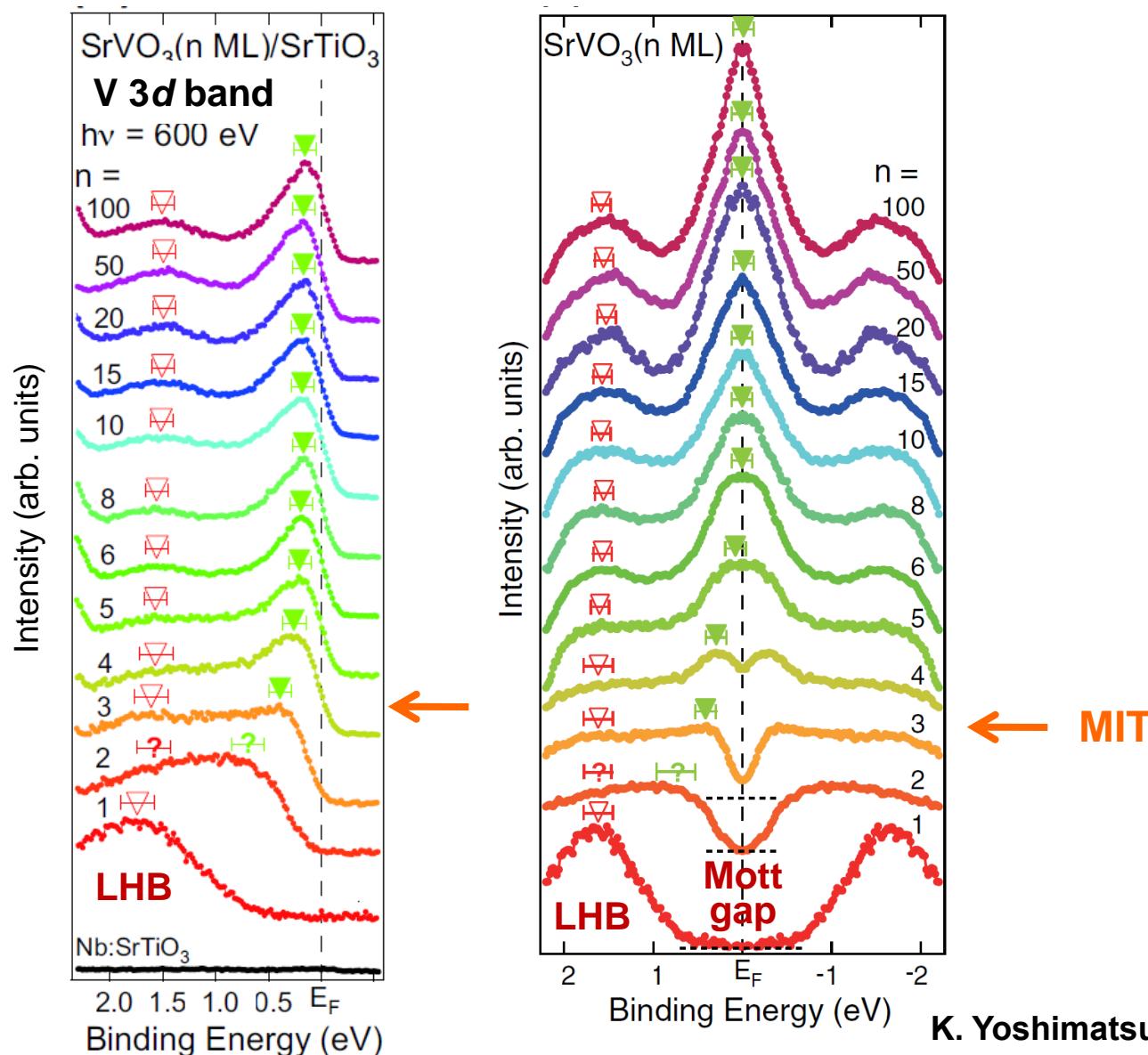


Total fluorescence yield (TFY) mode

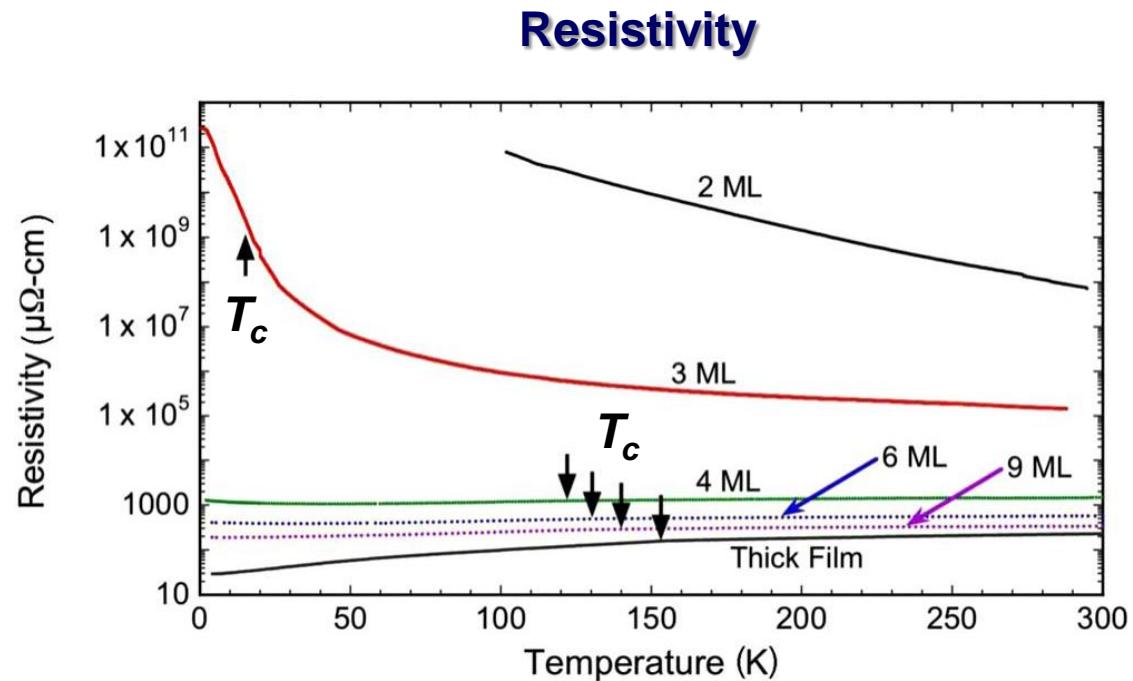
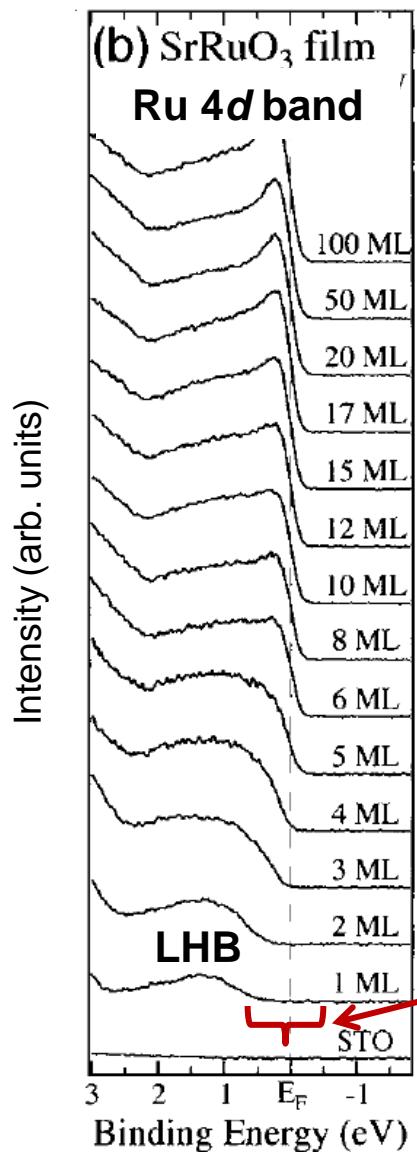
probing depth ~ 100nm ~ film thickness

Disadvantage: Self-absorption → Partial fluorescence yield (PFY) mode

Metal-to-insulator transition in SrVO_3 with decreasing film thickness



Metal-to-insulator transition in SrRuO_3 with decreasing film thickness

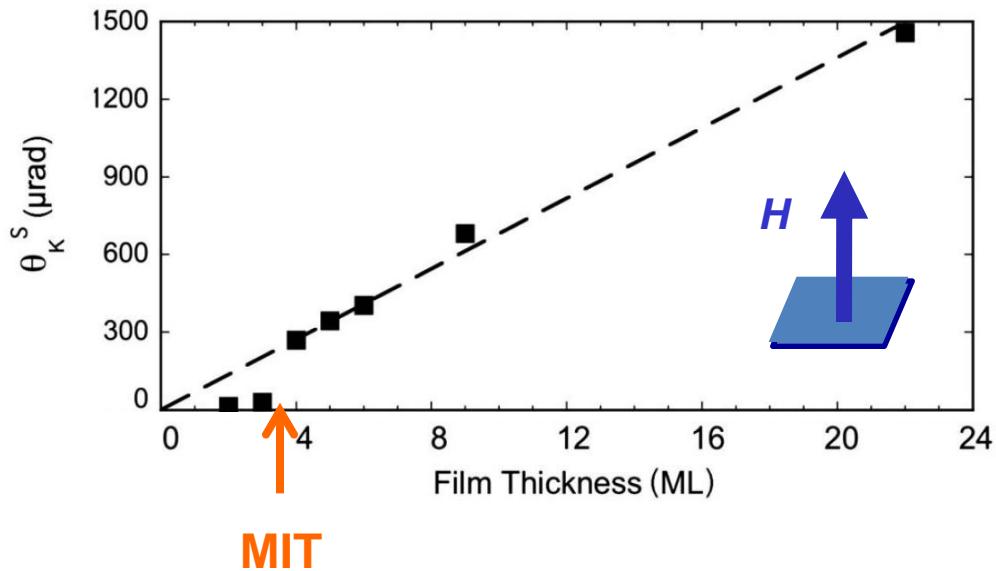


J. Xia et al., PRB '09

D. Toyota et al., APL '05

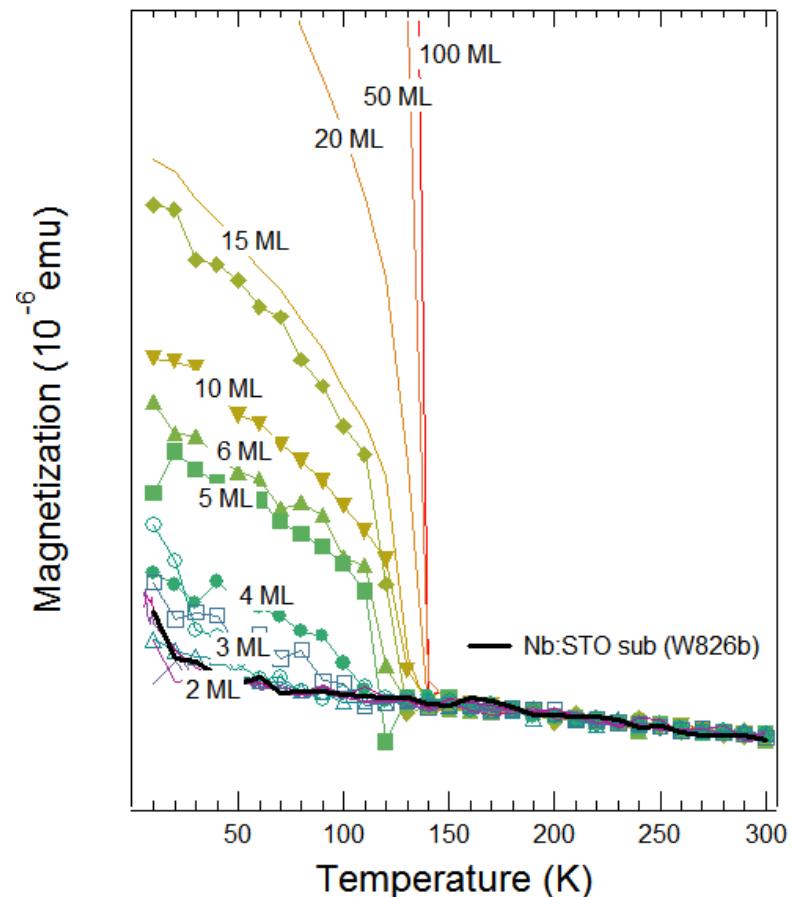
Concomitant ferromagnetic-to-paramagnetic transition in SrRuO_3 thin films

Polar Kerr effect



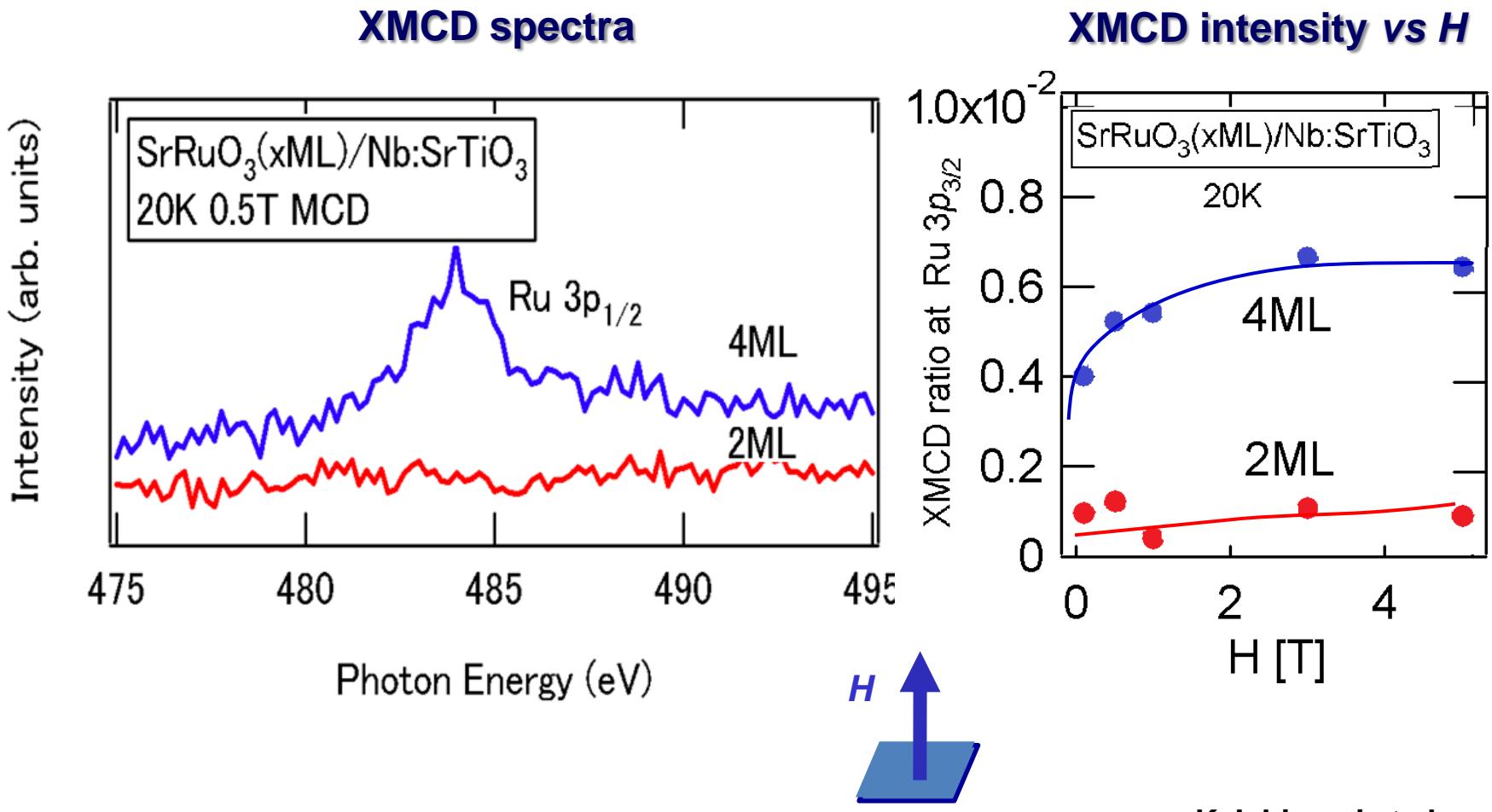
J. Xia et al., PRB '09

Magnetization measured by SQUID

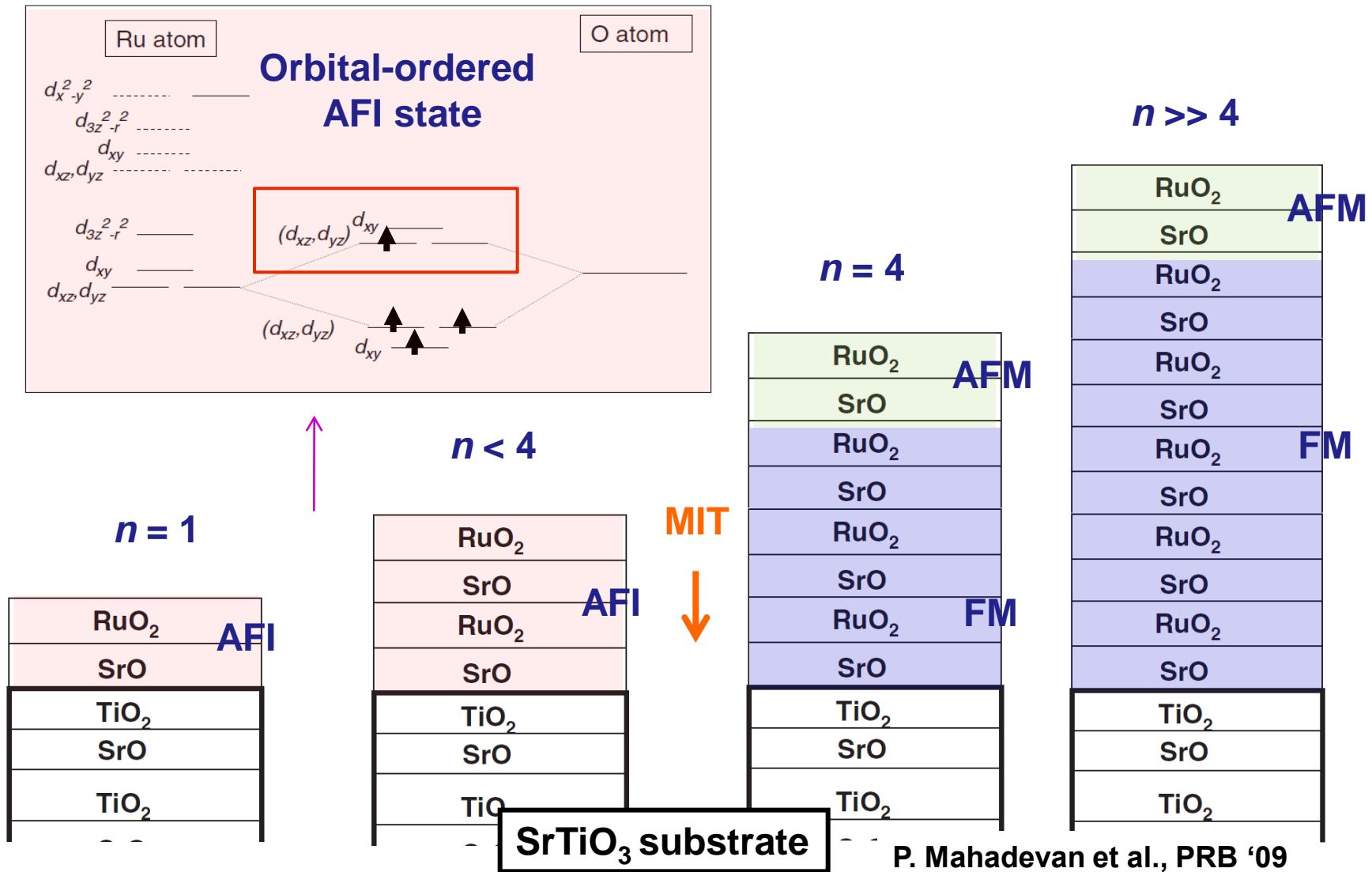


M. Takizawa et al.

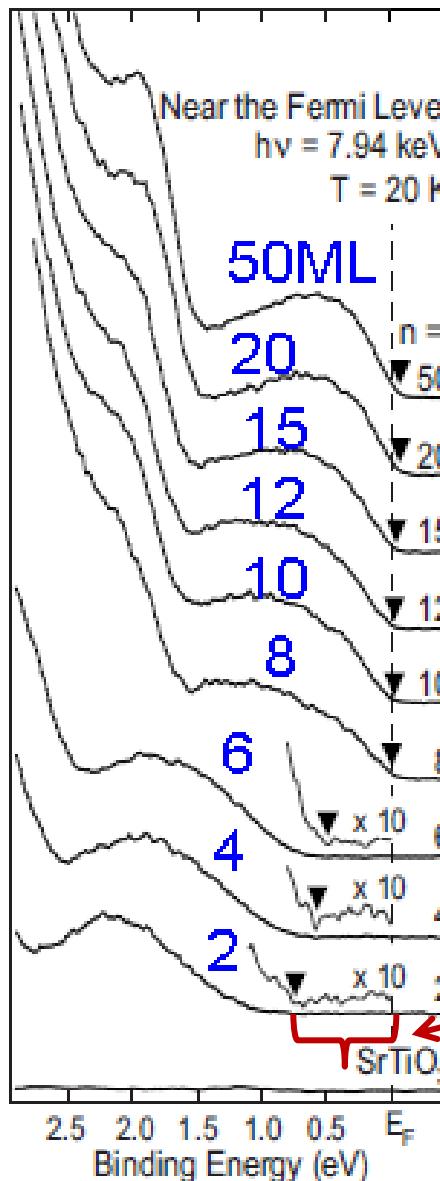
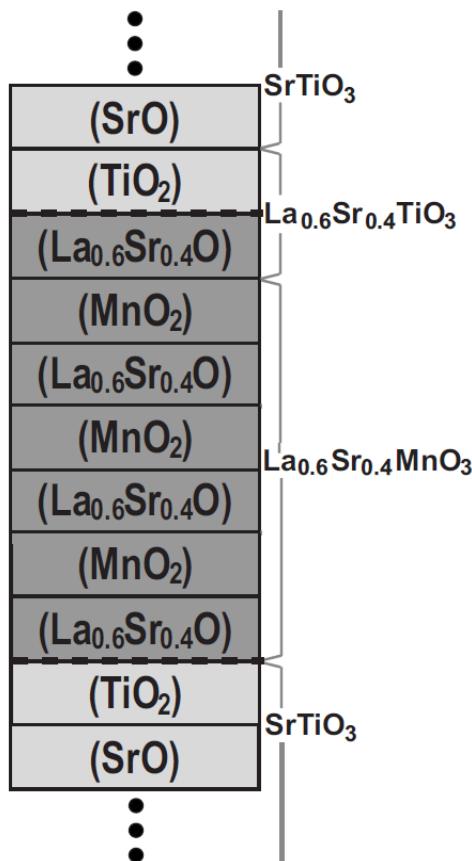
Ru 3p → 3d XMCD of SrRuO₃ thin films near critical thickness



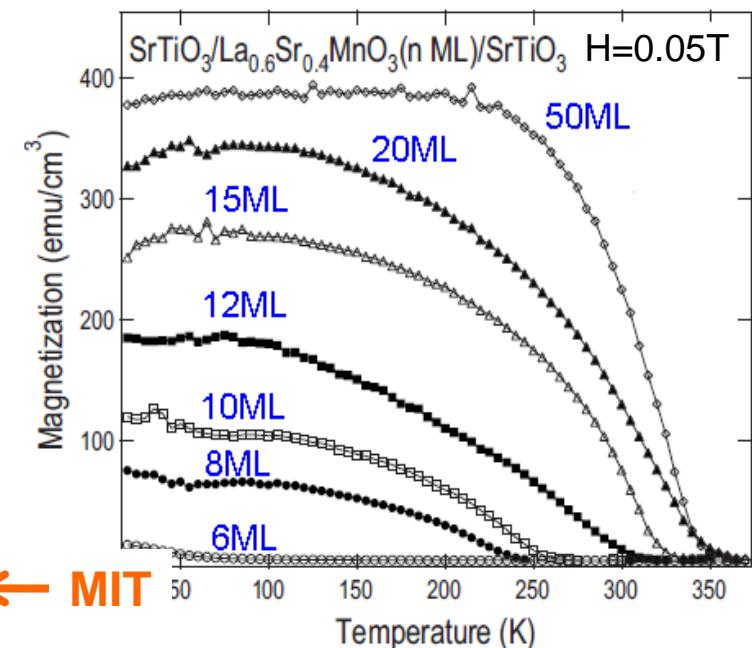
LDA+U calculation of SrRuO₃ thin films



Metal-to-insulator transition in $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ with decreasing film thickness



Magnetization measured by SQUID

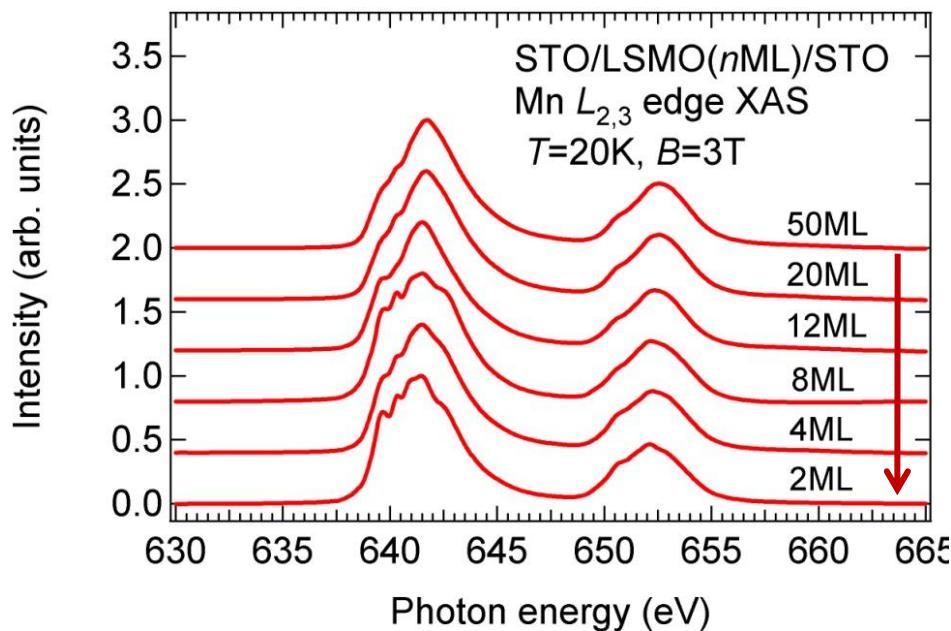


Mott gap

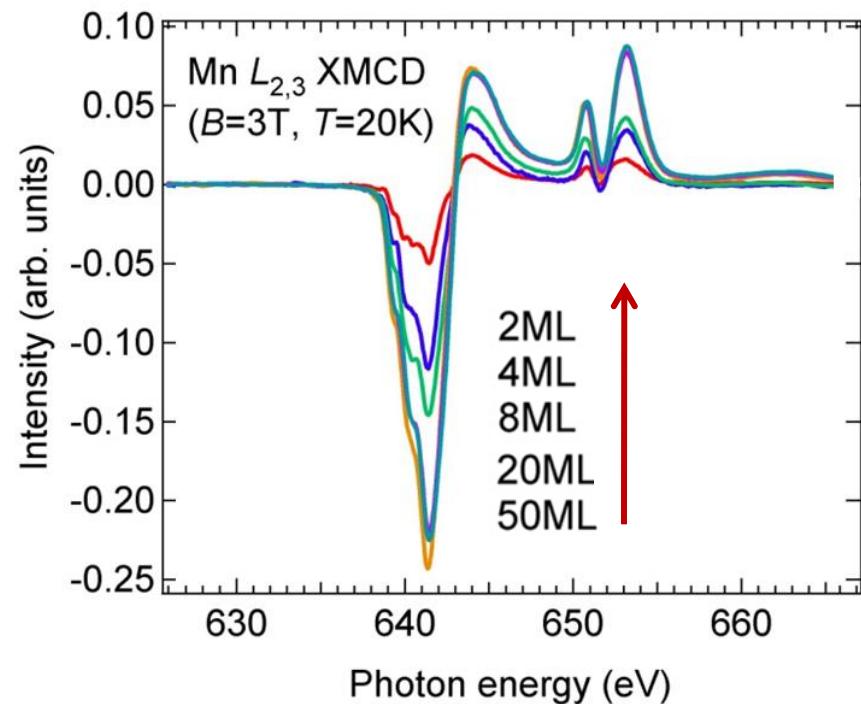


Mn $2p \rightarrow 3d$ XAS and XMCD of $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ thin films

X-ray absorption spectroscopy (XAS) spectra



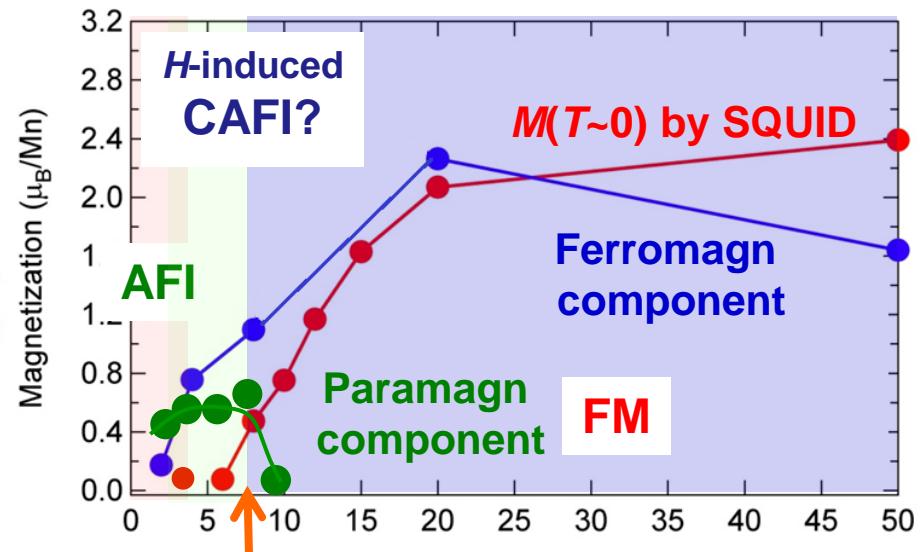
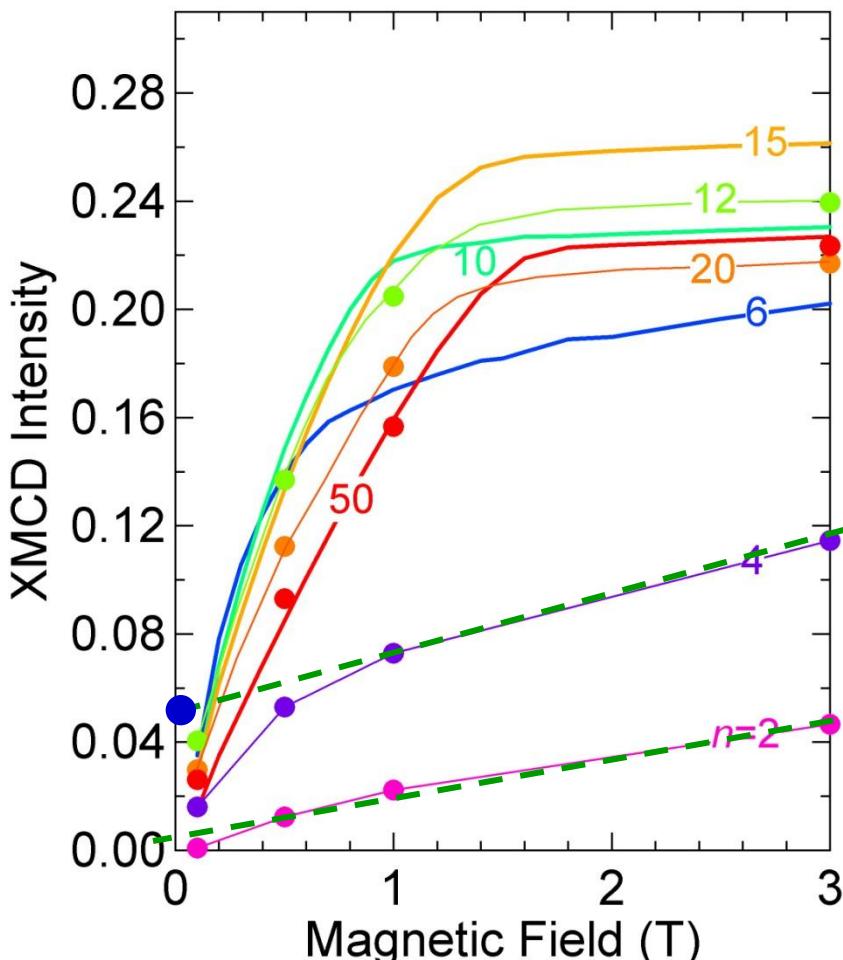
X-ray magnetic circular dichroism (XMCD) spectra



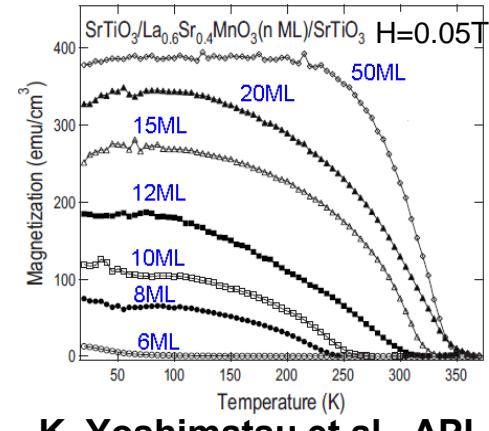
No change in $\text{Mn}^{3+}/\text{Mn}^{4+}$ ratio

Mn 2p \rightarrow 3d XMCD of $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ thin films

XMCD intensity vs H



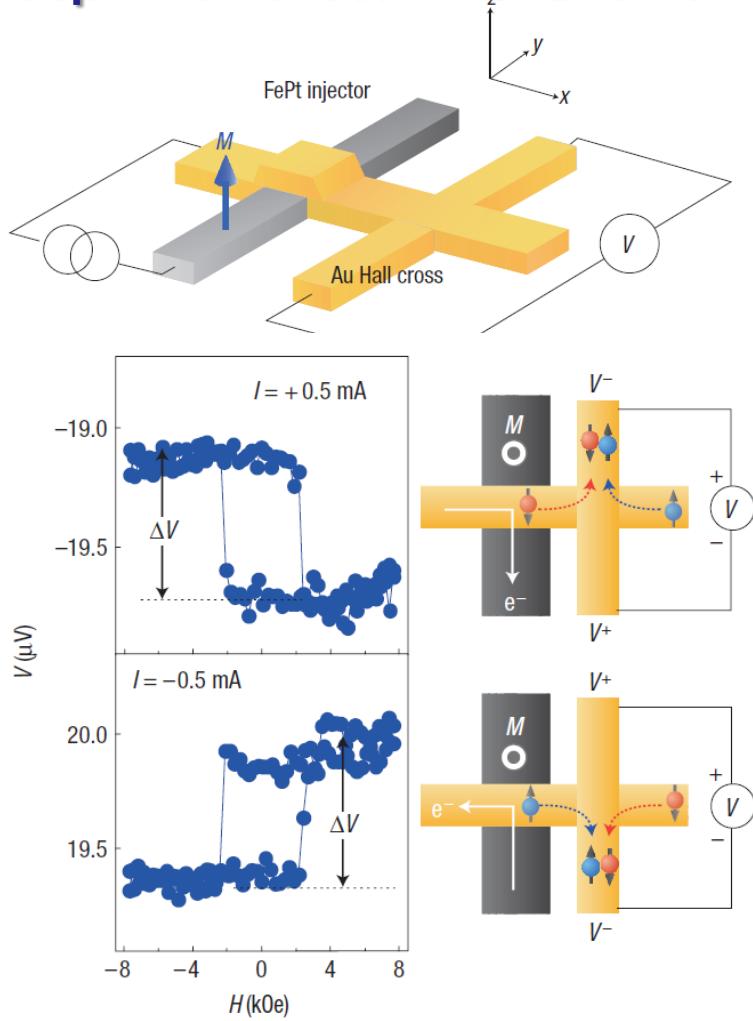
G. Shibata et al.



K. Yoshimatsu et al., APL '09

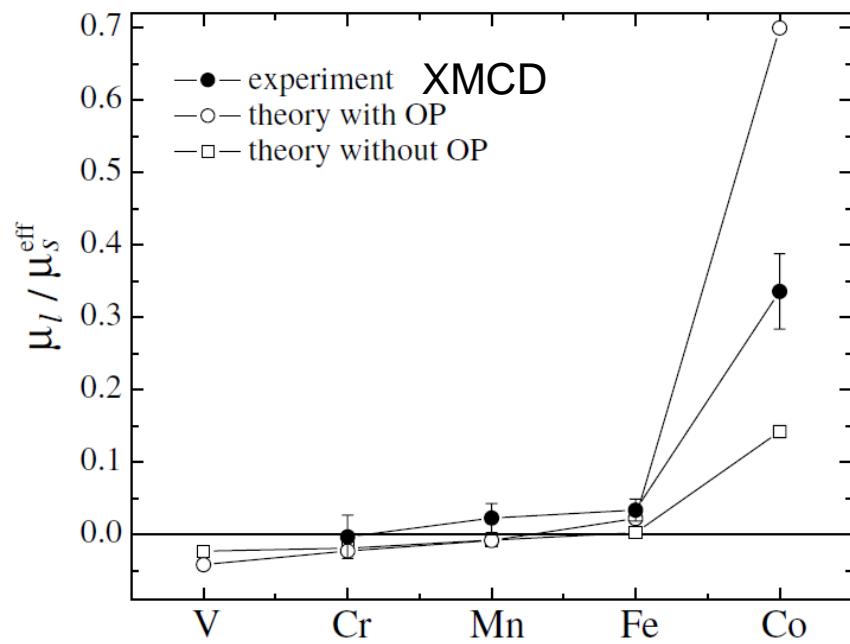
スピノンホール系におけるスピノン軌道相互作用

Giant spin Hall effect in FePt/Au devices



T Seki et al, Nat. mater. '08

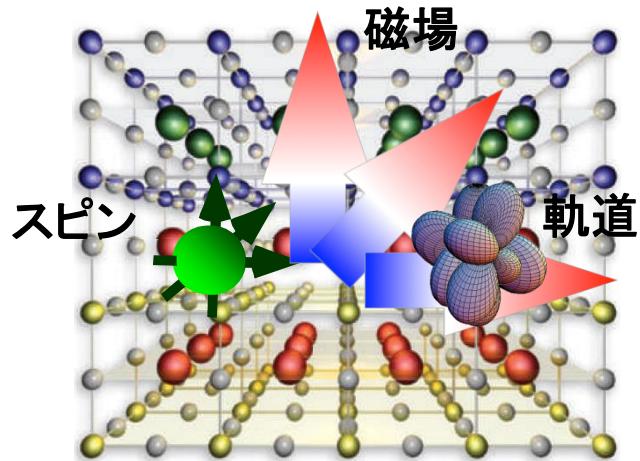
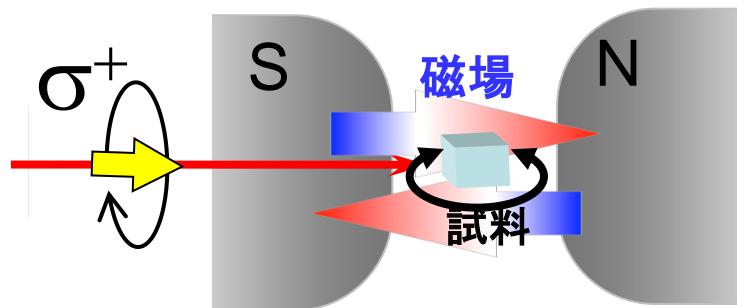
Orbital moment of TM atom in Au



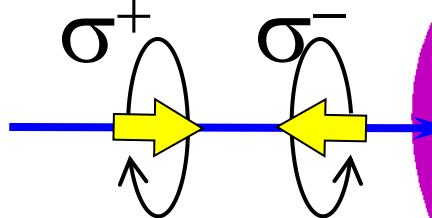
W. D. Brewer et al., PRL '04

ベクトル型マグネットと高速偏光スイッチング を用いたXMCD

従来のXMCD装置

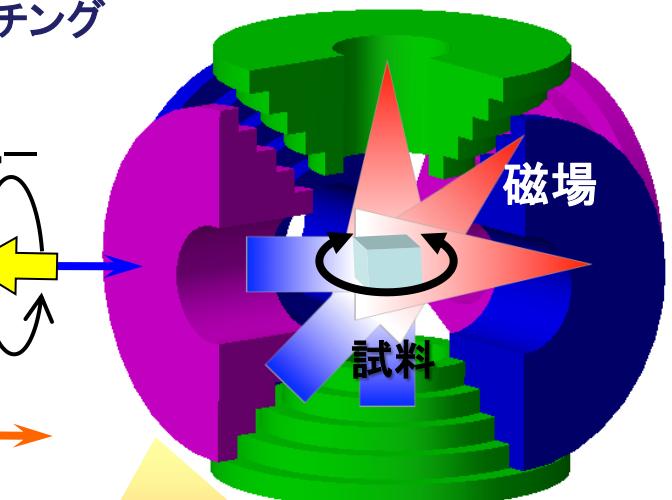


高速偏光スイッチング



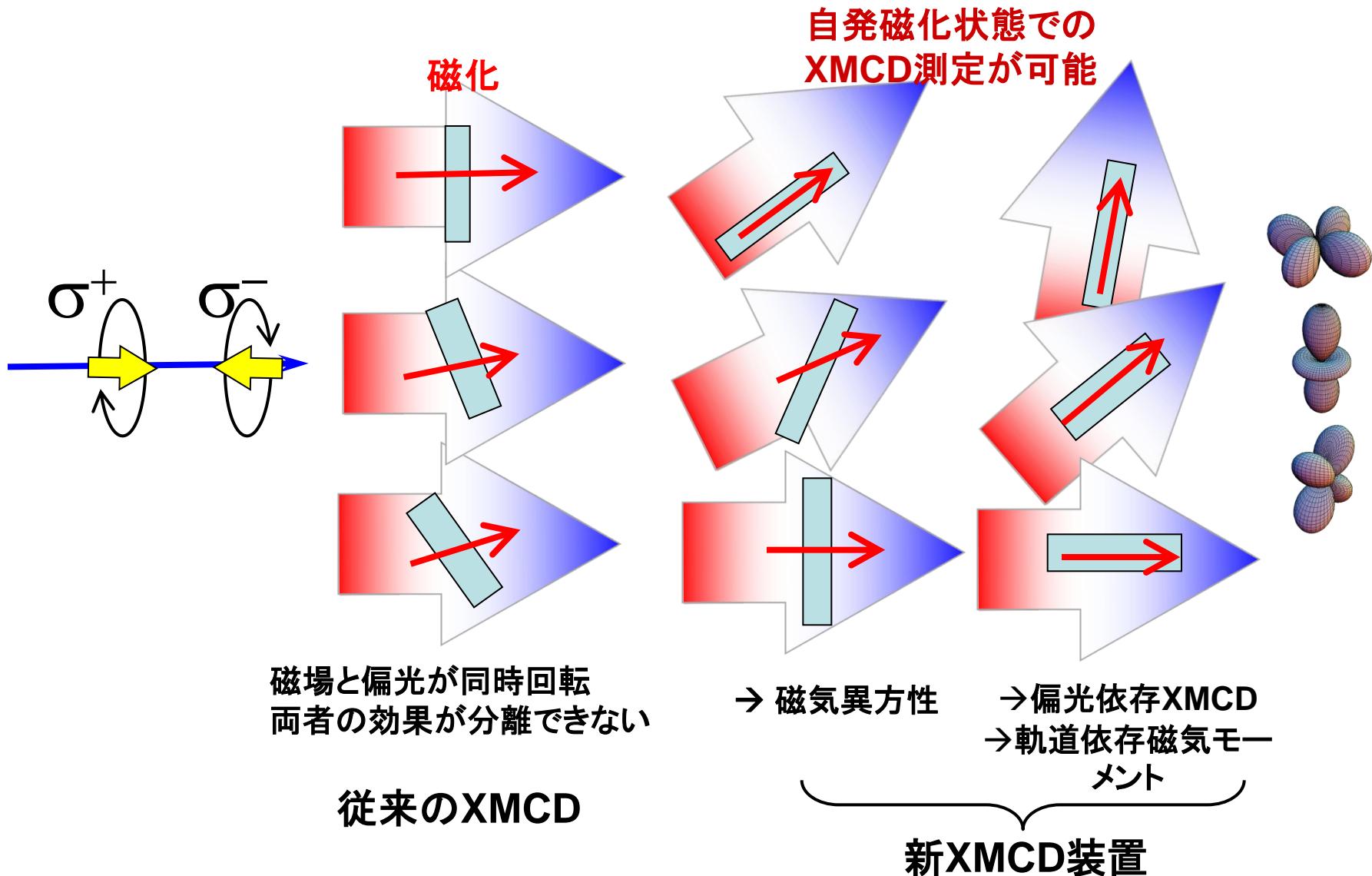
KEK-PF BL-16

ベクトル型超伝導
マグネット



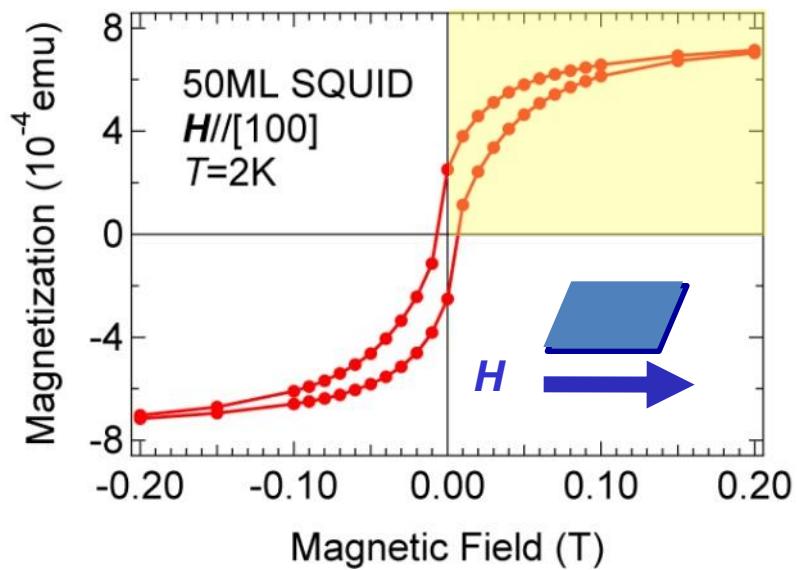
- 磁気異方性
- スピン-軌道状態
- 相競合, 相分離

方向可変磁場を用いたXMCD

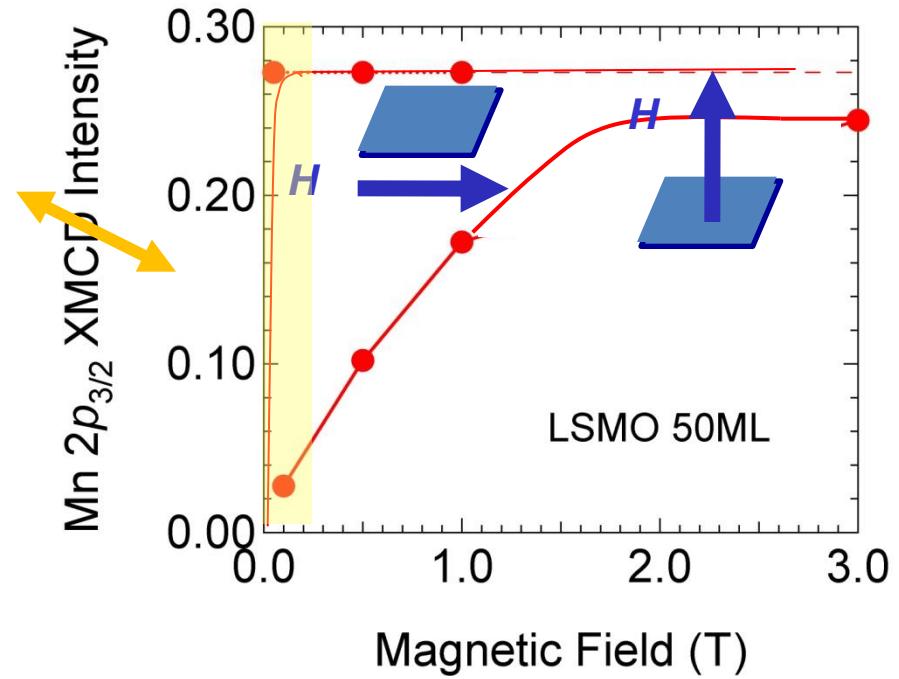


$\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ 薄膜の磁気異方性

Magnetization measured by SQUID

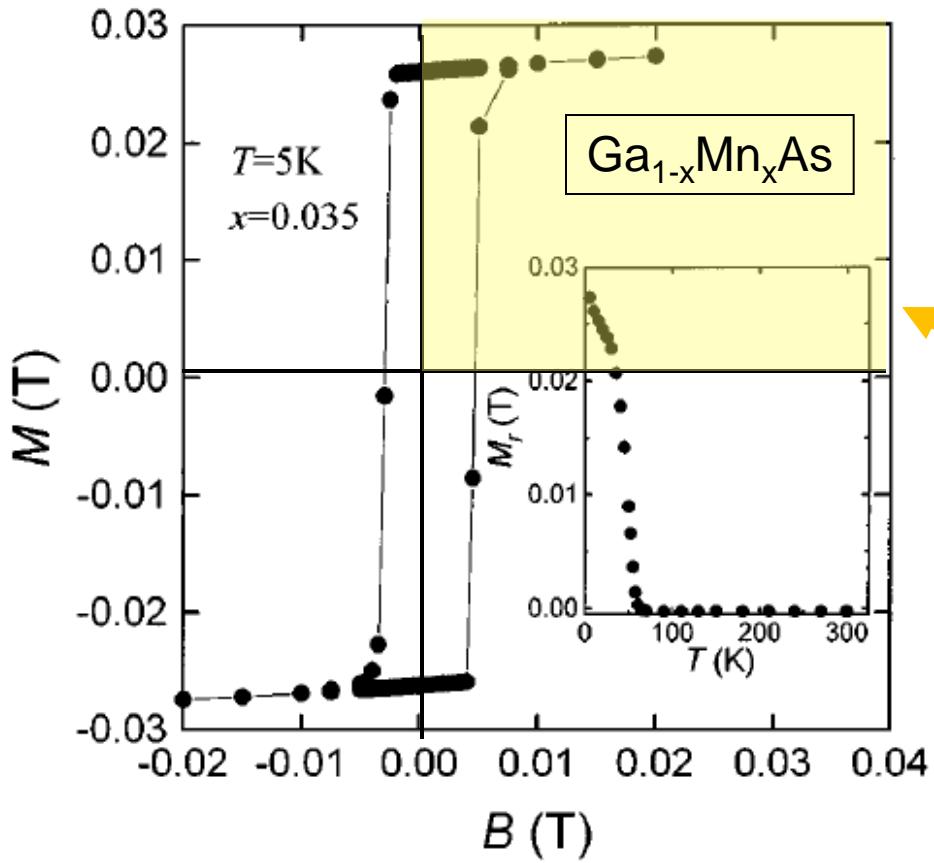


XMCD intensity vs H



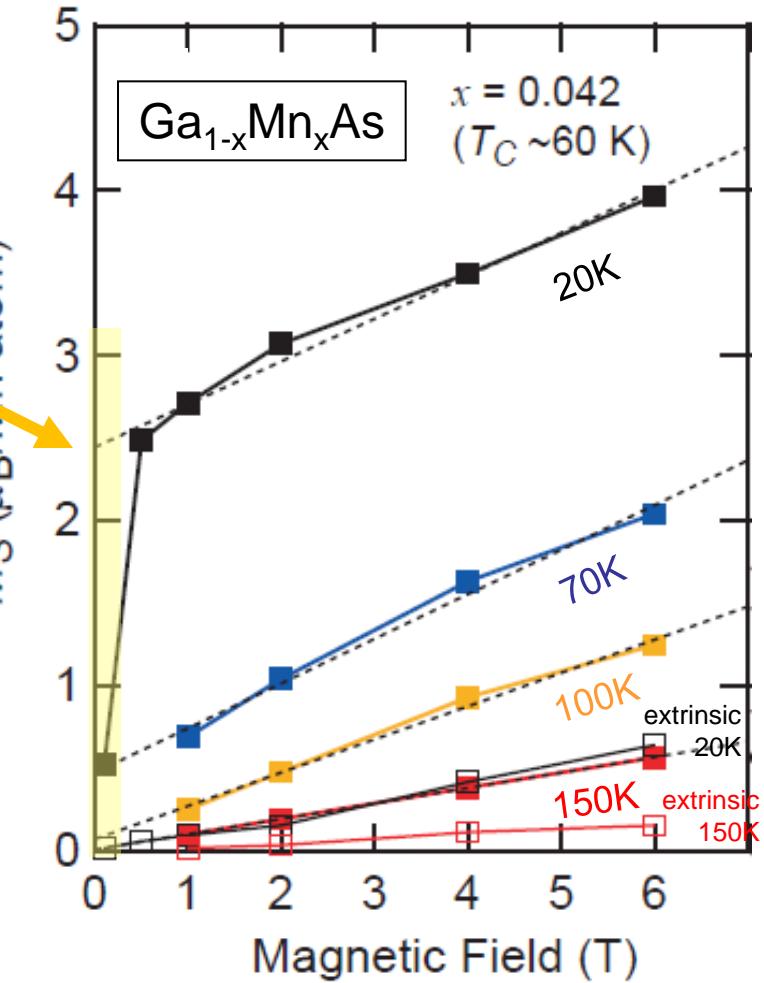
$Ga_{1-x}Mn_xAs$ (薄膜試料)の磁気異方性

Magnetization measured
by SQUID



H. Ohno e al. APL '96

XMCD intensity vs H

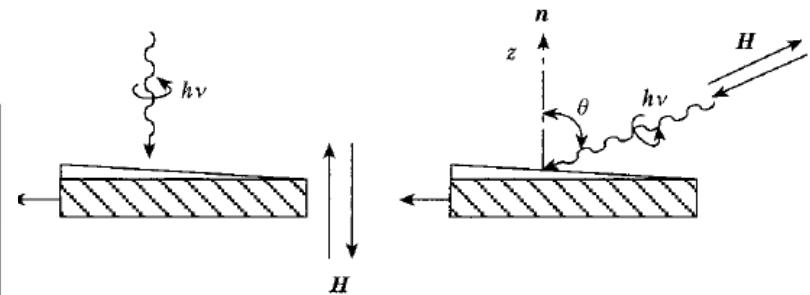


Y. Takeda et al. PRL '08

XMCDの角度依存性測定によるスピン・モーメント およびスピン分布四重極モーメントの決定

偏光X線を用いた d 軌道の異方性の検出

Orbital i	$\alpha = x, y \text{ or } z$ $4\langle Q_{\alpha}^{\frac{1}{2}} \rangle$	Linear Pol. $\vec{E} \parallel \alpha$ $4\left(1 - \frac{1}{4}\langle Q_{\alpha}^{\frac{1}{2}} \rangle\right)$	Circular Pol. $\vec{k} \parallel \alpha$ $4\left(1 + \frac{1}{8}\langle Q_{\alpha}^{\frac{1}{2}} \rangle\right)$	Circ. Dichroism $\vec{k}, \vec{H}_{\text{ext}} \parallel \alpha$ $4\left(1 + \frac{1}{2}\langle Q_{\alpha}^{\frac{1}{2}} \rangle\right)$
d_{yz}	 -8 16	 6 0	 3 6	 0 12
d_{zx}	 -8 -8	 6 6	 3 3	 0 12
$d_{3z^2-r^2}$	 -16 8	 8 2	 2 5	 -4 8
d_{xy}	 16 -8	 0 6	 5 3	 12 0
$d_{x^2-y^2}$	 16 -8	 0 6	 6 3	 12 0



$$m_{\text{orb}}^{\theta} = -\frac{4}{3} \left[\frac{\Delta I_{L_3} + \Delta I_{L_2}}{I_{L_3} + I_{L_2}} \right]_{\theta} n_h \mu_B$$

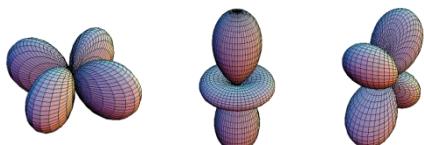
$$m_{\text{spin}} - 7 m_T^{\theta} = -2 \left[\frac{\Delta I_{L_3} - 2\Delta I_{L_2}}{I_{L_3} + I_{L_2}} \right]_{\theta} n_h \mu_B$$

$$m_T^{\perp} + 2m_T^{\theta} = 0$$

$$m_{\text{orb}}^{\theta} = m_{\text{orb}}^{\perp} \cos^2 \theta + m_{\text{orb}}^{\parallel} \sin^2 \theta$$

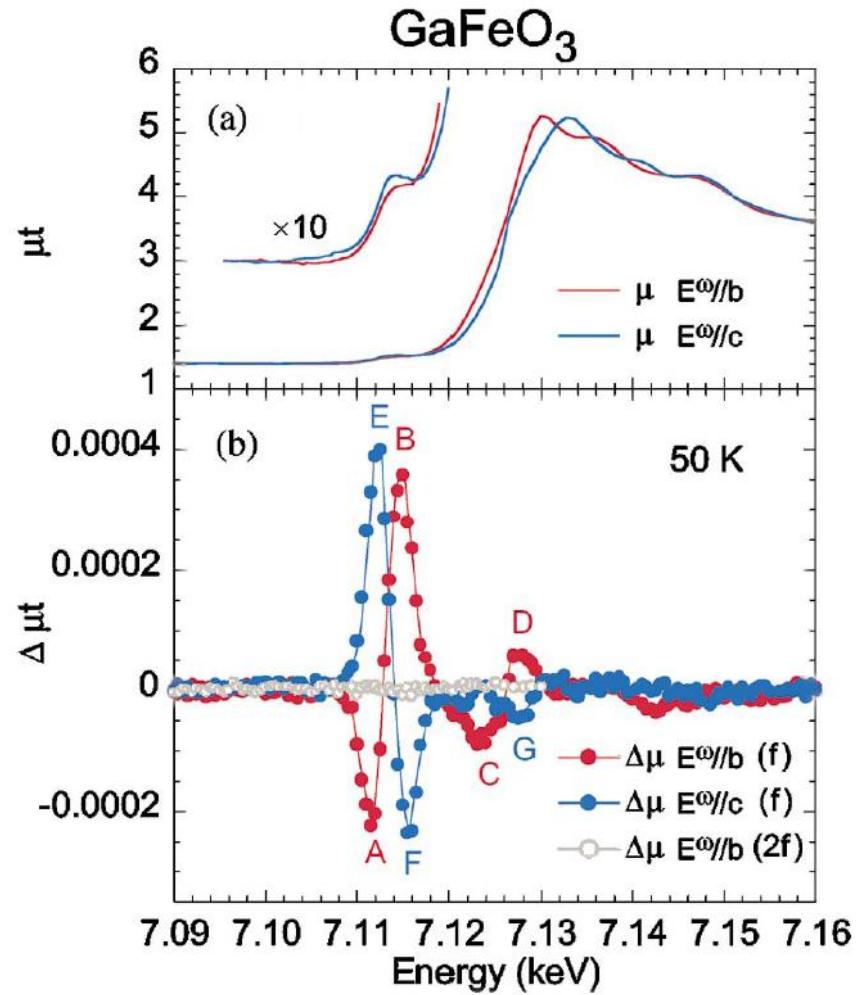
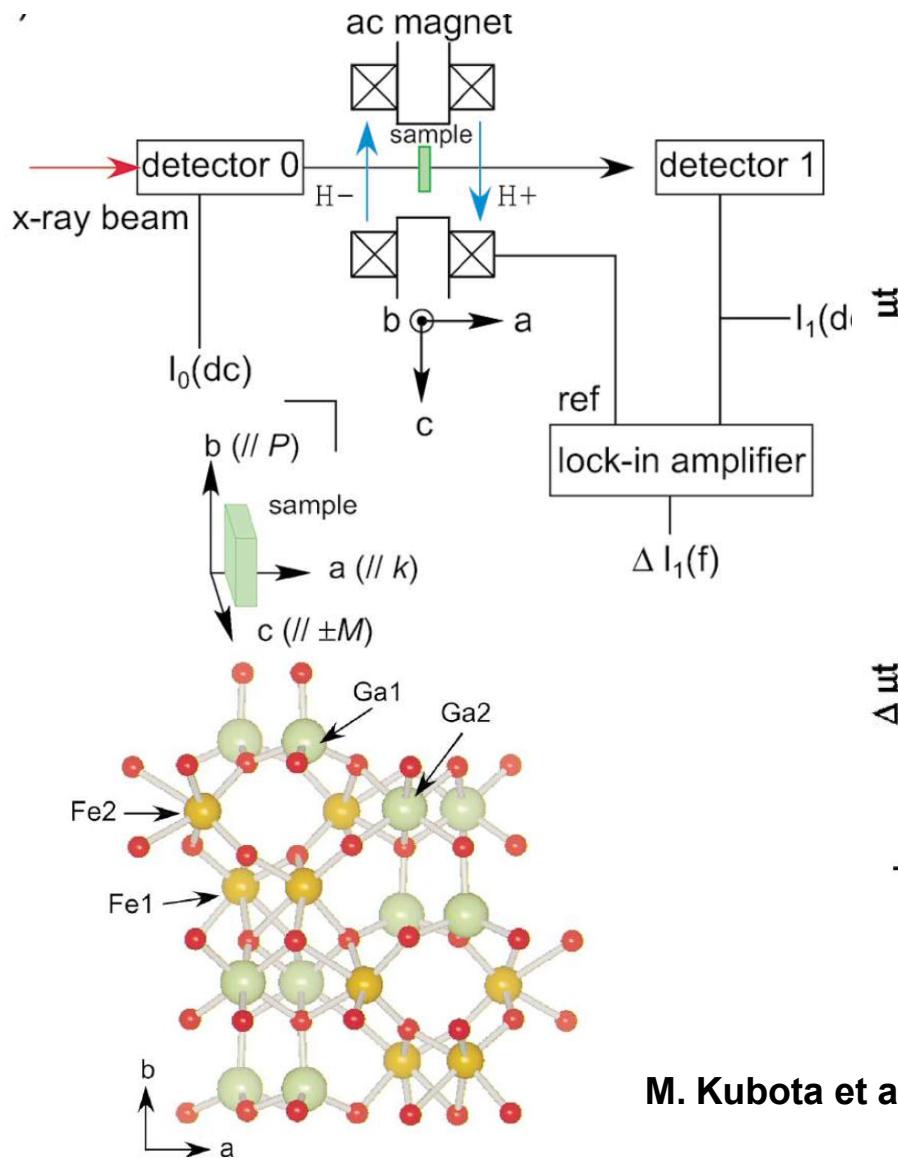
$$m_T^{\theta} = m_T^{\perp} \cos^2 \theta + m_T^{\parallel} \sin^2 \theta$$

$\rightarrow m_{\text{spin}}, m_{\text{orb}}^{\theta}, m_T^{\theta}$ の完全決定 ($\theta = \perp, \parallel$)



小出常晴:「新しい放射光の科学」
菅野, 藤森, 吉田編(講談社2000年)

軟X線非相反方向二色性(XNDD)の測定



M. Kubota et al., PRL '04

まとめ

- 研究対象
 - 希薄磁性半導体
 - 磁性薄膜：厚さ依存，界面，表面，基板圧力，電場，．．．
 - 磁性ナノ構造：細線，ナノワイヤー，ナノ粒子．．．
 - マルチフェロイック系，スピンドホール系
- 今後の発展
 - 磁場方向依存 → 磁気異方性(希土類磁石？)
 - 偏光依存
 - 強磁場
 - 時間分解
- 他の手法との組み合わせ
 - 共鳴軟X散乱
 - ．．．