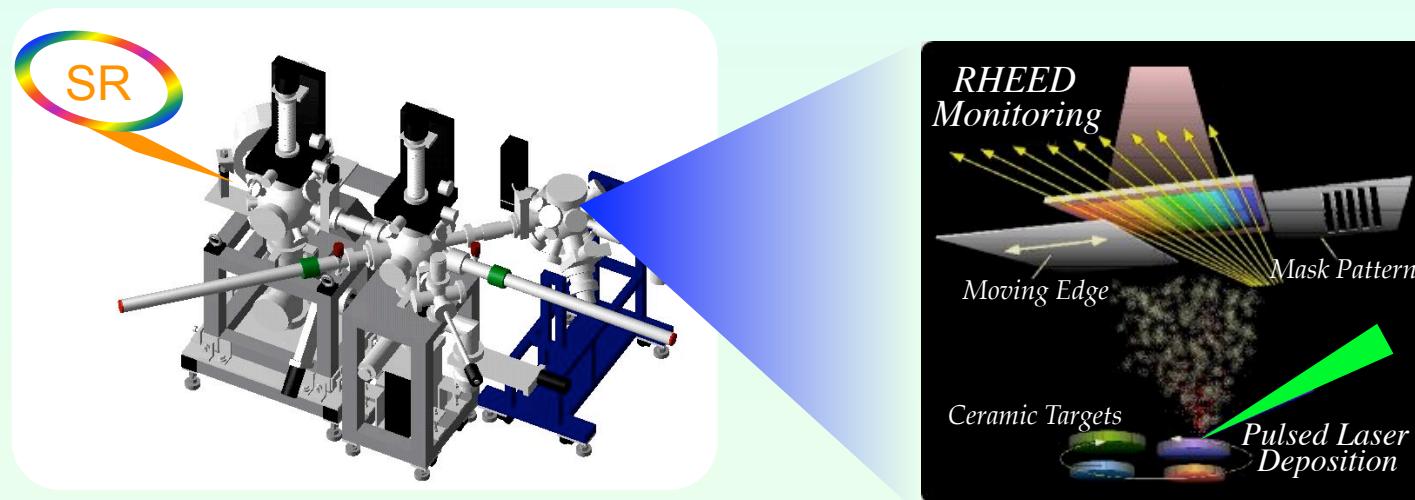


# 次元性制御強相関酸化物薄膜の 軟X線光電子分光

組頭 広志

KEK-PF (物構研&構造物性研究センター)  
JSTさきがけ



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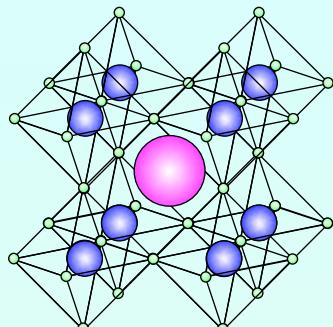
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*Oak Ridge National Laboratory, Oak Ridge Tennessee 37831-6071, USA*

# INTRODUCTION

*Transition Metal Oxides*



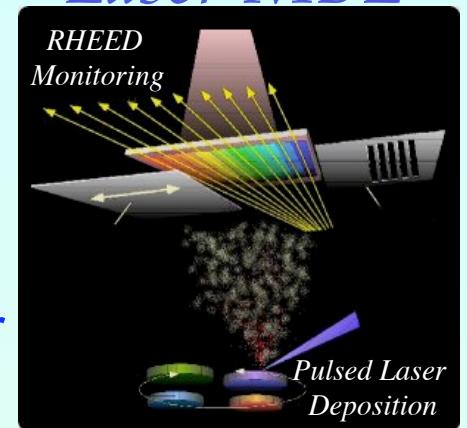
High-Tc  
CMR  
M-I Transition

Controlling unusual  
physical properties

**Charge**

*Charge Transfer*

*Laser MBE*



**Heterostructure**

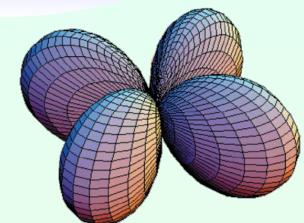
**Spin**

*Spin Exchange*



**Orbital**

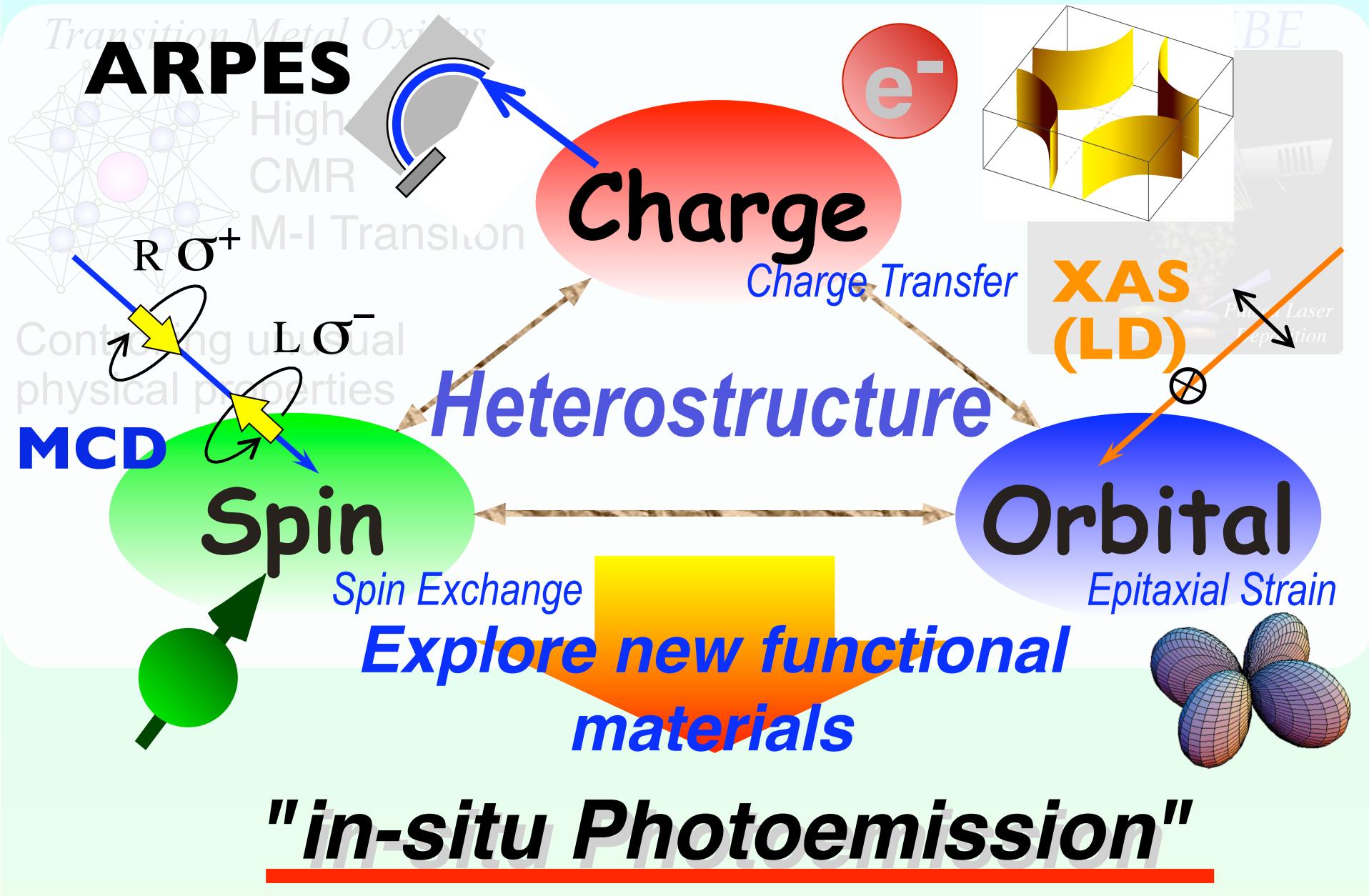
*Epitaxial Strain*



***Explore new functional  
materials***

***"in-situ Photoemission"***

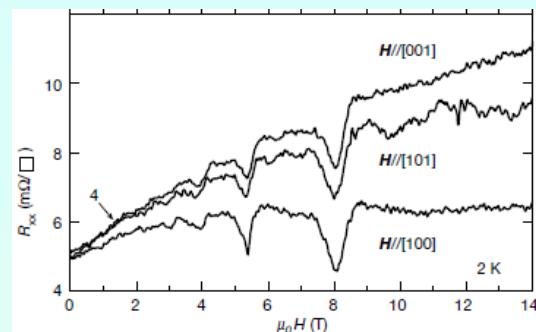
# Oxide Artificial Structures



# Appearance of metallic conductivity at the interface between the band insulators $\text{LaAlO}_3$ and $\text{SrTiO}_3$

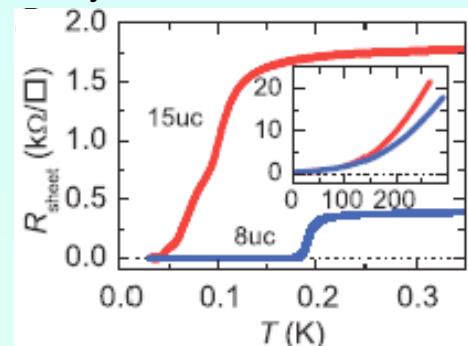
## High-mobility electron gas

(A. Ohtomo *et al*, Nature 2004)



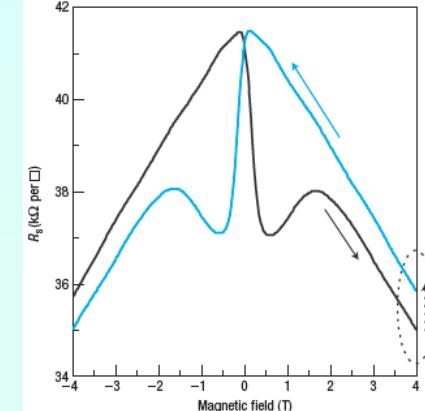
## Superconductivity

(N. Reyren *et al.*, Science 2007)



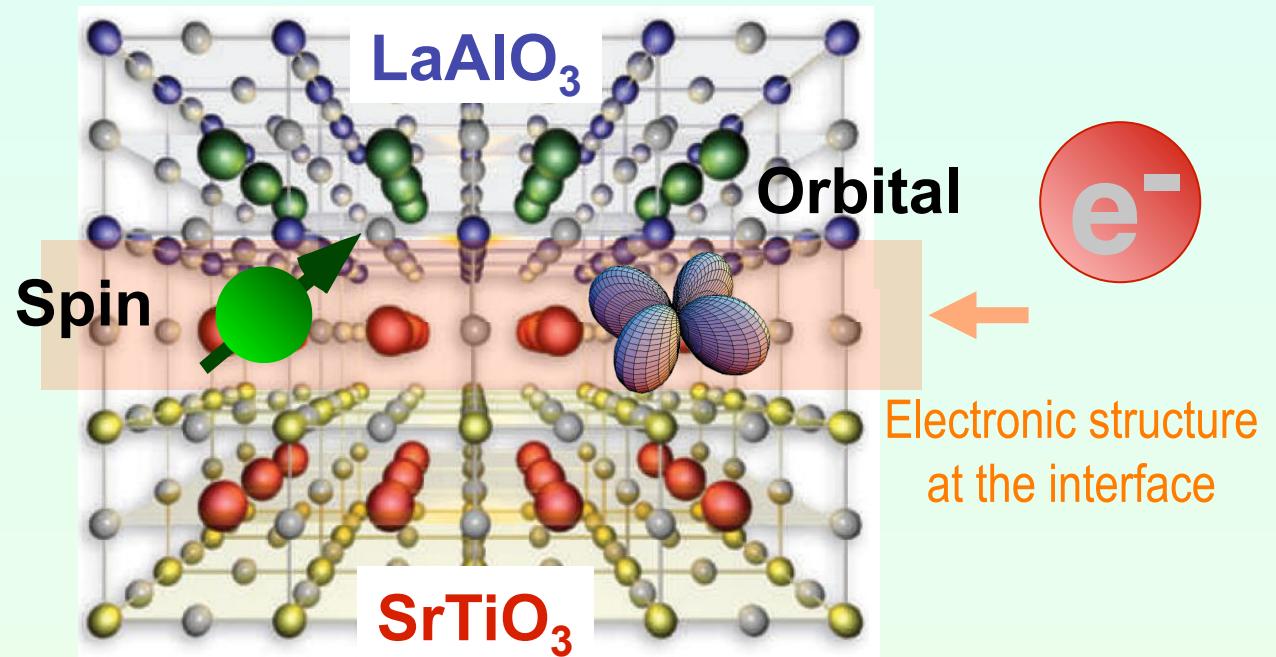
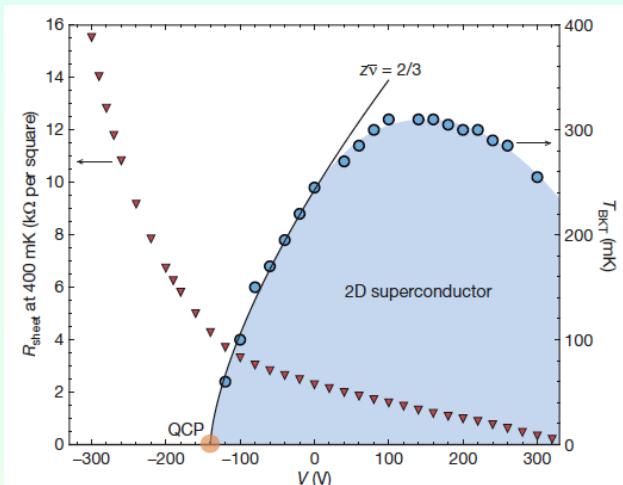
## Ferromagnetism?

(A. Brinkman *et al*, Nat. Mater. 2007)

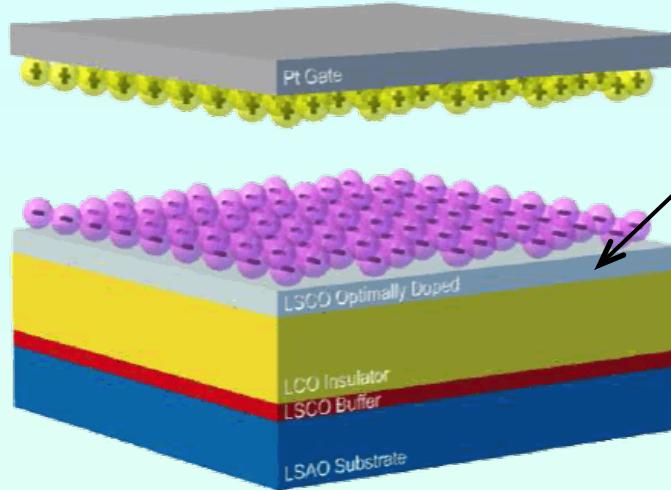


## Electric field control of the superconductivity

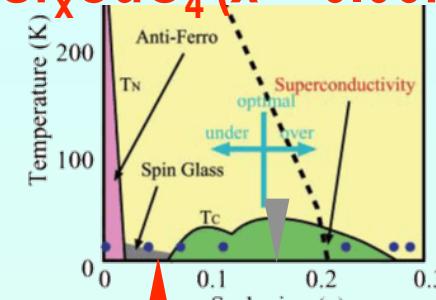
(A.D. Caviglia *et al*, Nature 2008)



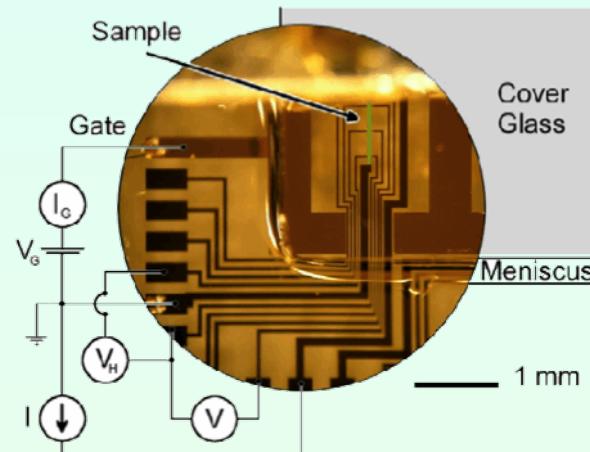
# Superconductor–Insulator Transition in $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ Driven by Electric Field



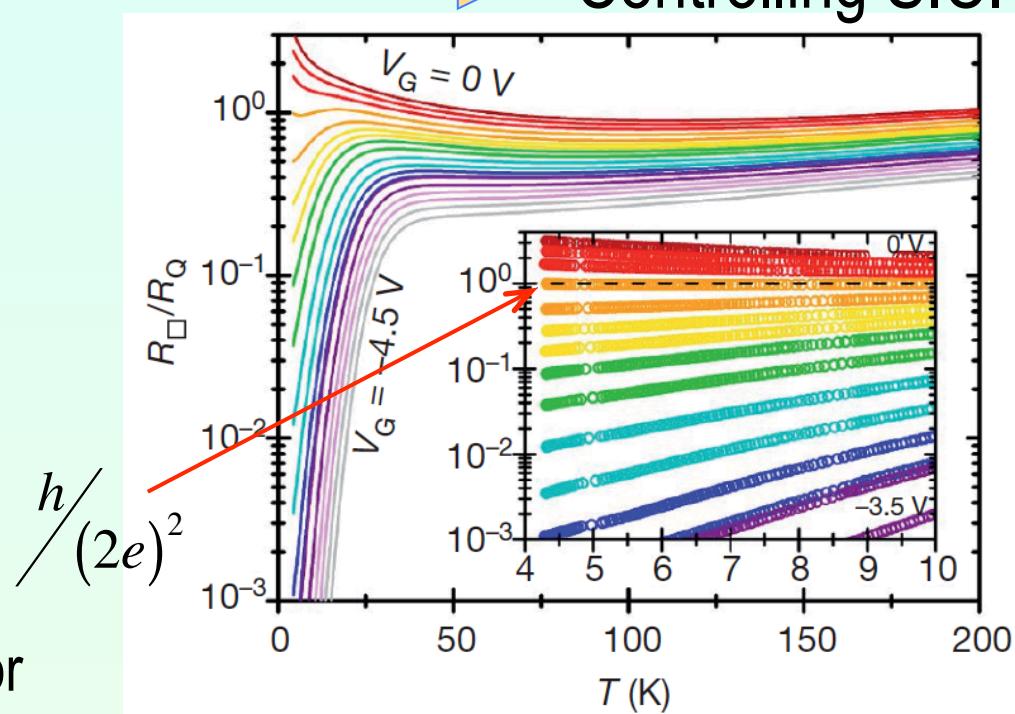
1 u.c.  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  ( $x = 0.06$ : Under doped)



Controlling S.C.



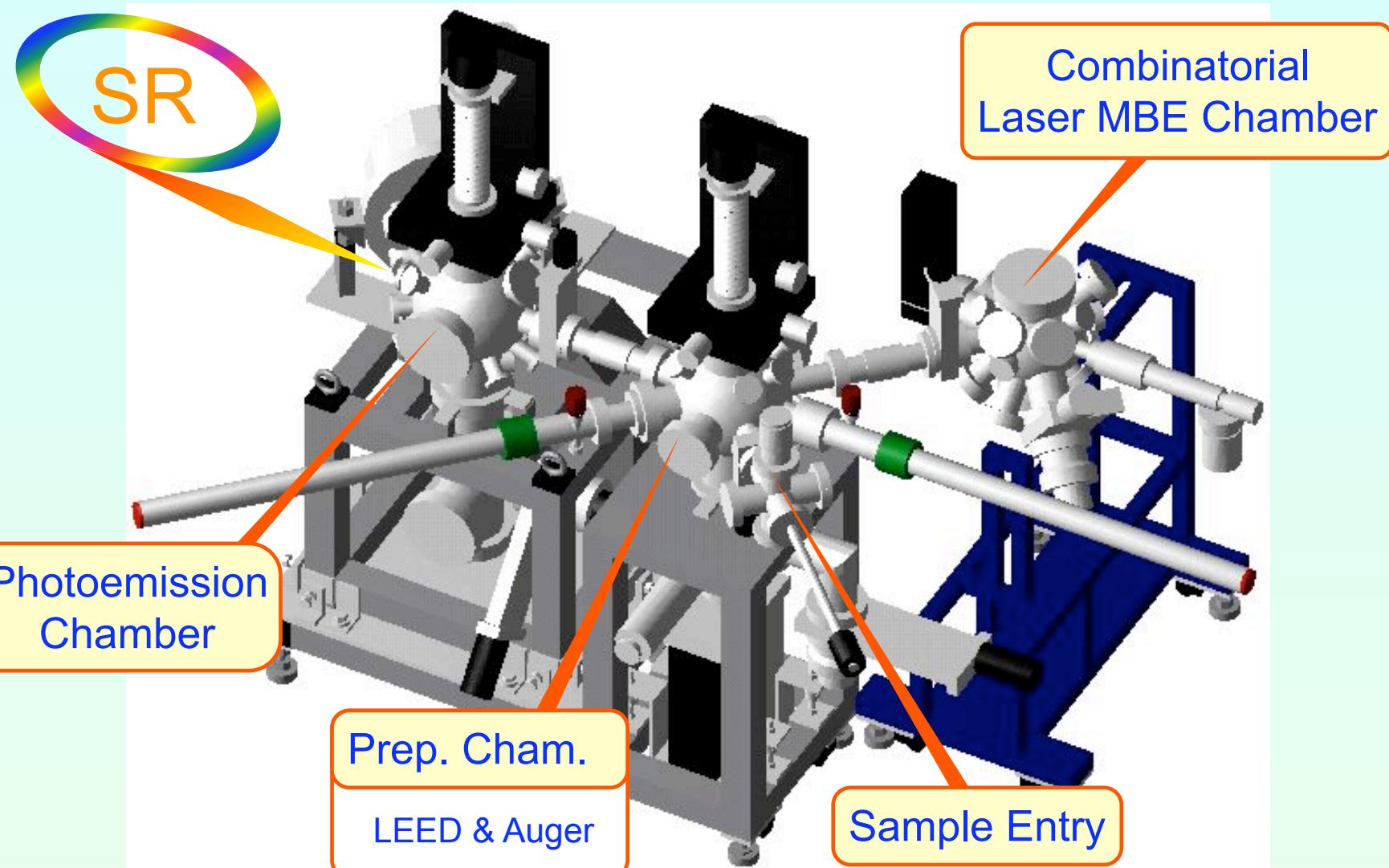
Electrolyte Double-Layer Transistor  
(Field Effect Transistor)



A.T. Bollinger et al., Nature 472, 458 (2011).

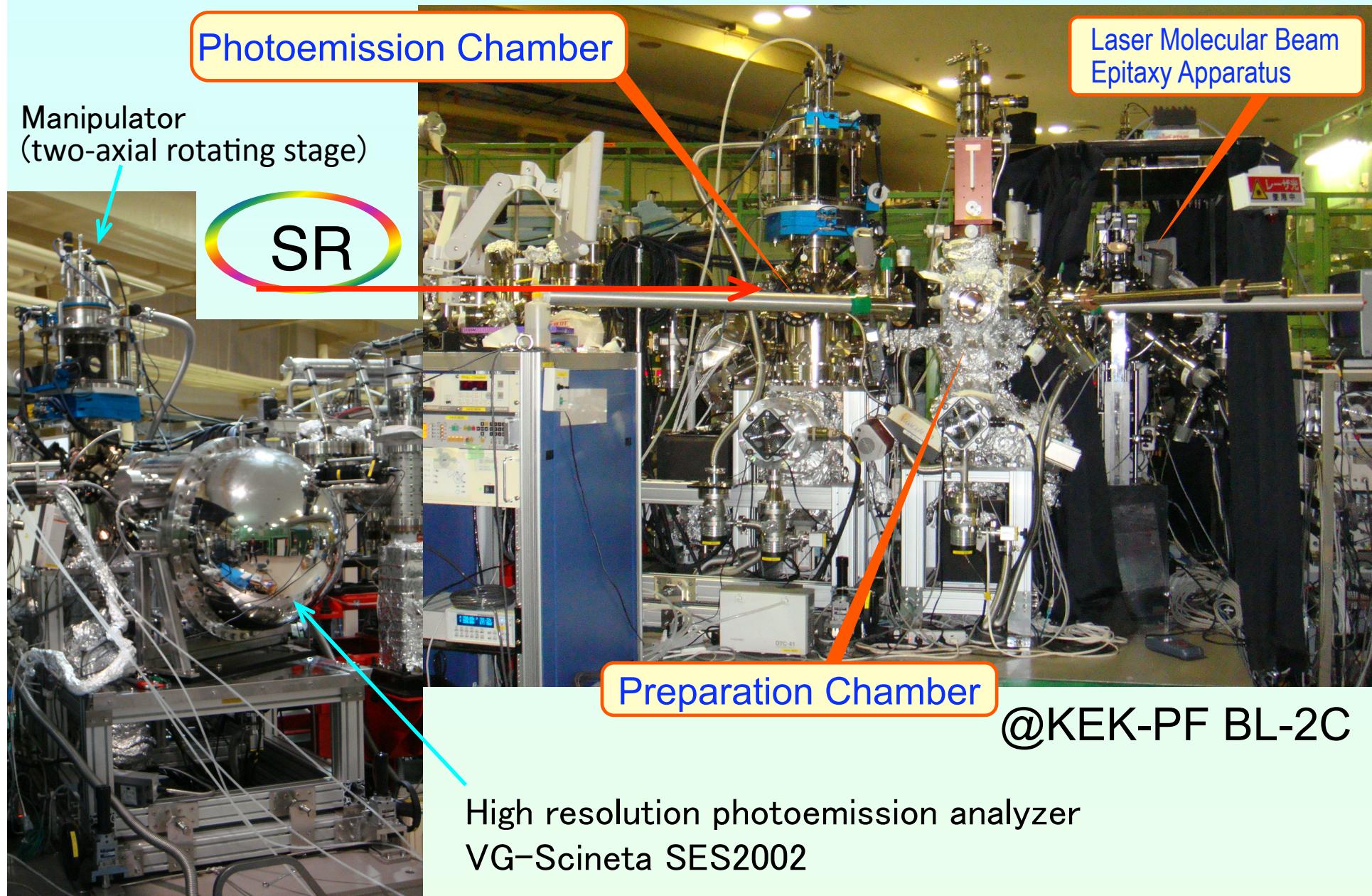


# In-situ PES + Laser MBE system

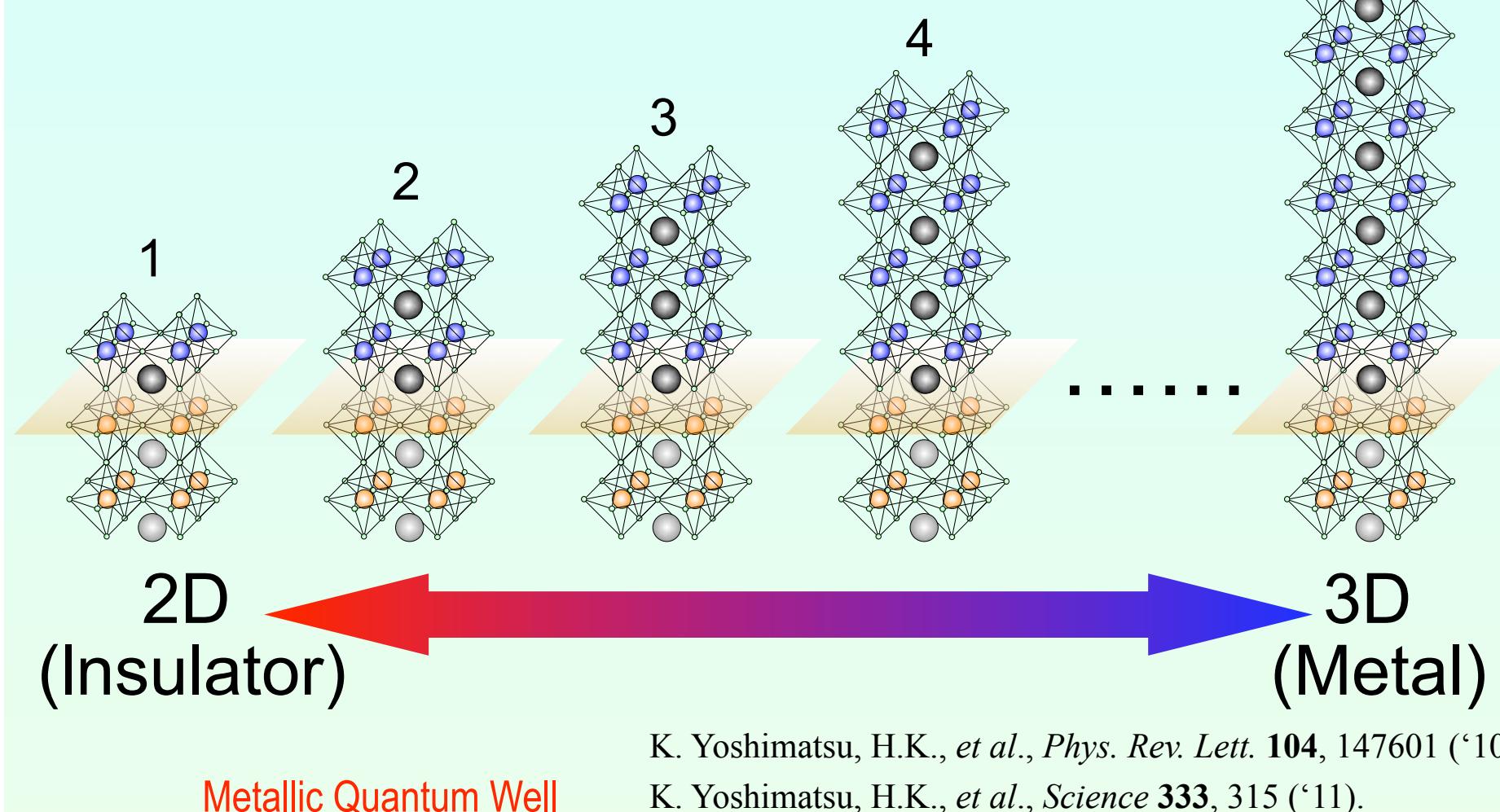




# New *in-situ* PES + Laser MBE system



# Dimensional-Crossover-Driven Metal-Insulator Transition in $\text{SrVO}_3$ ultrathin Films



K. Yoshimatsu, H.K., et al., *Phys. Rev. Lett.* **104**, 147601 ('10).

K. Yoshimatsu, H.K., et al., *Science* **333**, 315 ('11).

# Metal Insulator Transition

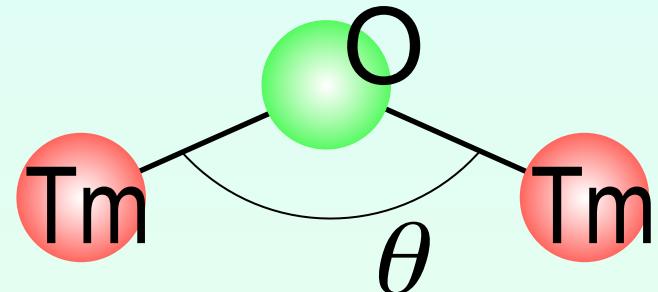
Mott-Hubbard theory

$U \gg W$  Insulator

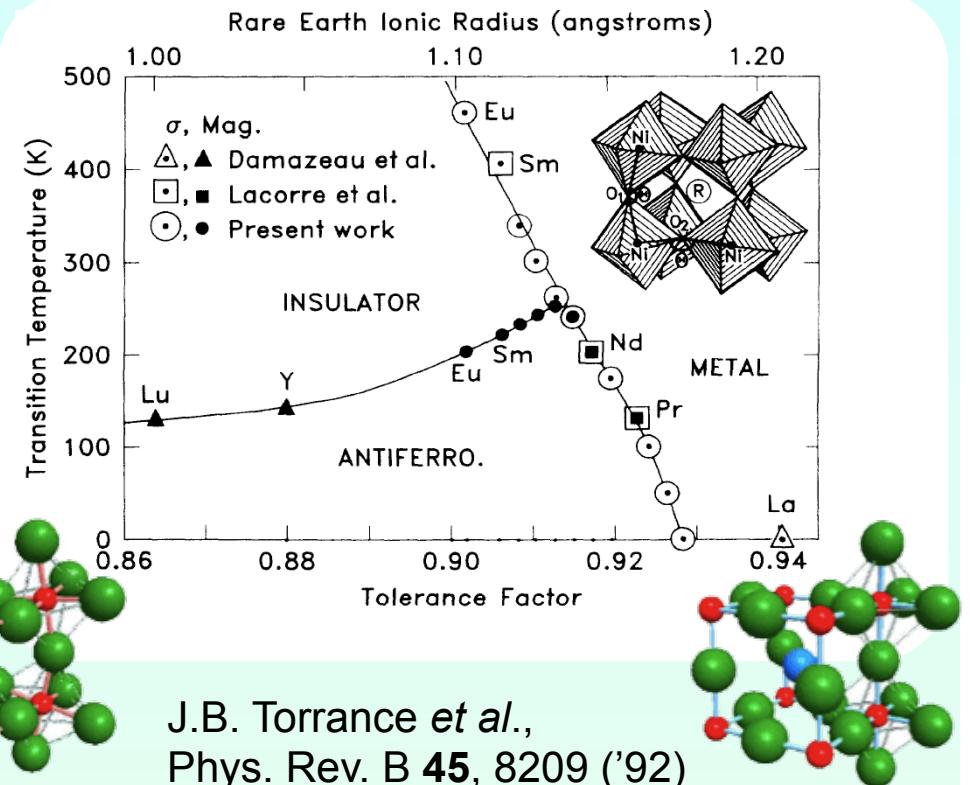
$U \ll W$  Metal

$U$ : on-site Coulomb repulsion

$W$ : one-electron band width



$$W \propto \cos^2 \theta$$



J.B. Torrance et al.,  
Phys. Rev. B **45**, 8209 ('92)

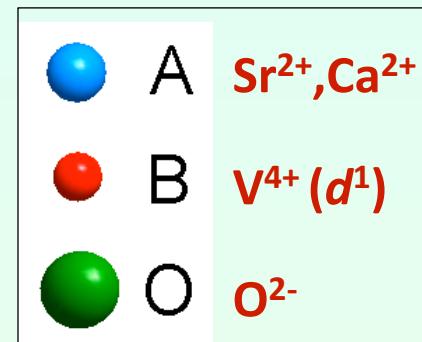
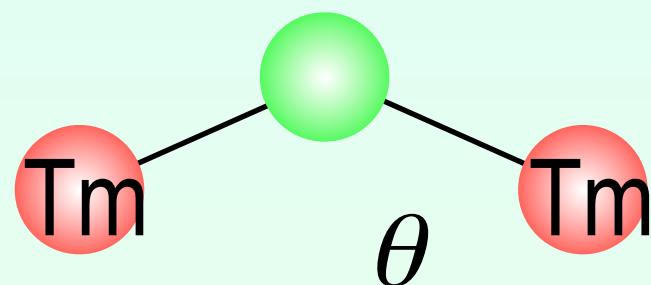
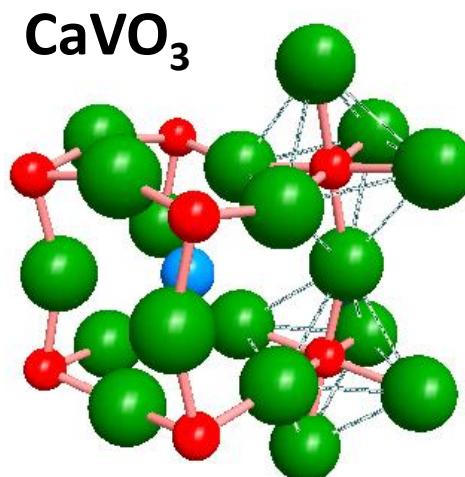
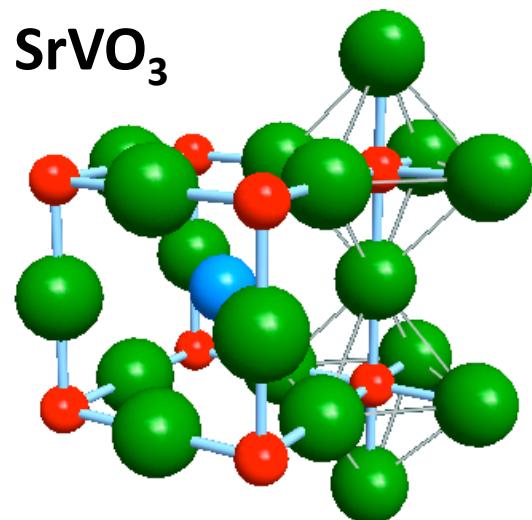
For bulk materials, MIT has been intensively studied by chemical substitution of constituent ions with ones having a smaller ion radius. However, such chemical substitution always induce randomness in a solid.

# Bandwidth control in $\text{Sr}_{1-x}\text{Ca}_x\text{VO}_3$

$$W = 12t$$

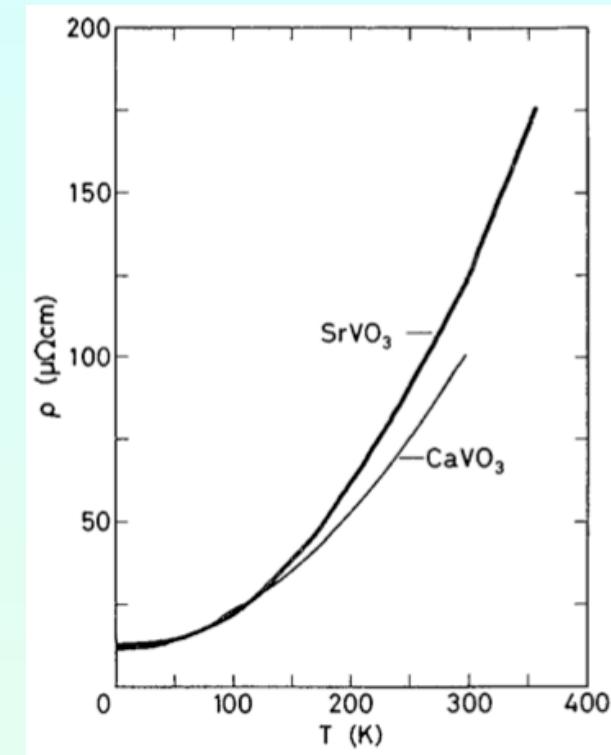
$$W = 10.6t$$

$$W \propto \cos^2 \theta$$



$$\theta : \begin{array}{l} \text{SrVO}_3 \ 180^\circ \\ \text{CaVO}_3 \sim 160^\circ \end{array}$$

$\text{Sr}_{1-x}\text{Ca}_x\text{VO}_3$  is always metal in the entire  $x$  range.

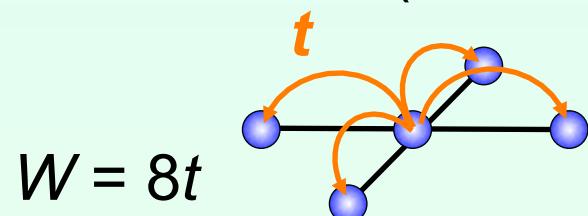
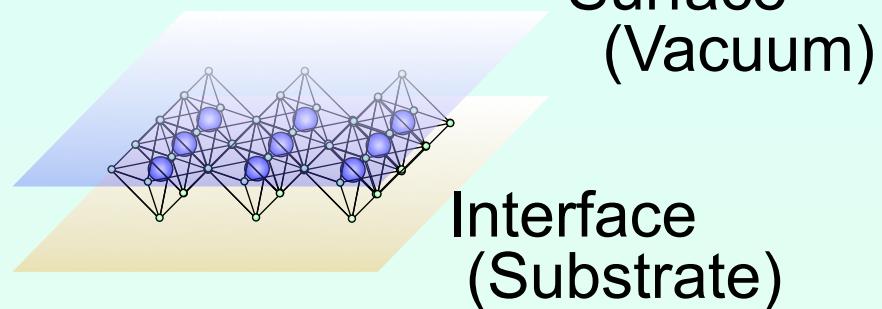


M. Onoda *et al.*, Solid State Commun. **79**, 281 (1991).

# Our Approach: Dimensional crossover occurring in an $\text{A}_{n+1}\text{B}_n\text{O}_{3n+1}$ Ruddlesden-Popper series

2D

1 ML

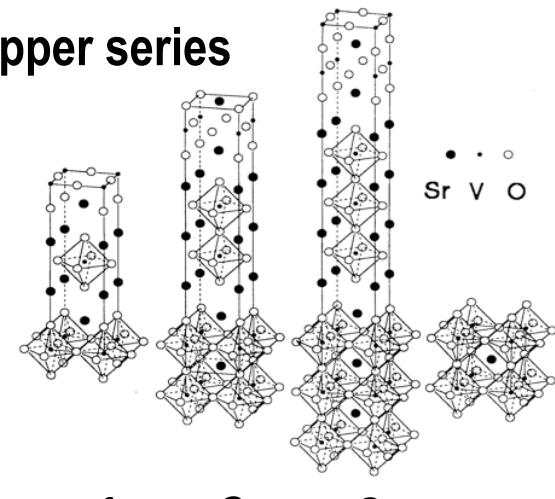


Insulator

Metal

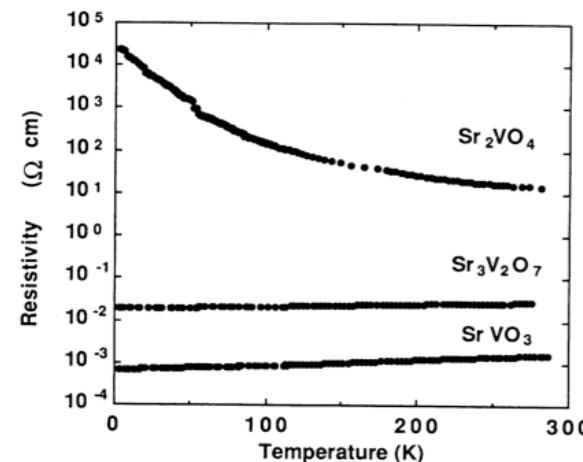
Dimensional crossover from 3D to 2D

$\text{A}_{n+1}\text{B}_n\text{O}_{3n+1}$   
Ruddlesden-Popper series



$n = 1$   
Insulator

$n=1$      $n=2$      $n=3$      $n=\infty$



A. Nozaki *et al.*, Phys. Rev. B **43**, 181 (1991).

$n > 2$   
Metal

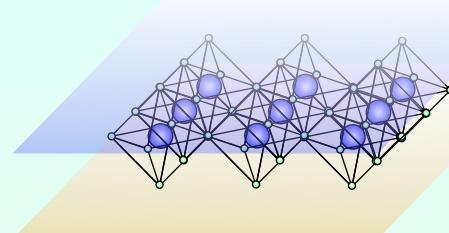
# Our Approach: Dimensional crossover occurring in an artificial structure

## Digitally-controlled $\text{SrVO}_3$ ultrathin films

**2D**

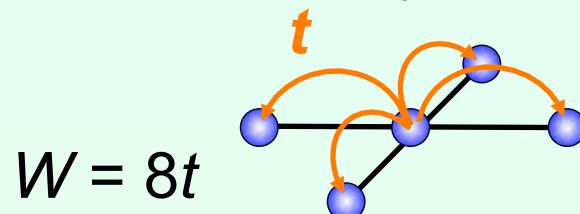
$\text{SrVO}_3$ : Paramagnetic metal  
 $3d^1$  system

**1 ML**



Surface  
(Vacuum)

Interface  
(Substrate)

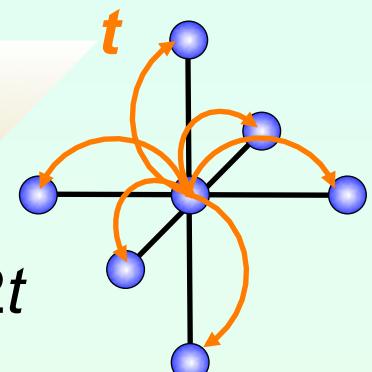
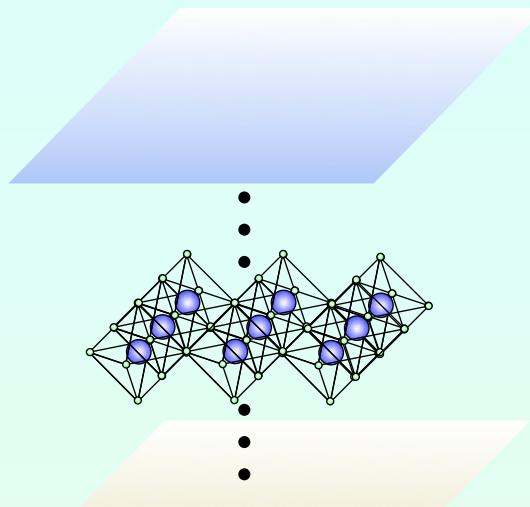


$$W = 10.6t$$

( $\text{CaVO}_3$  : Metal)

**3D**

$\infty$  MLs



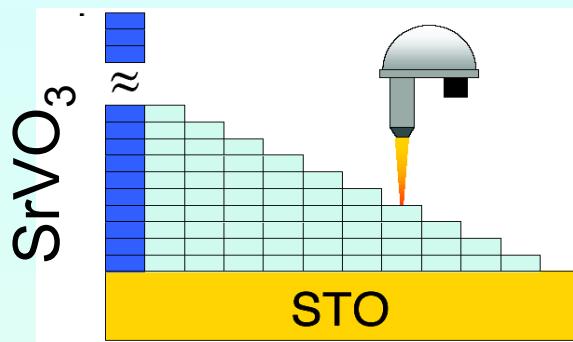
Insulator

Metal

Dimensional crossover from 3D to 2D

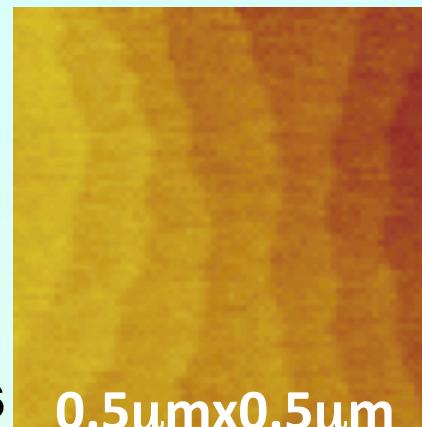
# Atomically-flat Surface and Abrupt Interface

Preconditions to thickness dependent experiments

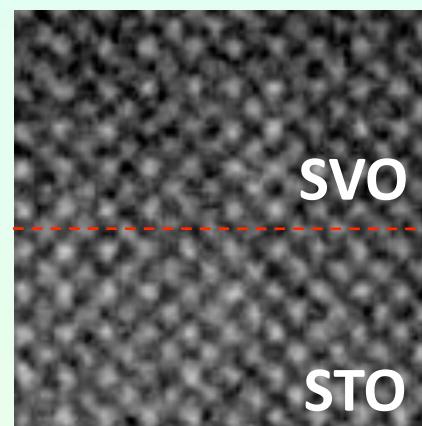


- ✓ Precise digital control of thickness (RHEED oscillation)
- ✓ Atomically flat surface (AFM image)
- ✓ Chemically abrupt interface (TEM image)
- ✓ Coherent growth of thin film (Reciprocal space mapping)

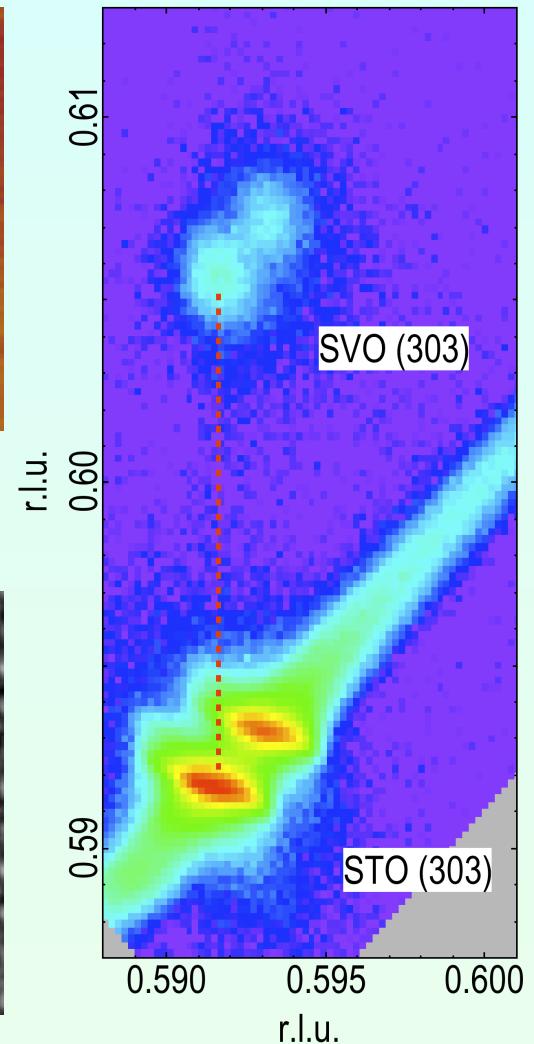
AFM image



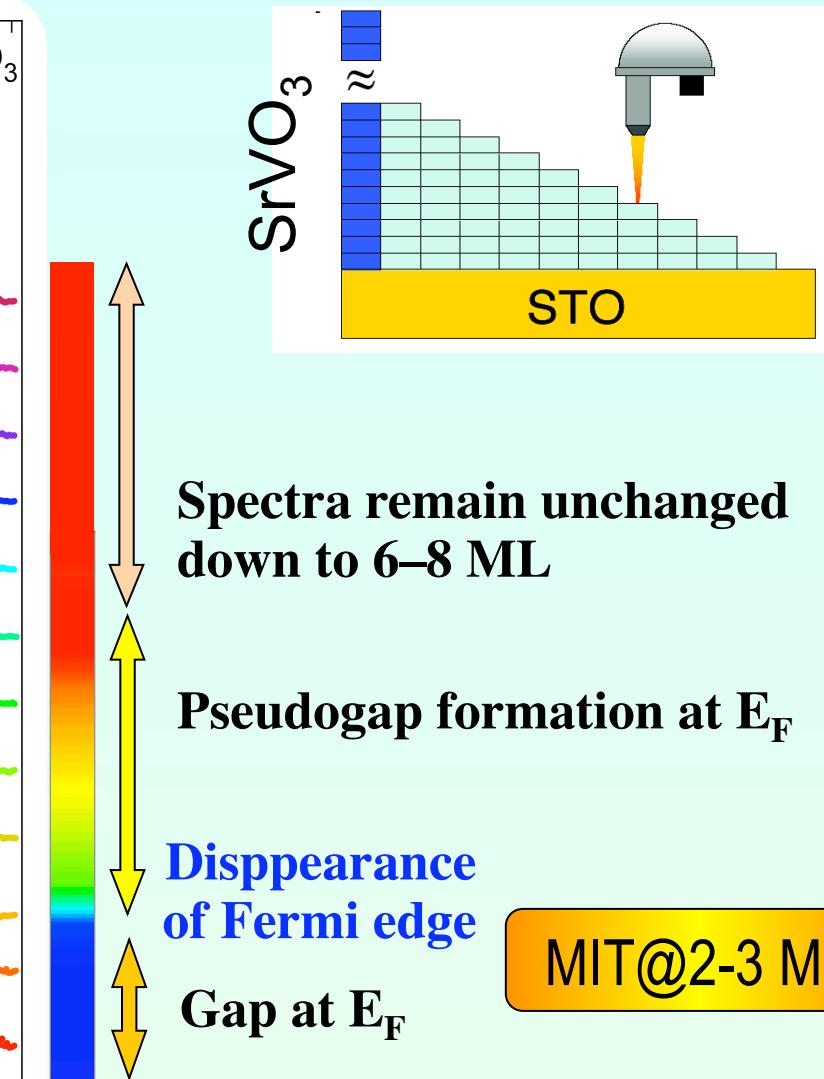
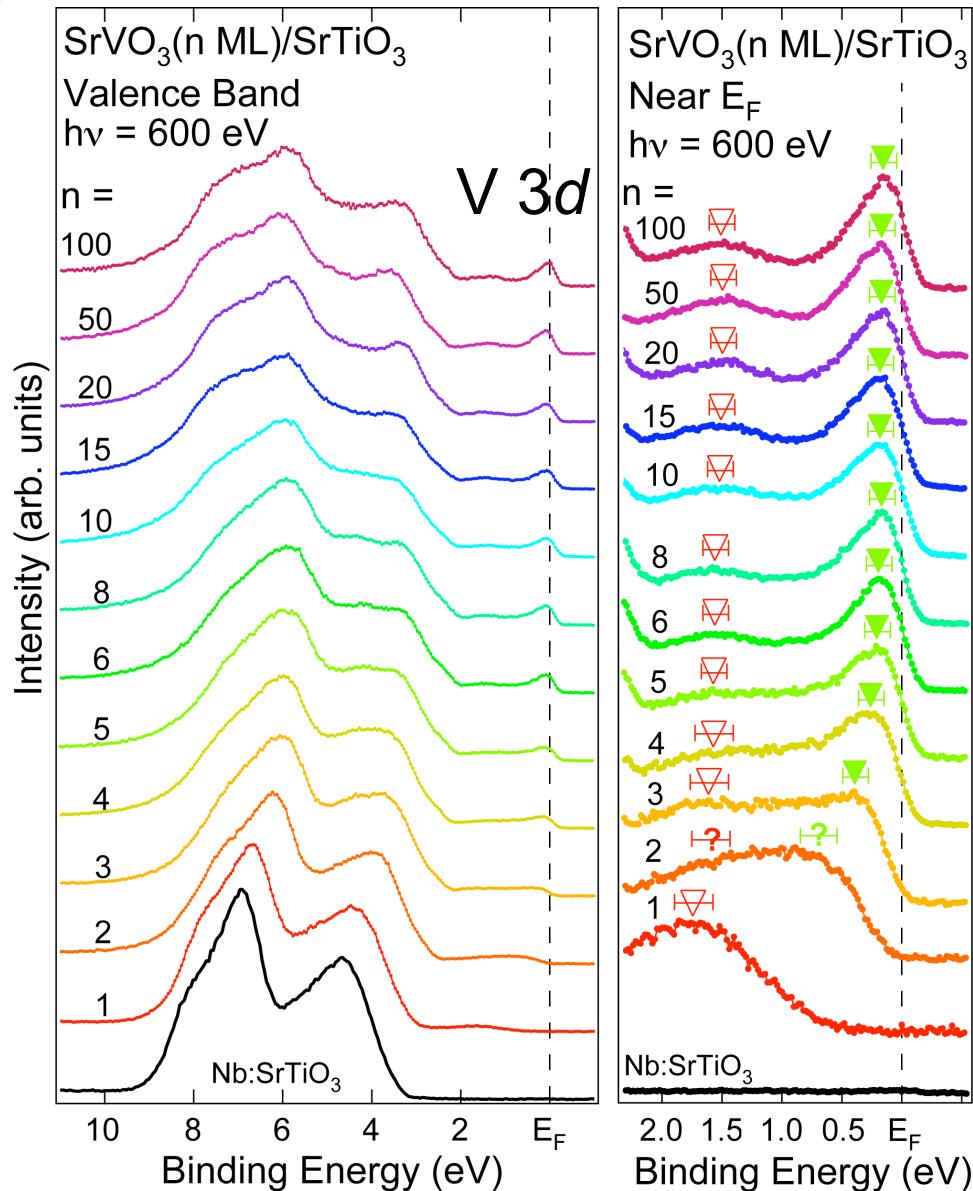
TEM image



4cXRD RSM around (303)

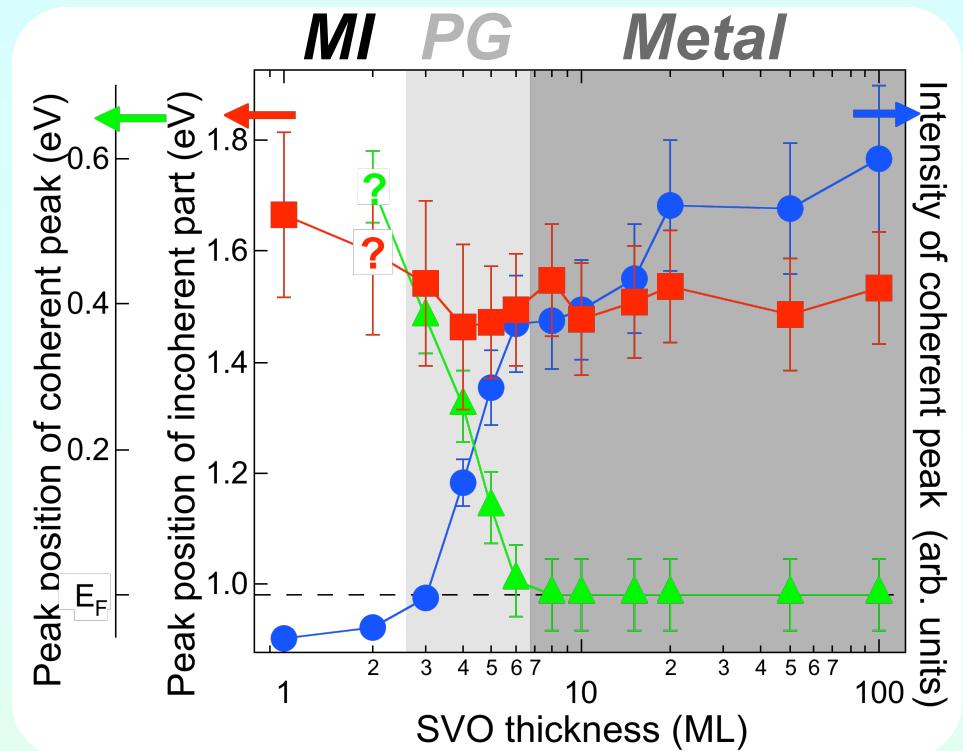
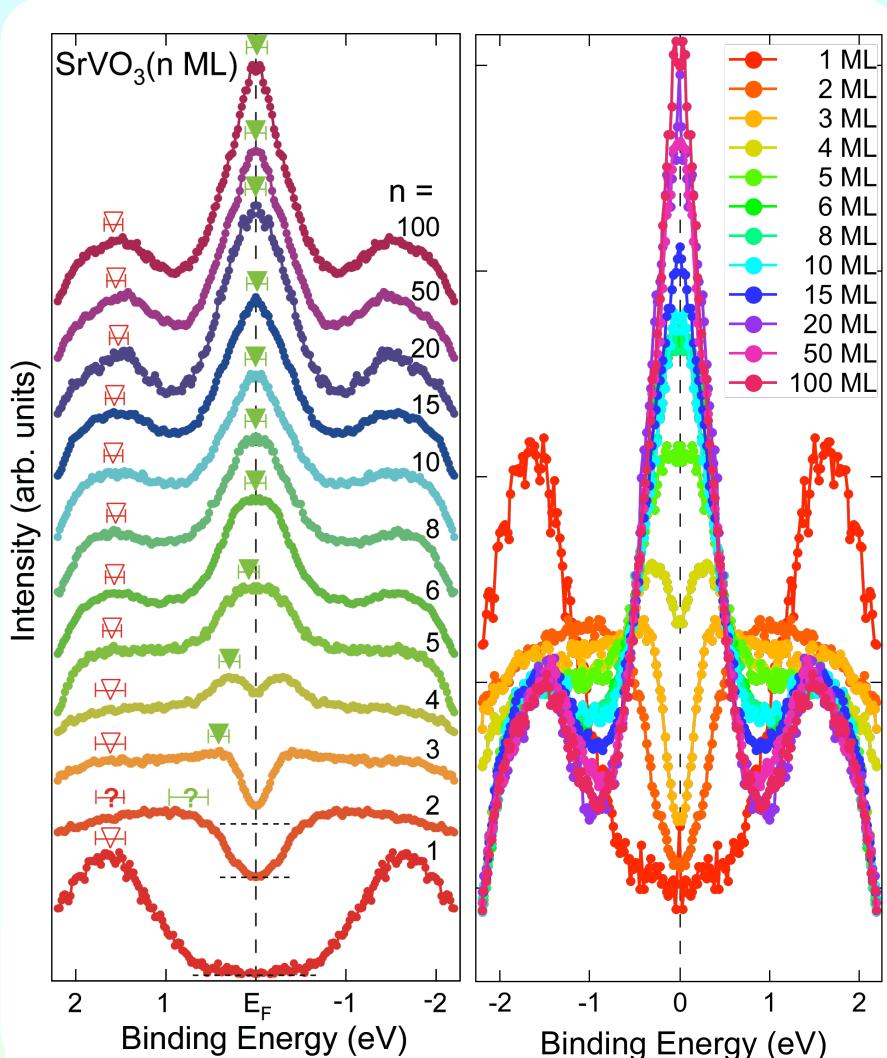


# *In situ* PES Spectra of $\text{SrVO}_3$ thin Films



Metal-insulator transition occurs in the SVO thickness of 2–3 ML.

# Dimensional-Crossover-Driven MIT in $\text{SrVO}_3$ Films

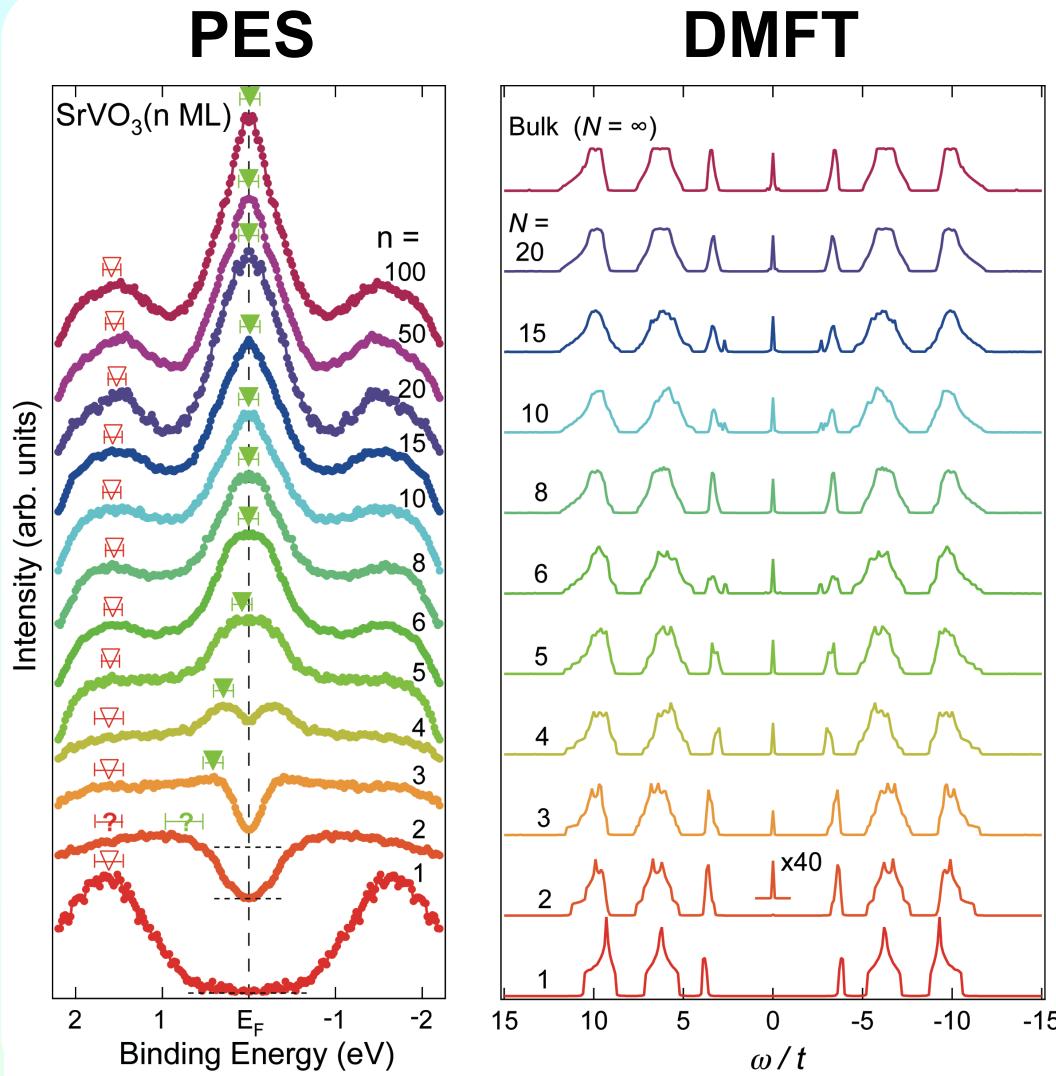


Quasiparticle peak intensity steeply increasing with increasing film thickness ( $W$  increasing).

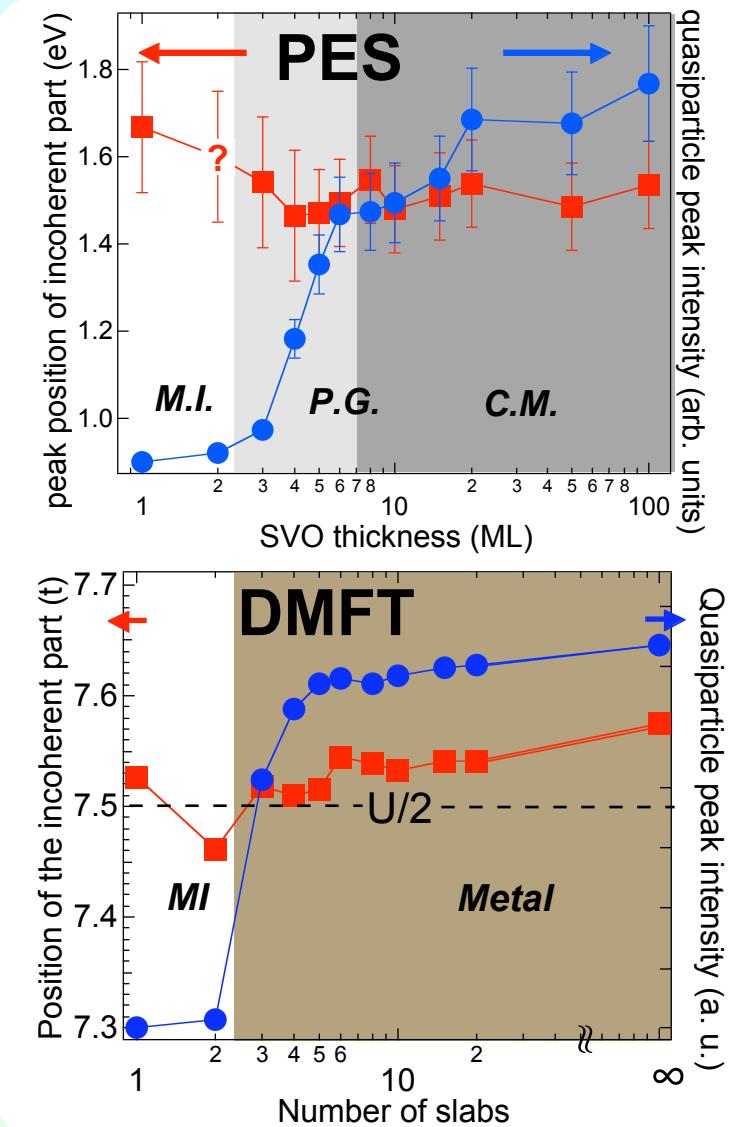
Peak position of the incoherent peak remains unchanged ( $U$  is constant).

Dimensional-Crossover-Driven (from 3D to 2D) MIT in  $\text{SrVO}_3$  Ultrathin Films

# Comparison between PES and Layer DMFT Cal.



K. Yoshimatsu, H.K., et al., Phys. Rev. Lett. **104**, 147601 ('10).



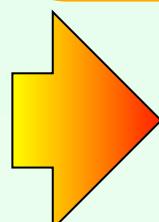
**Spectral behavior is well reproduced by layer DMFT calculation.**

# Summary: $\text{SrVO}_3$ ultrathin films

We have investigated the change occurring in the electronic structure of digitally controlled  $\text{SrVO}_3$  ultrathin films across the dimensional-crossover-driven MIT using *in-situ* photoemission spectroscopy.

## We have found

1. The V 3d derived states near  $E_F$  remain unchanged down to 6–8 ML.
2. In the film thickness of 3–6 ML, a pseudogap is formed at  $E_F$  owing to the spectral weight transfer from the coherent part to incoherent part.
3. The pseudogap finally evolves into an energy gap below 2 ML, indicating the occurrence of MIT at a critical thickness of 2–3 ML.
4. These spectral behaviors are well reproduced by layer DMFT calculations.

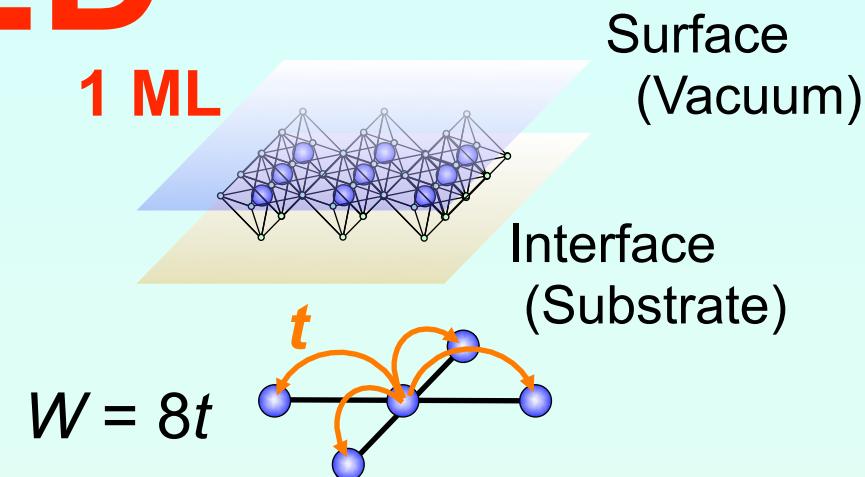


The observed MIT is caused by the reduction of bandwidth due to dimensional-crossover from 3D to 2D.

# Summary

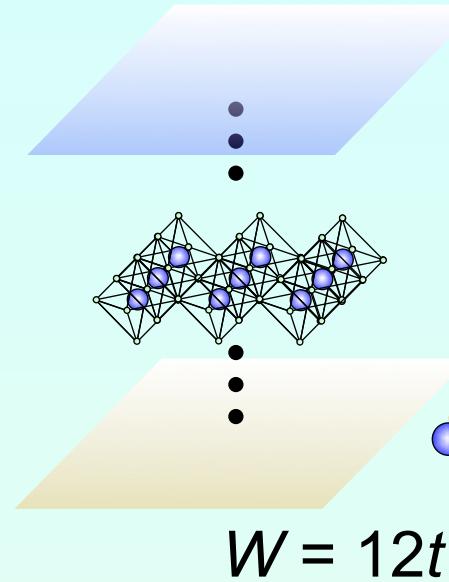
**2D**

**1 ML**



**3D**

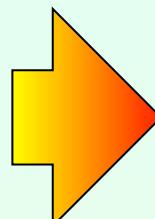
**$\infty$  MLs**



**Insulator**

**Metal**

Dimensional crossover from 3D to 2D

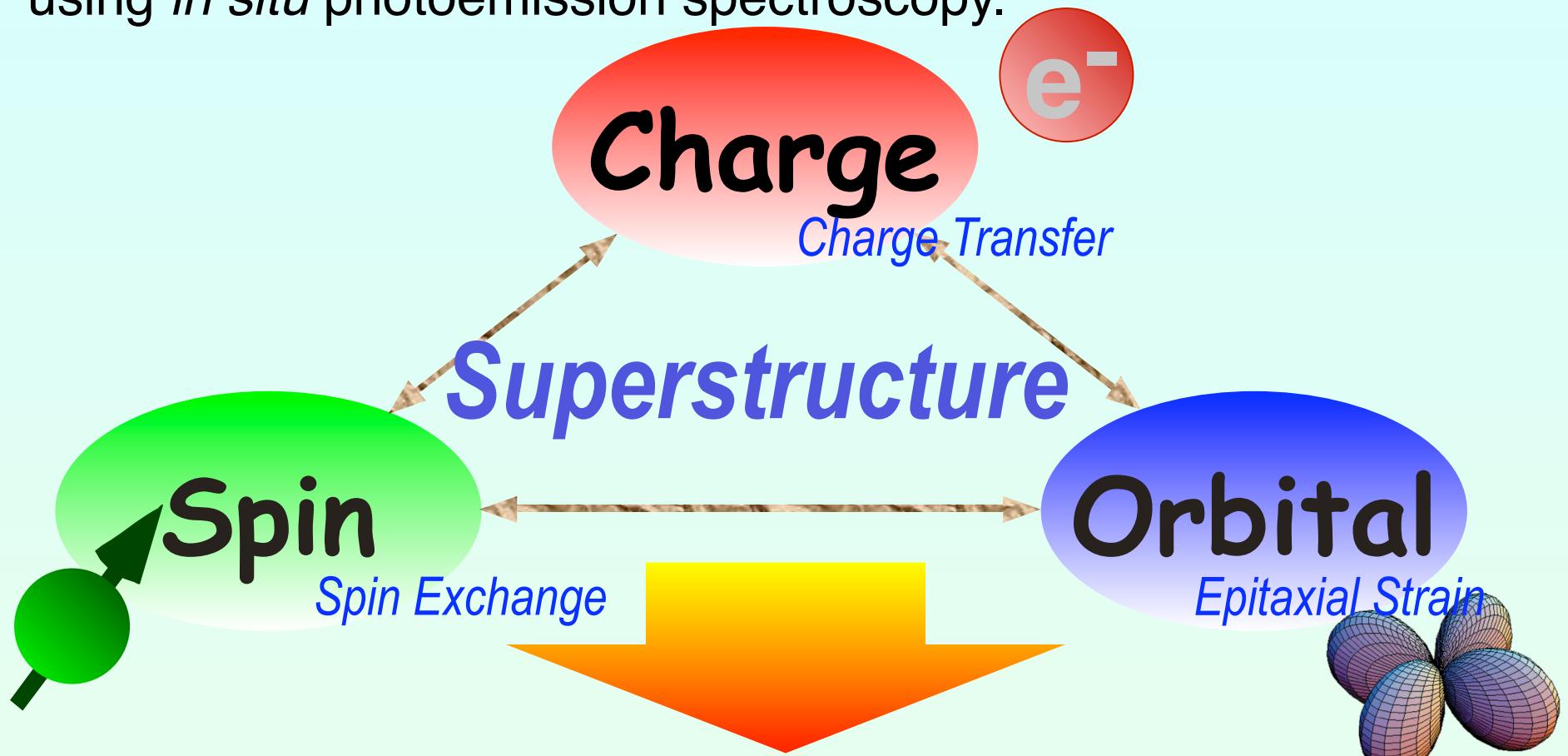


The observed MIT is caused by the reduction of bandwidth due to dimensional-crossover from 3D to 2D.

Ubiquity of thickness-dependent MIT in conductive oxide films are derived from the dimensional-crossover.

# Concluding Remarks

We studied the electronic states of oxide superstructures by using *in situ* photoemission spectroscopy.



Photoemission studies using oxide superstructures enable us to pave a new way for the better understanding of the physics of strongly correlated oxides.