

Positron and Spintronics

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When both positrons and electrons are spin-polarized, electron-positron momentum distribution exhibits asymmetry upon their mutual spin reversal. Annihilation of positronium formed on metal surface also shows spin-reversal asymmetry. These properties are demonstrated to be useful in studying ferromagnetic band structure and surface magnetism, respectively. Here, we call positron annihilation spectroscopy, which particularly uses the spin dependent annihilation process, spin-polarized positron annihilation spectroscopy (SP-PAS). In the spintronics field, SP-PAS will be a potential tool in revealing spin-related phenomena, such as magnetoresistance, current-induced spin polarization, spin-injection, vacancy-induced magnetism, half-metal band structures and so on. To promote spintronics study with SP-PAS, spin-polarized positron beam is needed. In this talk, I will report the development of spin-polarized positron beam, some fundamental aspect of SP-PAS and its applications to spintronics study performed so far.

陽電子とスピントロクス

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陽電子と電子の両方がスピン偏極している場合、電子－陽電子運動量分布は互いのスピン反転に対して非対称性を示す。金属表面で形成されるポジトロニウムの消滅も、同様にスピン反転非対称性を示す。これらの特性は、強磁性バンド構造や表面磁性の研究に有用であることが示されている。ここで、スピンに依存した陽電子消滅過程を利用する陽電子消滅法を特にスピン偏極陽電子消滅法と呼ぶことにする。近年急速に進展しているスピントロニクス分野において、スピン偏極陽電子消滅法は各種のスピン現象（磁気抵抗、電流誘起スピン分極、スピン注入、空孔誘起強磁性、ハーフメタルバンド構造など）を解明する上で、有用なプローブになると期待される。スピン偏極陽電子消滅法を用いてスピントロニクス研究を推進するためには、スピン偏極陽電子ビームが必要である。本講演では、スピン偏極陽電子ビームの開発、及び、スピン偏極陽電子消滅法の基礎とこれまで講演者等が行った幾つかの応用研究について報告し、将来の展開を模索したい。

**陽電子＋スピントロニクス
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◆About SP-PAS

◆Our Research and Development

◆3d and 4f ferromagnets

◆Spin Hall effect on Pt surfaces

◆Summary

◆About SP-PAS

Historical background
ACAR : Ferromagnetic band structures since 1957
(USA, Netherland, Switzerland, UK, India, Japan ...)
Surface Ps: Surface magnetism in 1980
(USA)

Present needs in materials science

Spintronics R & D

Bulk/Thin film
Half Metals

Vacancy-induced magnetism
DMS, SnO₂, CeO₂...

Surface/Interface
Magnetoresistance
Spin Hall effect
FM NM Spin-injection

◆About SP-PAS

Case A Internal annihilation

After Berkov 1967

MD in positive and negative fields

$$N_{\text{e-e}}(P_z) = \frac{1}{4} \sum_{i=1}^{\infty} \left[\frac{(\pm P_z)N_i^+(\mu_i)}{\lambda^0} + \frac{(\mp P_z)N_i^-(\mu_i)}{\lambda^0} \right]$$

Spin dependent lifetime

$$\lambda^{(\mp)} = \frac{1}{2} \sum_{i=1}^{\infty} [\lambda_2 w_i^{(\mp)} + \lambda_3 (w_i^{(\mp)} + 2w_i^{(0)})]$$

Wavefunction overlapping

$$N_i^{\pm}(\mu_i) = \int_{-\infty}^{\infty} N_i^{\pm}(P_z) dP_z$$

Differential MD between maj. and min. bands

$$\sum_{i=1}^{\infty} [N_i^+(\mu_i) - N_i^-(\mu_i)] \propto \Delta N \cdot P \frac{\lambda^0 - \lambda^{\pm}}{\lambda^0 + \lambda^{\pm}} \Sigma N$$

$$\Delta N = N_{+}(P_z) - N_{-}(P_z) \quad P^0 = \frac{N_{+}^0 - N_{-}^0}{N_{+}^0 + N_{-}^0}$$

Asymmetry of 3γ fraction

$$\Delta N = N_{+}(P_z) - N_{-}(P_z) \quad P^0 = \frac{N_{+}^0 - N_{-}^0}{N_{+}^0 + N_{-}^0}$$

Annihilation Lifetime

$$\Delta L(t) = L_{+}(t) - L_{-}(t) \neq 0$$

Lifetime spectra in positive and negative fields

$$L_{\text{e-e}}(t) = \frac{\lambda_2}{4} \sum_{i=1}^{\infty} [w_i^{\pm} (\pm P_z) \exp(-\lambda_i^0 t) + w_i^{\mp} (\mp P_z) \exp(-\lambda_i^0 t)]$$

Case B Surface Ps annihilation

c.f. Gidley 1980

Ps fractions

$$F_{\text{Ps}} = (1 - P_z P_{\perp} \cos \phi) / 4$$

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$$F_{\text{Ps}} = (1 + P_z + P_{\perp} \cos \phi + P_z P_{\perp} \cos \phi) / 4$$

$$F_{\text{Ps}} = (1 - P_z - P_{\perp} \cos \phi + P_z P_{\perp} \cos \phi) / 4$$

Total 3γ fraction

$$F_{\text{3}\gamma} = \varepsilon(1)(F_{\text{Ps}} + F_{\text{1}\gamma}) + \varepsilon(0)F_{\text{1}\gamma}$$

Doppler meas.

$$R = \frac{(1 - F_{\text{Ps}})^2 R_0 + F_{\text{Ps}}^2 R_1 P_z / P_0}{1 - F_{\text{Ps}}^2 + F_{\text{Ps}}^2 P_z / P_0}$$

$$\Delta R = R - R_0 \propto F_{\text{Ps}}^2 \text{ for small } F_{\text{Ps}}$$

Lifetime meas.

$$I_{\text{1}\gamma\text{tot}} = F_{\text{Ps}}^2$$

Asymmetry of 3γ fraction

$$A = \frac{\Delta R(+)-\Delta R(-)}{\Delta R(+) + \Delta R(-)} = \frac{I_{\text{1}\gamma\text{tot}}(+)-I_{\text{1}\gamma\text{tot}}(-)}{I_{\text{1}\gamma\text{tot}}(+) + I_{\text{1}\gamma\text{tot}}(-)}$$

$$= \frac{2\varepsilon(1)-\varepsilon(0)}{2\varepsilon(1)+\varepsilon(0)} P_z P_{\perp} \cos \phi$$

Surface polarization directly determined

◆Our R&D

Spin-Polarized Positron Source & Beam

- Highly Spin-Polarized Positrons
 ${}^{22}\text{Na}$ P=70%
 ${}^{68}\text{Ge}-{}^{68}\text{Ga}$ P=94%
- Energy tunability
10eV(0.1nm)–10keV(100nm)
- Polarization switchability
Longitudinal/Transverse

Foundation of SP-PAS method

- Influence of magnetic field on positron annihilation
- Theoretical calculation

3d(Fe, Co, Ni), 4f(Gd, Tb, Dy)...
First principles calculation

Application of SP-PAS to Spintronics Studies

- Spin Hall effect
- Spin-injection
- Heusler alloys
- Vacancy-induced magnetism
- etc

DBAR measurement

◆Our R&D

Flux $\sim 5 \times 10^3 \text{ e}^+/\text{s}$
P: 30~35%

◆SP-PAS on 3d and 4f ferromagnets

Field reversal asymmetry of momentum distribution

$$N_{+}(P_z) - N_{-}(P_z) = \frac{P}{2} \sum_{i=1}^{\infty} \left[\frac{N_i^+(\mu_i)}{\lambda^0} - \frac{N_i^-(\mu_i)}{\lambda^0} \right]$$

◆SP-PAS on 3d and 4f ferromagnets

Intensity Fe > Co > Ni
Intensity Intensity 2.2 : 1.6 : 0.4
M_S (μ_B) 2.2 : 1.7 : 0.6

Ferromagnetic band structure PRB83(2011)100406(R). PRB83(2012)024417.

