

テラヘルツ時間領域分光法と 基礎科学への応用

Department of Physics, Kyoto University
Koichiro Tanaka

SCOPE project
2003-2007



Advanced Terahertz Technologies
adapted for Optical Communication



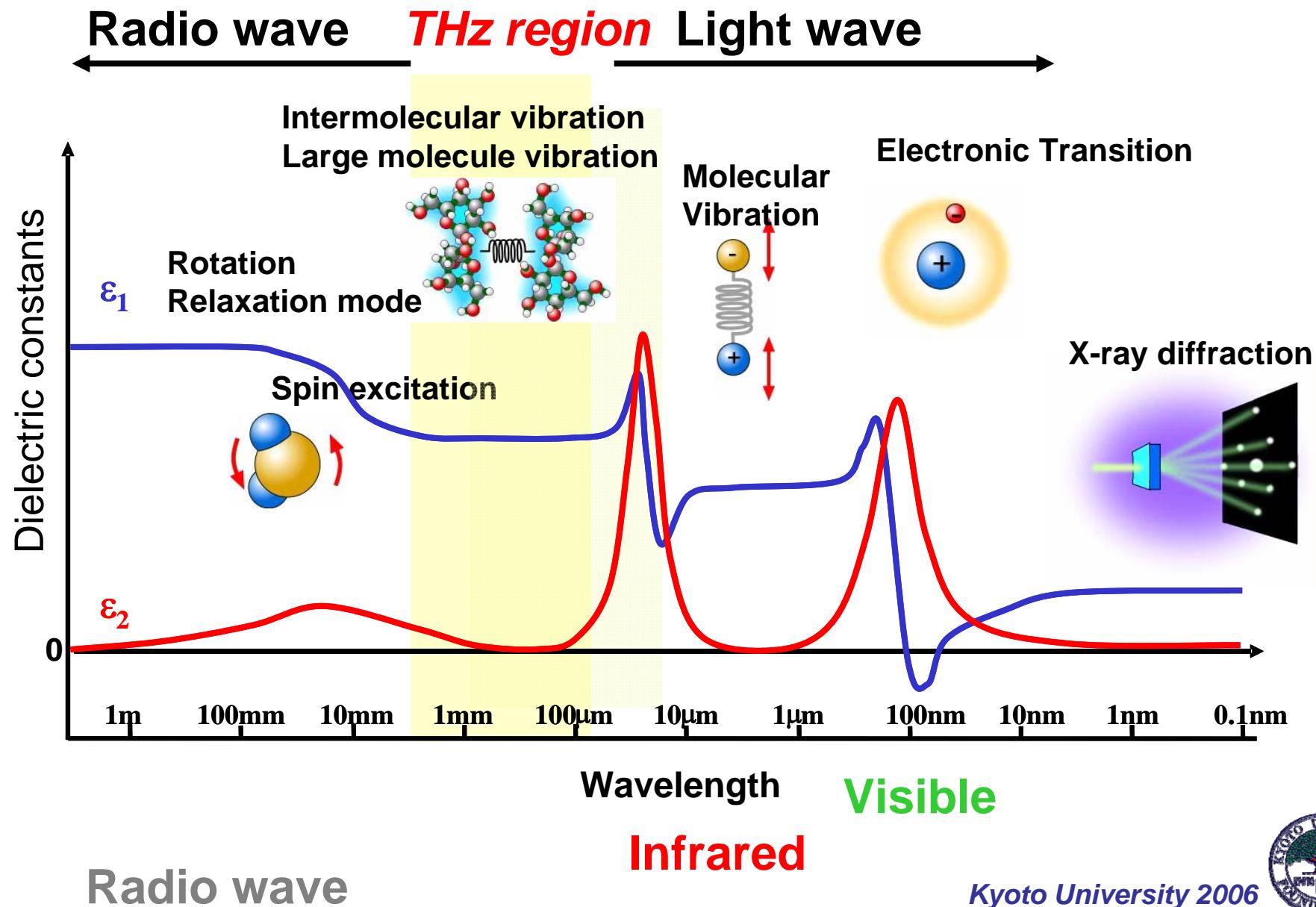
JSPS project
2006-2010

Exploitation of Organic Electronic Materials
of Potential Dynamic Switches for Non-
equilibrium Condensed Matter Sciences

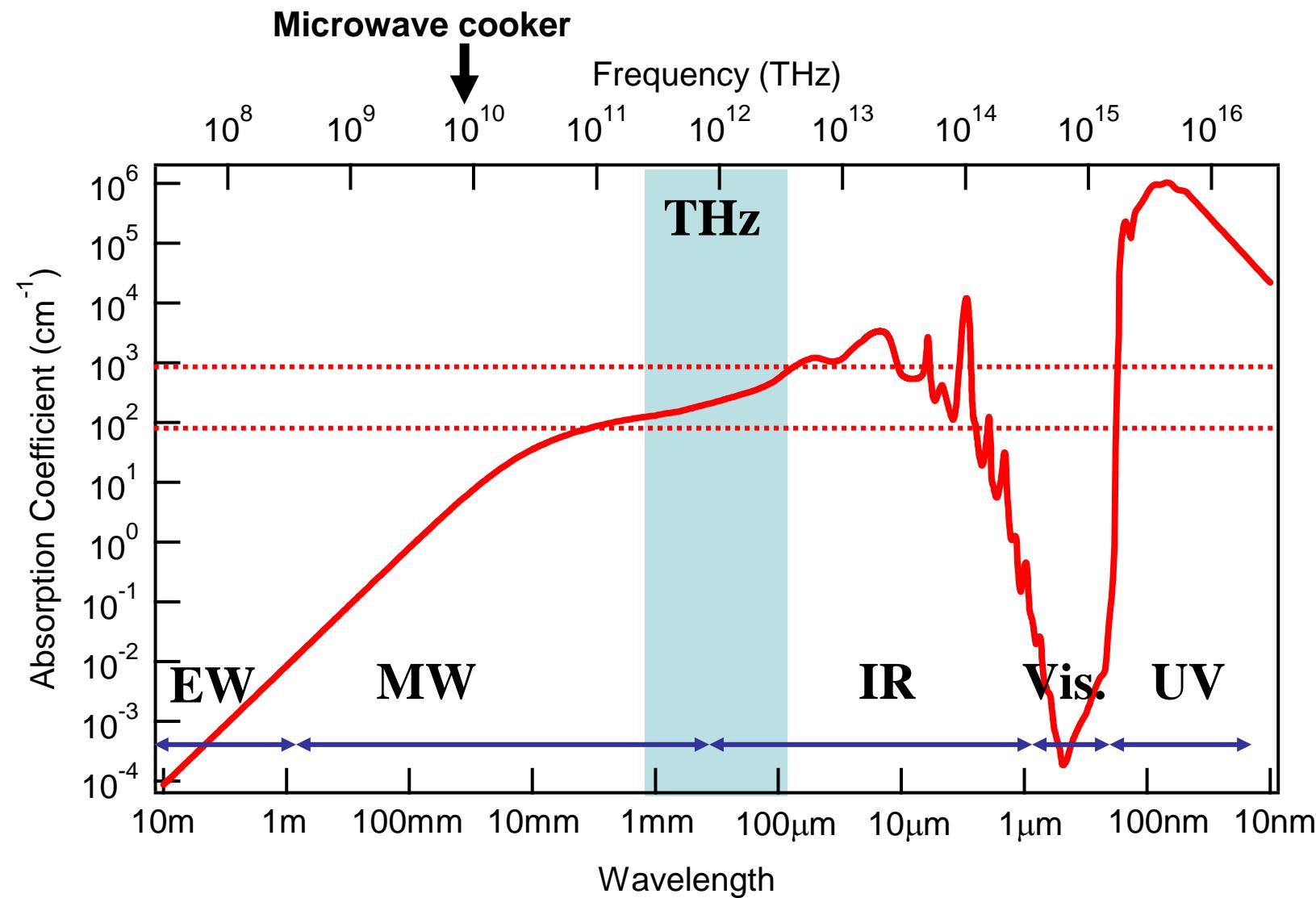
講演内容

- テラヘルツテクノロジーとは何か？
- 時間領域テラヘルツ分光法の基礎
- 時間領域テラヘルツ分光法による物性測定
 - 有機材料、強誘電体、半導体、超伝導体
 - 水、水溶液
- 非線型テラヘルツ分光の必要性と戦略

Importance of terahertz frequency region

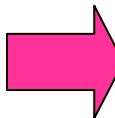
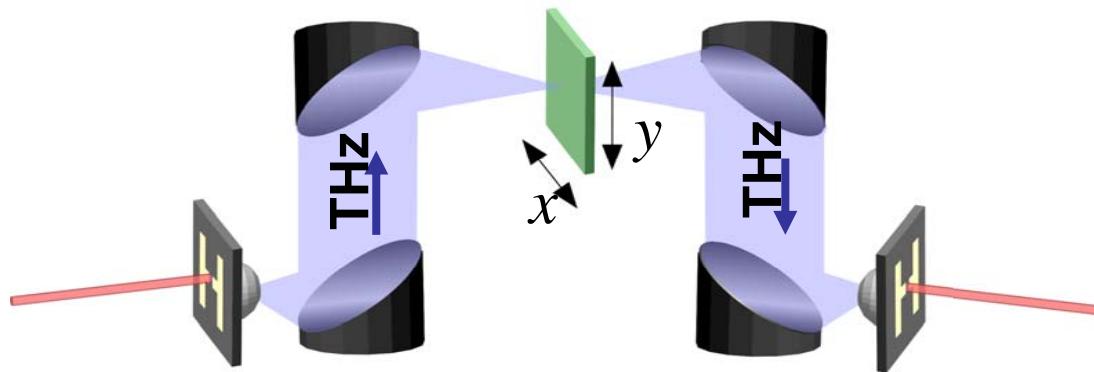


Absorption of Liquid Water



Jackson, Classical Electrodynamics

THz-Imaging



Much of current THz research revolves around spectral specificity and transmission properties. The THz frequency



B. B. Hu and M. C. Nuss, Opt. Lett. 20, 1716 (1995)

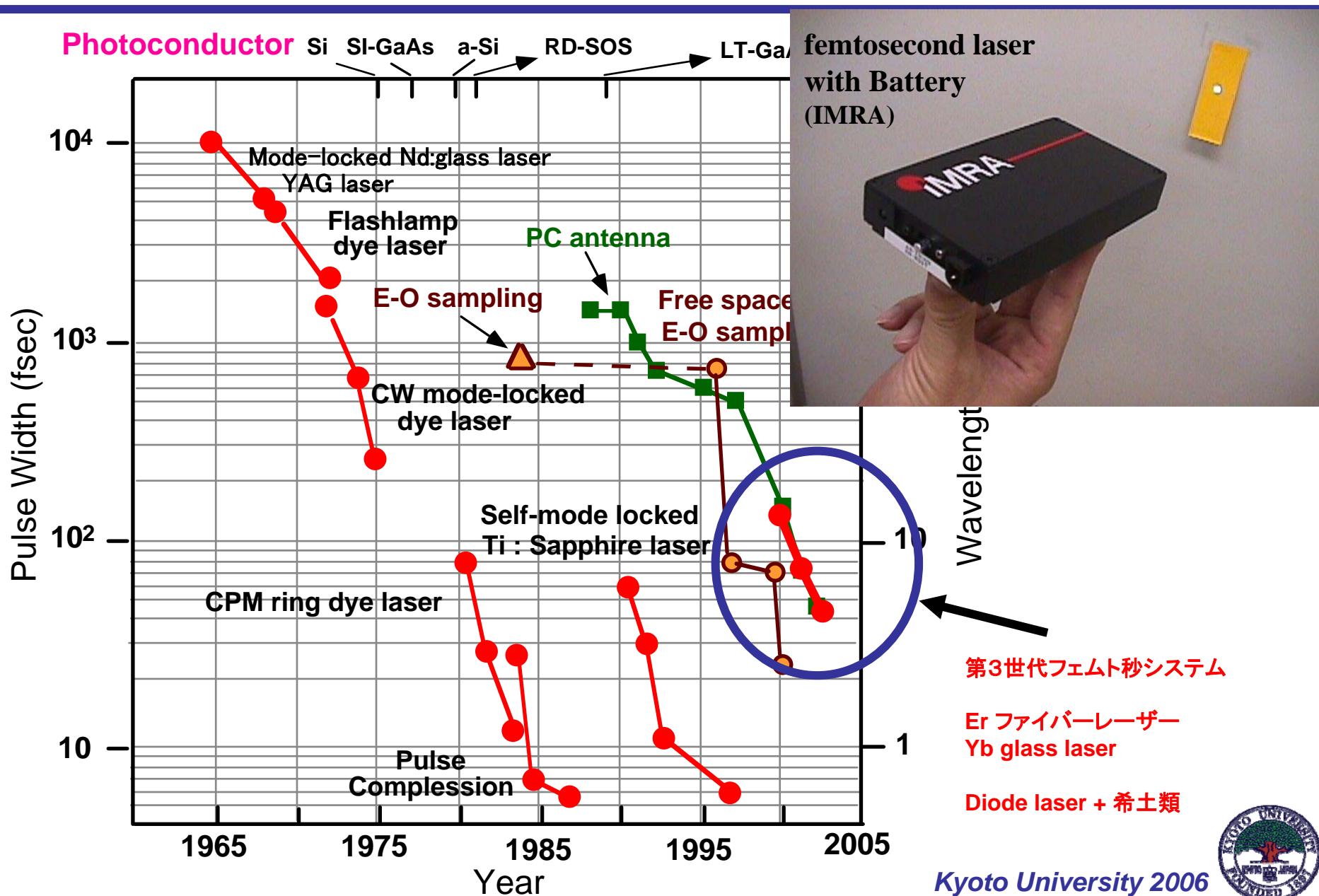
Kawase (Nagoya/Riken), Nikon

講演内容

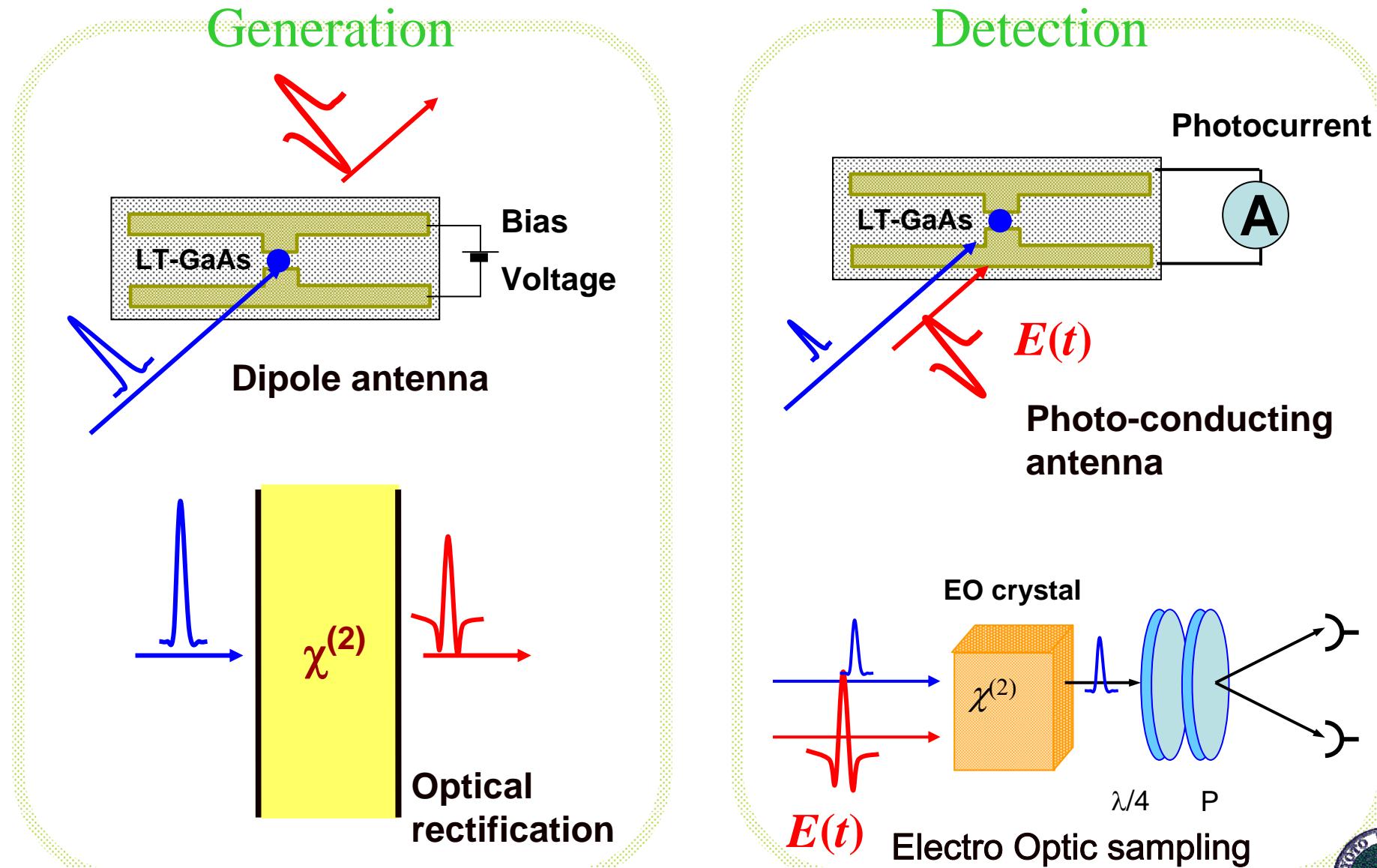
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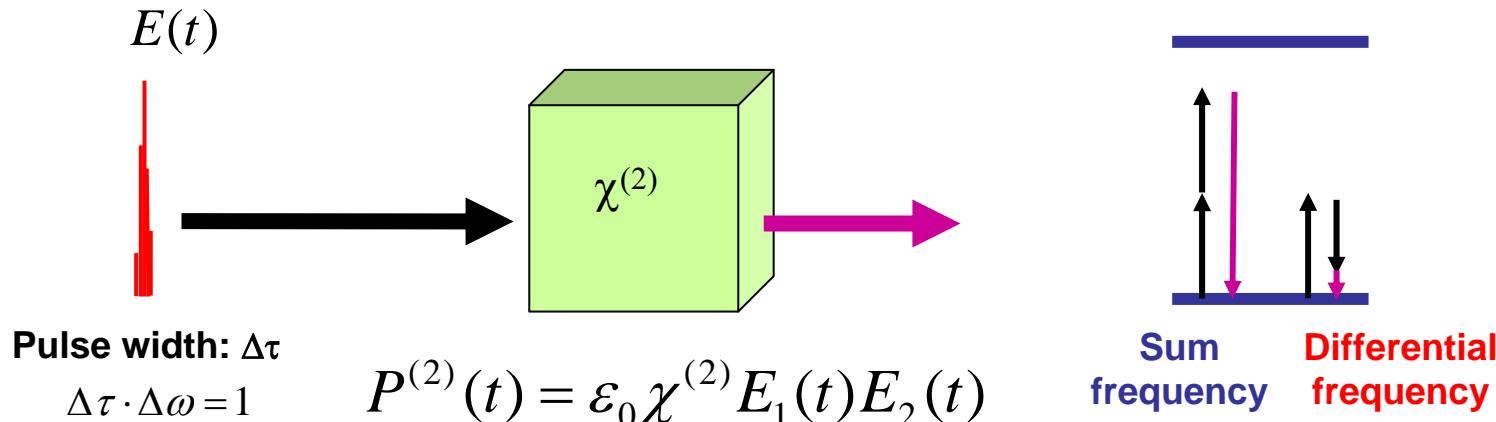
THz-TDSの進展 - フェムト秒レーザーの進歩



THz-detection



THz-wave generation with second-order non-linear optical process



$$E_1(t) = E_2(t) = A(e^{i\omega t} + e^{-i\omega t})$$

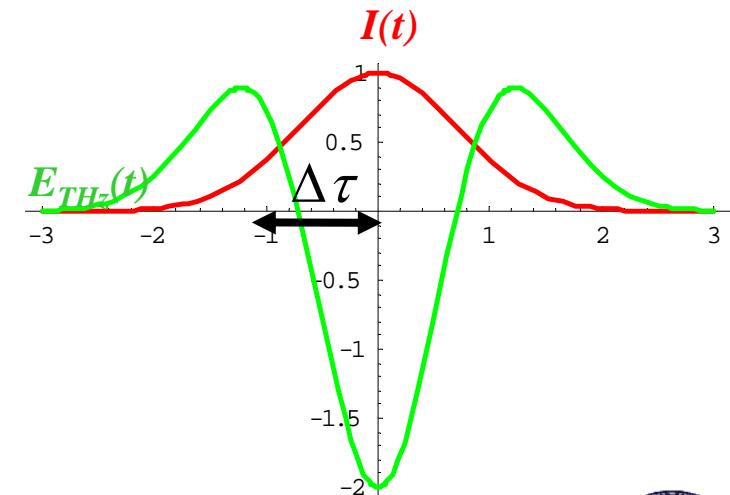
$$P^{(2)}(t) = A^2 (\underbrace{\chi^{(2)}(2\omega, \omega, \omega) e^{i2\omega t}}_{\text{SHG}} + \underbrace{\chi^{(2)}(0, \omega, -\omega)}_{\text{Optical rectification}})$$

$$P^{(2)}(2\omega) \quad P^{(2)}(0)$$

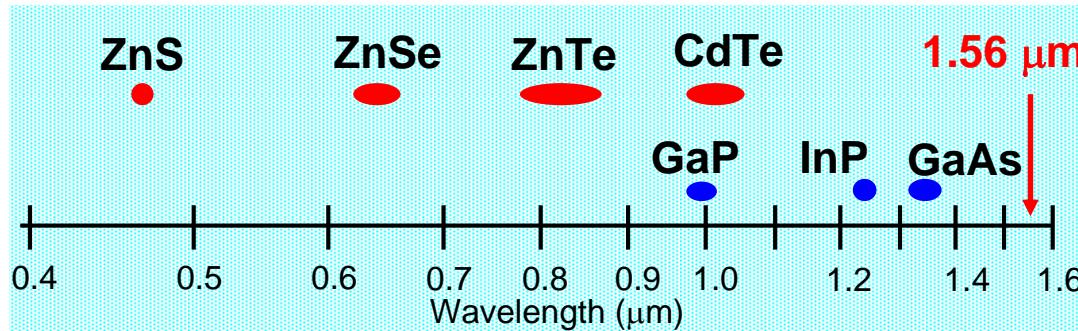
$$\Delta\tau \cdot \Delta\omega = 1$$

$$\Omega < \Delta\omega$$

$$P^{(2)}(t) = A^2 \chi^{(2)}(\Omega, \omega + \Omega, -\omega) e^{i\Omega t}$$



Phase matching condition



Phase match condition:

$$k(\omega+\Omega) - k(\omega) - k(\Omega) = 0$$

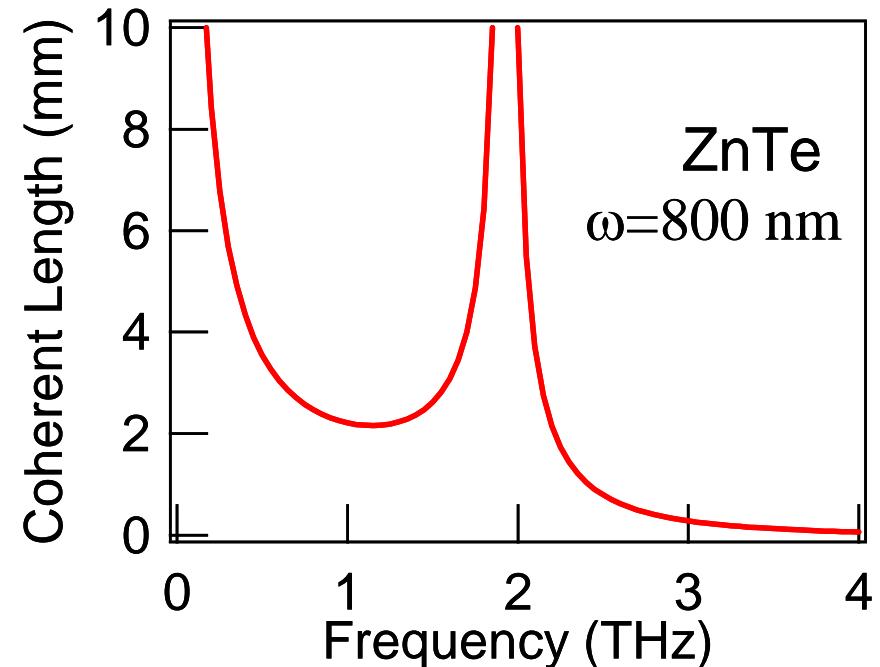
Including dispersion:

refractive index
in THz region = optical Index of
group velocity

$$V_{phase}^{THz} = V_{group}^{visible}$$

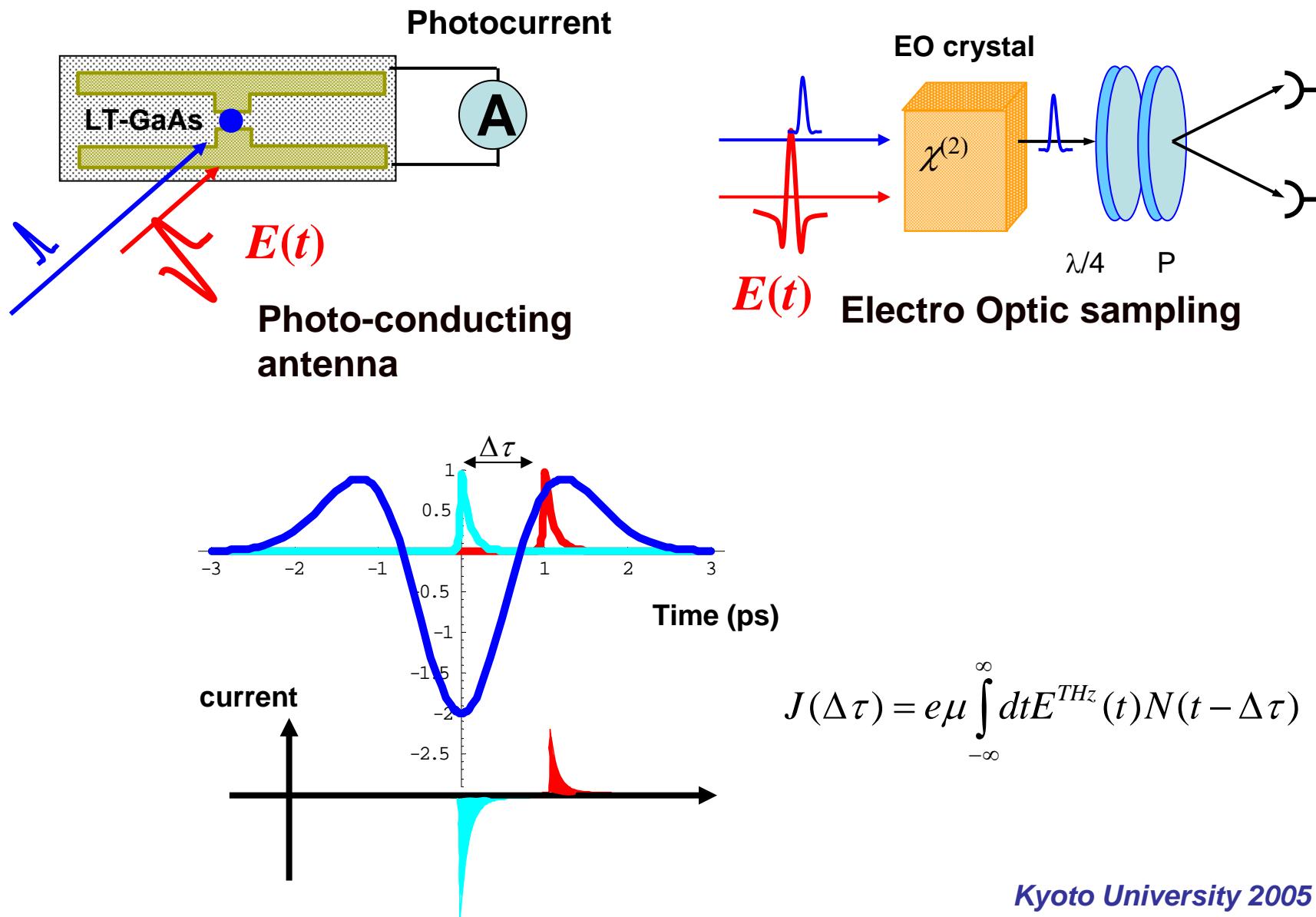
Coherent length:

$$l_c = \frac{c}{2f|n_{THz} - n_g|}$$

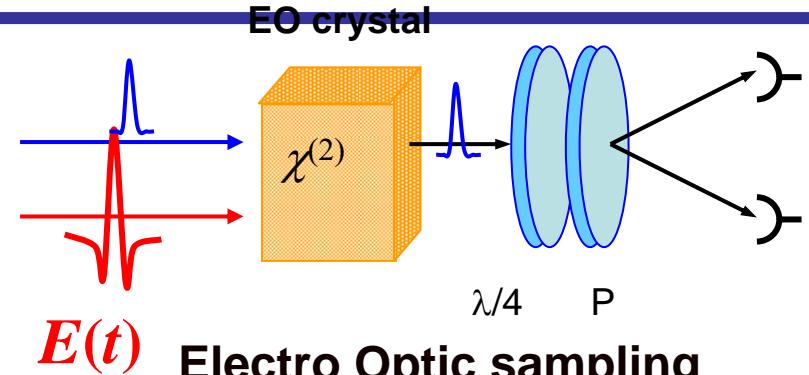


Nahata et al., APL **69**, 2321 (1996).

THz-wave detection



Non-linear THz spectroscopy



State of the art in our laboratory

(Kumiko YAMASHITA, 2005) :

– Estimation was made by EO sampling technique.

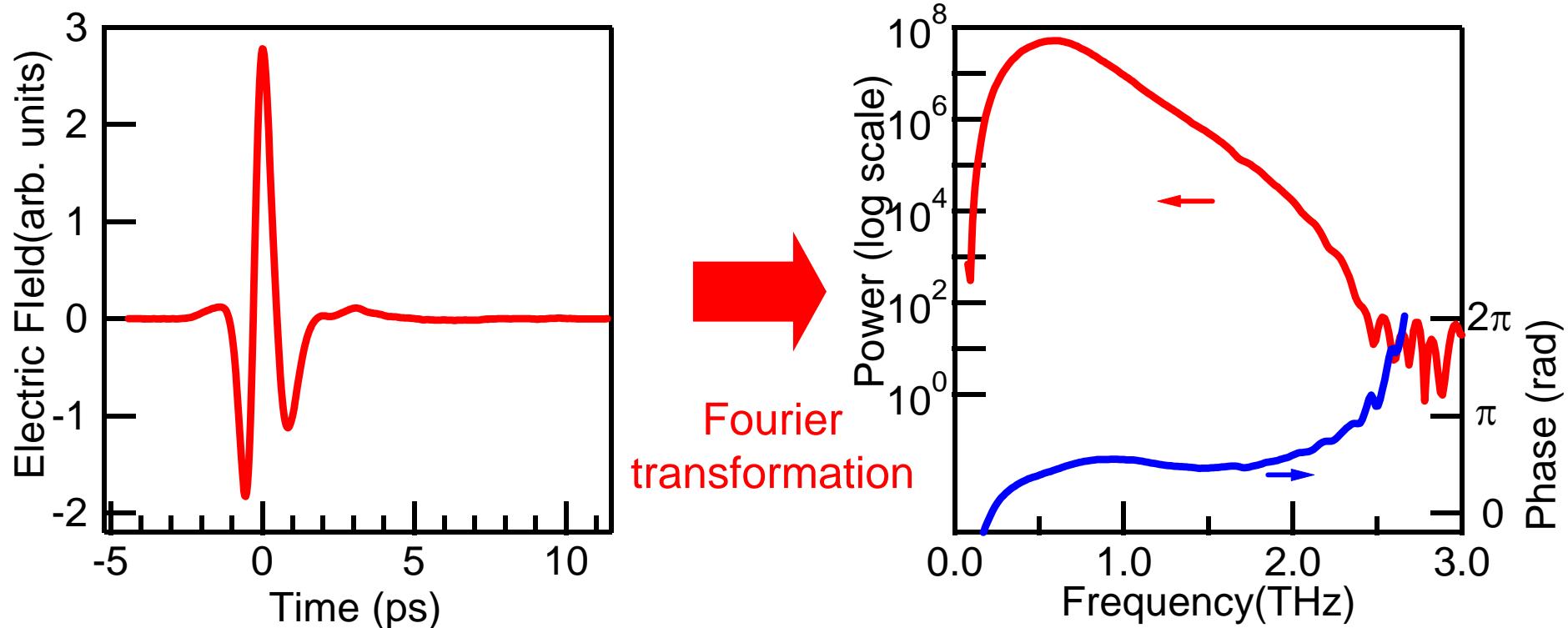
$$\frac{\Delta I}{I} = -\Gamma = -2\pi \frac{l}{\lambda} n^3 r_{41} |E_{THz}| = 0.2$$

Optical pump: $\lambda = 800\text{nm}$ 485 mW (485 $\mu\text{J}/\text{pulse}$)
ZnTe($l=1\text{mm}$ 、 $n=3.2$ 、 $r_{41}=4 \times 10^{-12}\text{m/V}$)

$$P_{THz} = 100 \text{ pJ/pulse}, 10 \text{ nJ/cm}^2, 10 \text{ kW/cm}^2, |E_{THz}| \sim 1 \text{kV/cm} = 100 \text{kV/m}$$

*Typically, for the non-linear spectroscopy in visible region,
we need a MW/cm² class laser.*

Time-domain spectroscopy (TDS) is a powerful tool in THz region

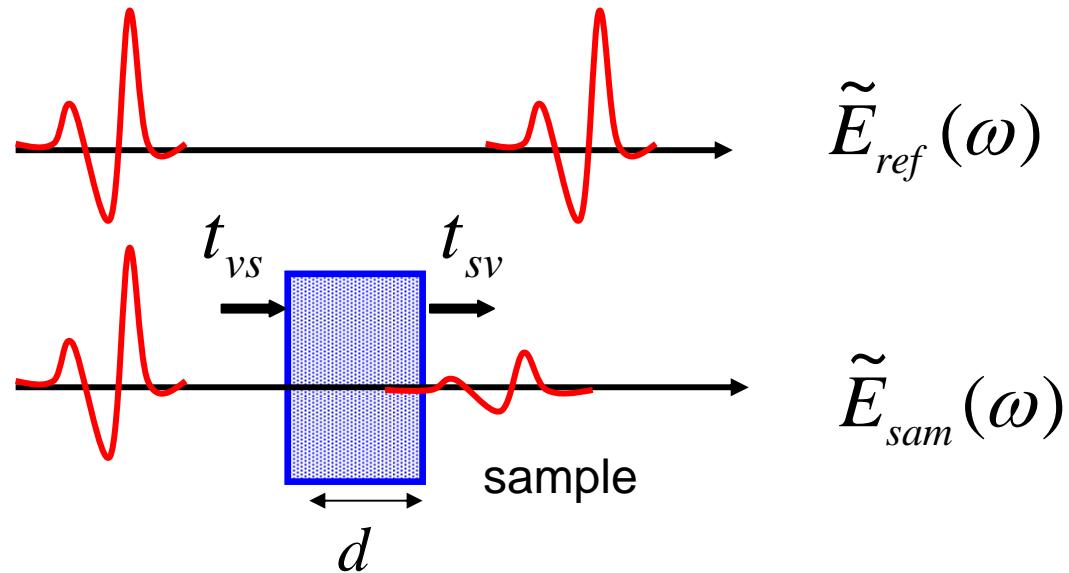


$$\tilde{E}(\omega) = E(\omega)e^{-i\phi(\omega)} = \frac{1}{2\pi} \int E(t)e^{i\omega t} dt$$

- ✓ Pulse measurements \Rightarrow High sensitivity. cf. FTIR
- ✓ Electric field measurements \Rightarrow Complex dielectric constants.

$E(\omega), \phi(\omega)$

Terahertz Time-Domain Spectroscopy (THz-TDS)



Complex refractive index

$$\begin{aligned}\tilde{t}(\omega) &= \frac{\tilde{E}_{sam}(\omega)}{\tilde{E}_{ref}(\omega)} \\ &= t_{vs} \cdot t_{sv} \exp\left[i \frac{(\tilde{n}(\omega) - 1)\omega d}{c}\right] \\ &= t_{vs} \cdot t_{sv} \exp\left[i \frac{(n - 1)\omega d}{c}\right] \exp\left[-\frac{\kappa\omega d}{c}\right]\end{aligned}$$

$$\tilde{n} = n + i\kappa$$

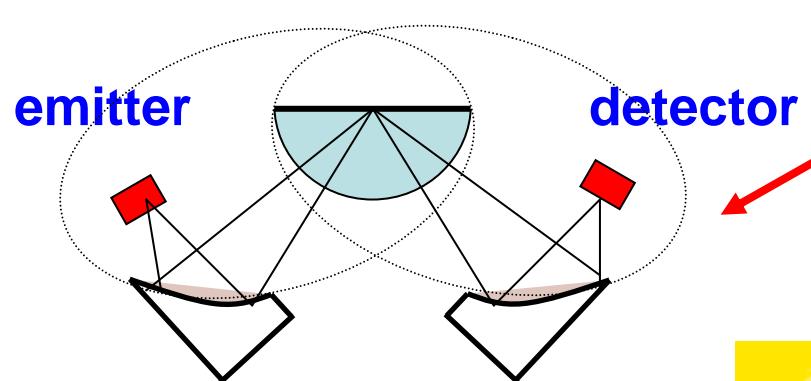
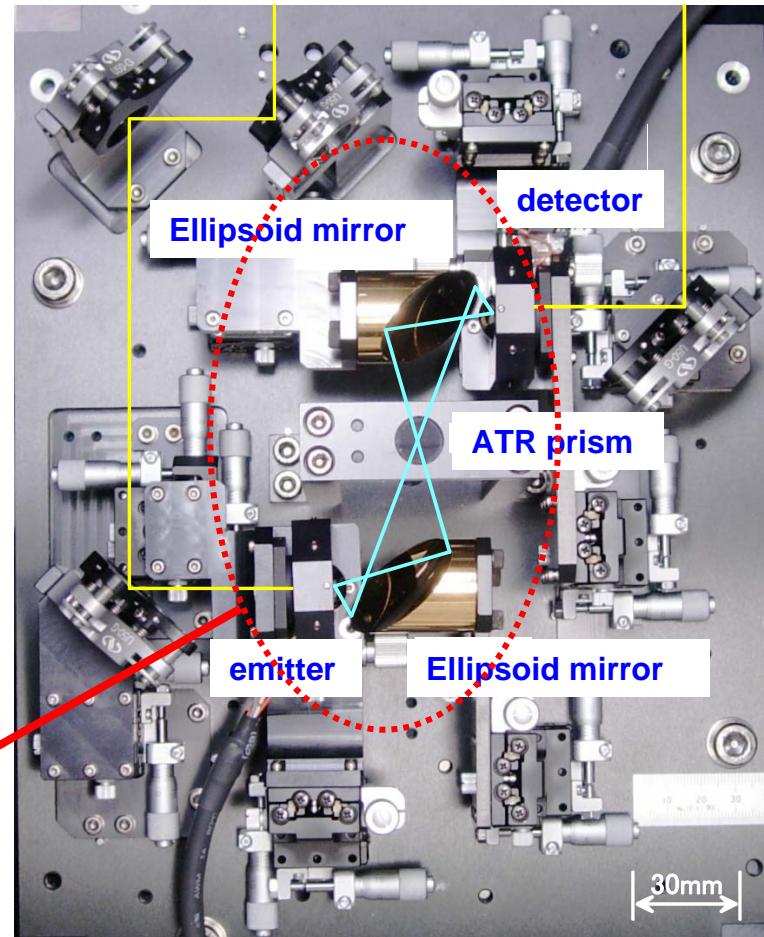
t_{vs}, t_{sv} : Fresnel coefficients

$$t_{vs} = \frac{2}{\tilde{n}(\omega) + 1} \quad (\textit{vacuum} \rightarrow \textit{sample})$$

$$t_{sv} = \frac{2\tilde{n}(\omega)}{\tilde{n}(\omega) + 1} \quad (\textit{sample} \rightarrow \textit{vacuum})$$



Compact THz-ATR spectrometer



Double-ellipsoid mirrors with
a semispherical ATR prism.



Advanced Terahertz Technologies
adapted for Optical Communication

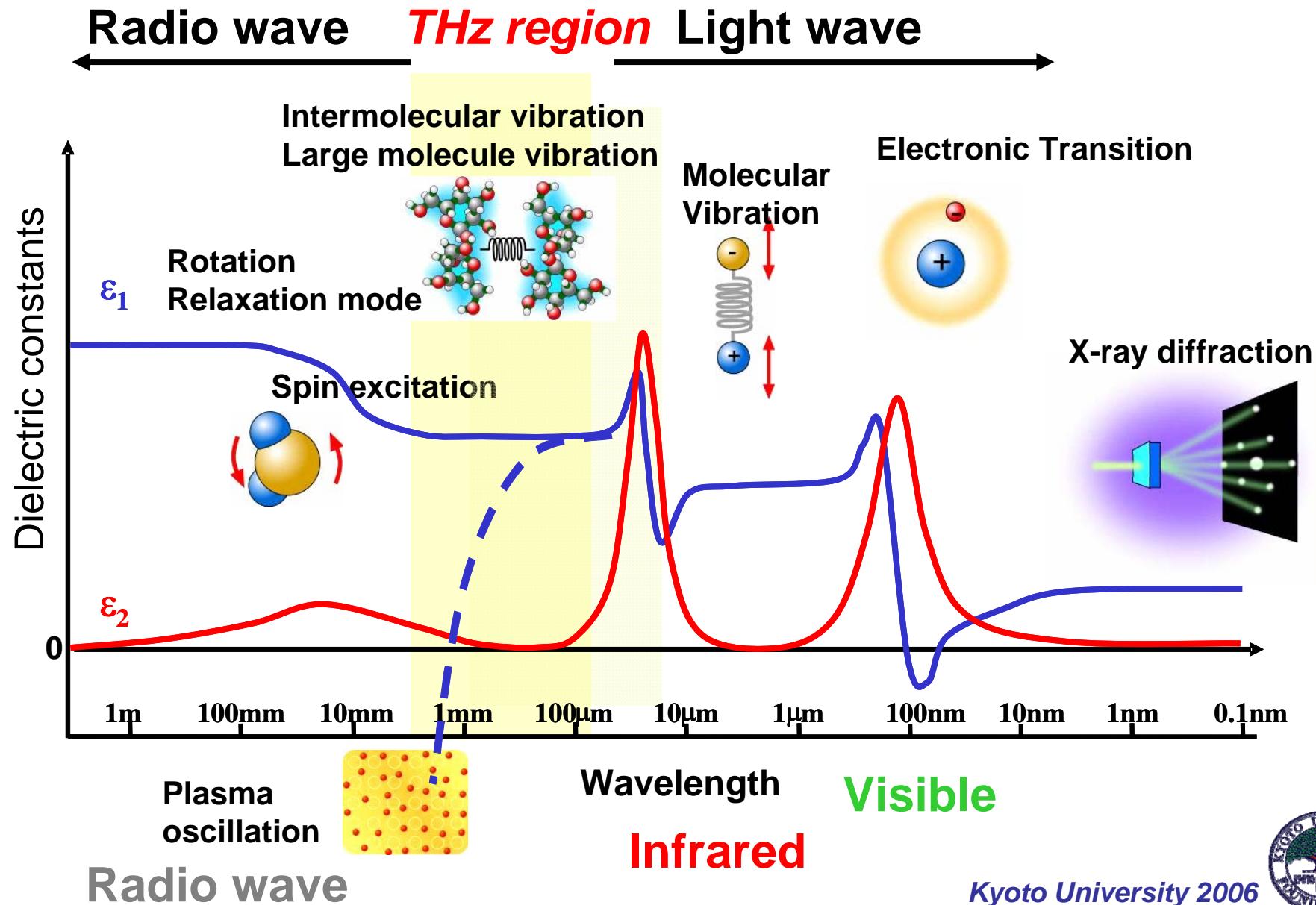


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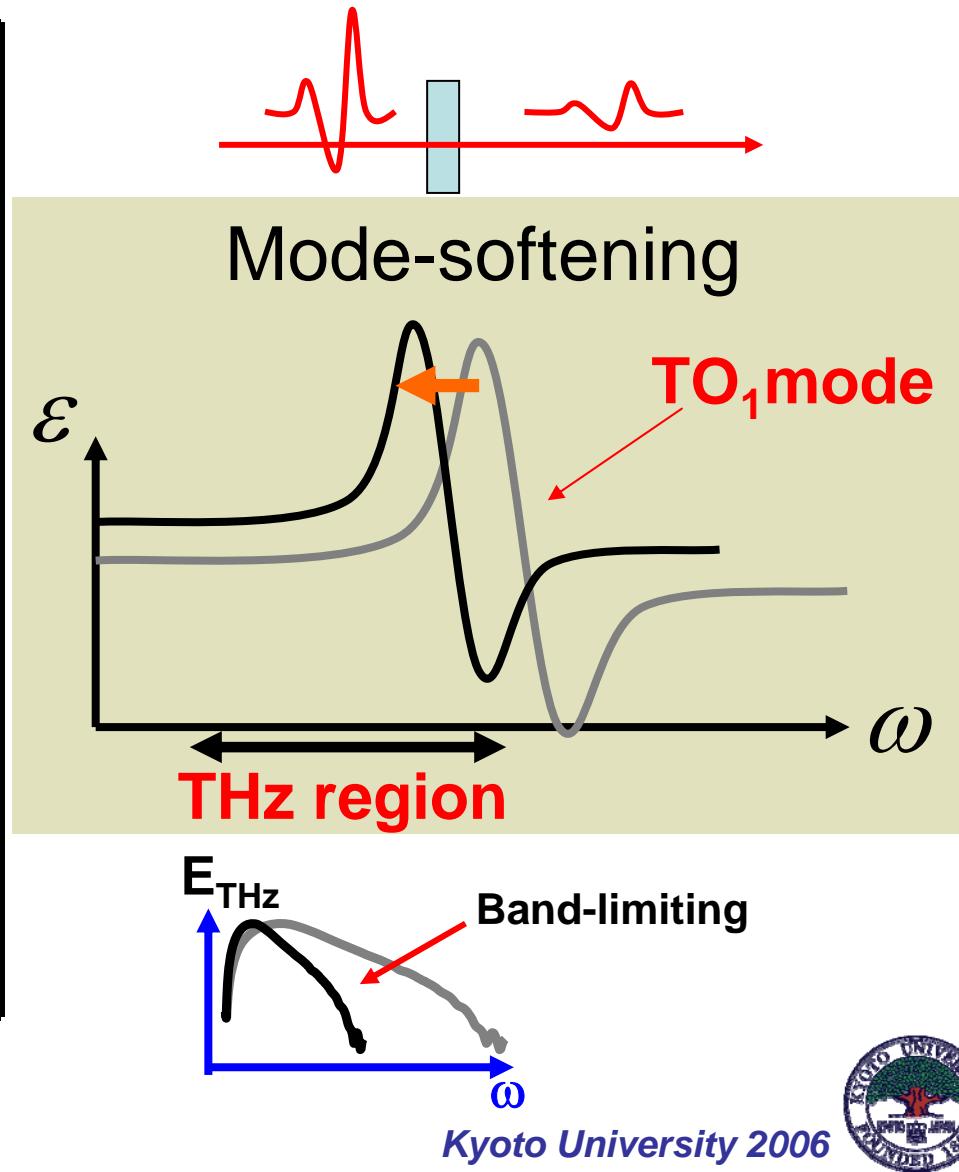
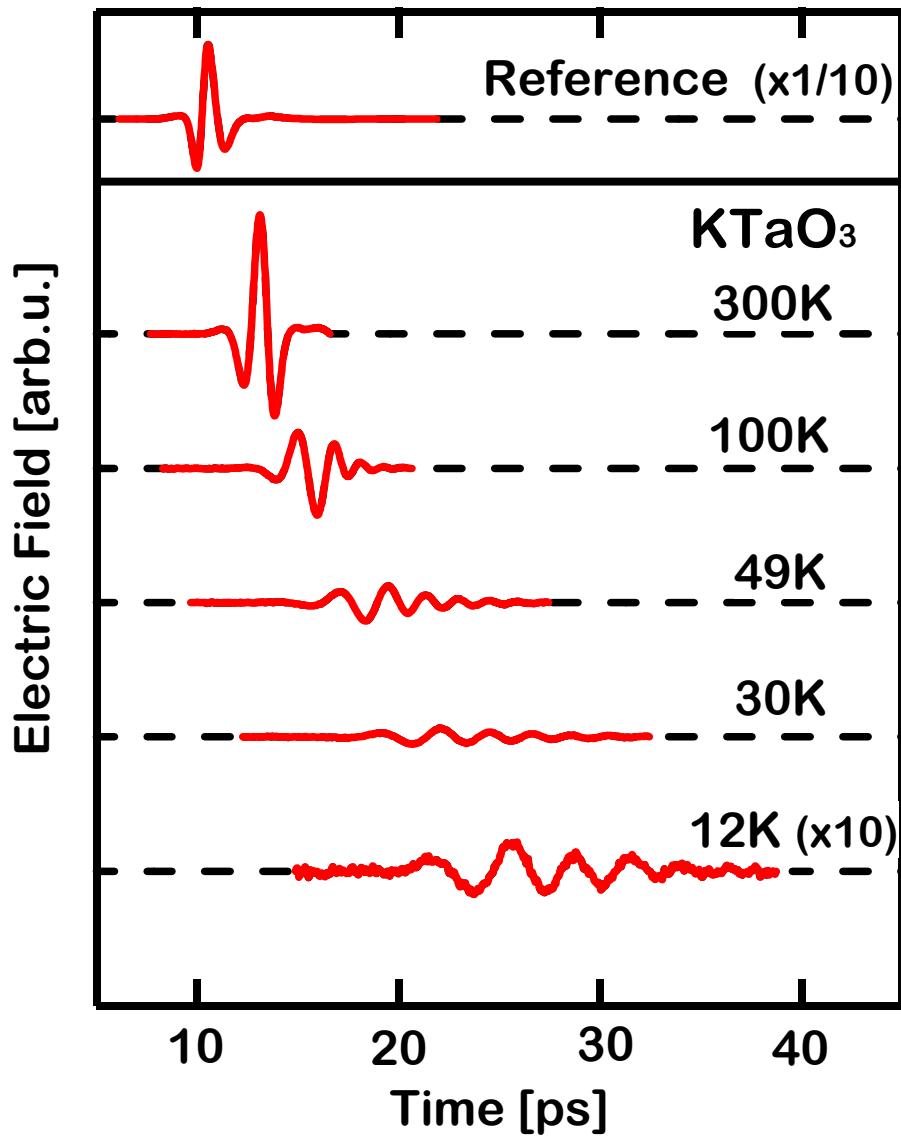


Importance of terahertz frequency region



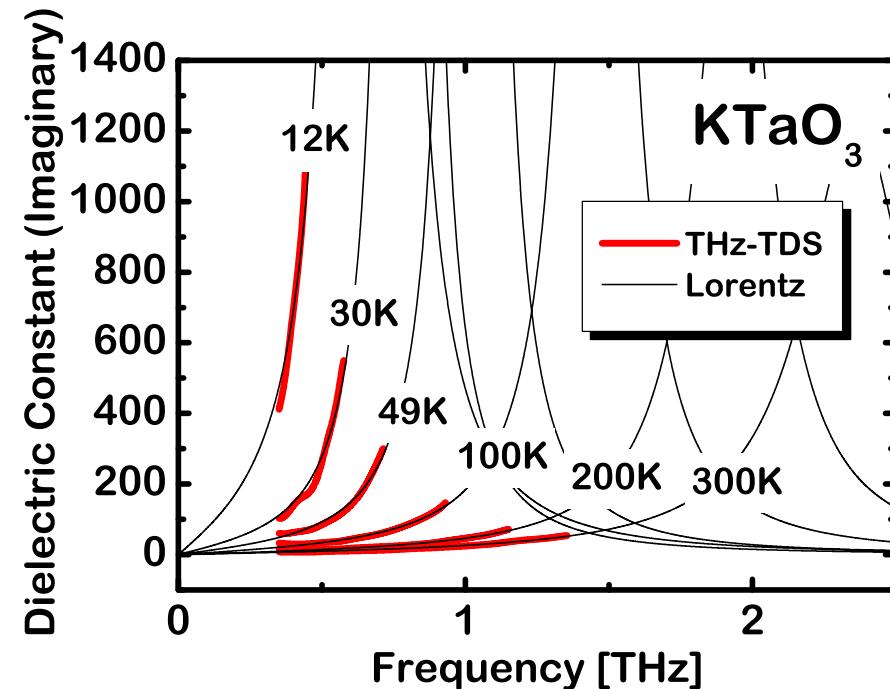
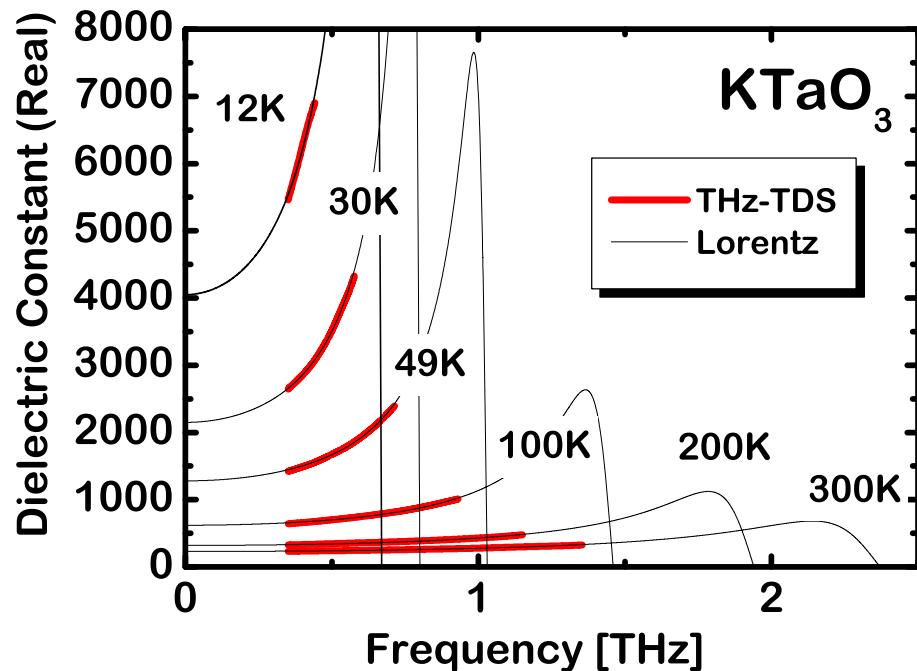
Mode-softening in $KTaO_3$ crystal

Ichikawa, Tanaka et al.: Physical Review B 71(2005) 086509.



Dielectric function in KTaO_3

Ichikawa et al.: Physical Review B 71(2005) 086509.



A single Lorentz Model

$$\tilde{\varepsilon}(\omega) = \varepsilon_{\infty} + \frac{(\varepsilon_0 - \varepsilon_{\infty})\Omega_0^2}{\Omega_0^2 - \omega^2 - i\gamma_0\omega}$$

Fitting parameters

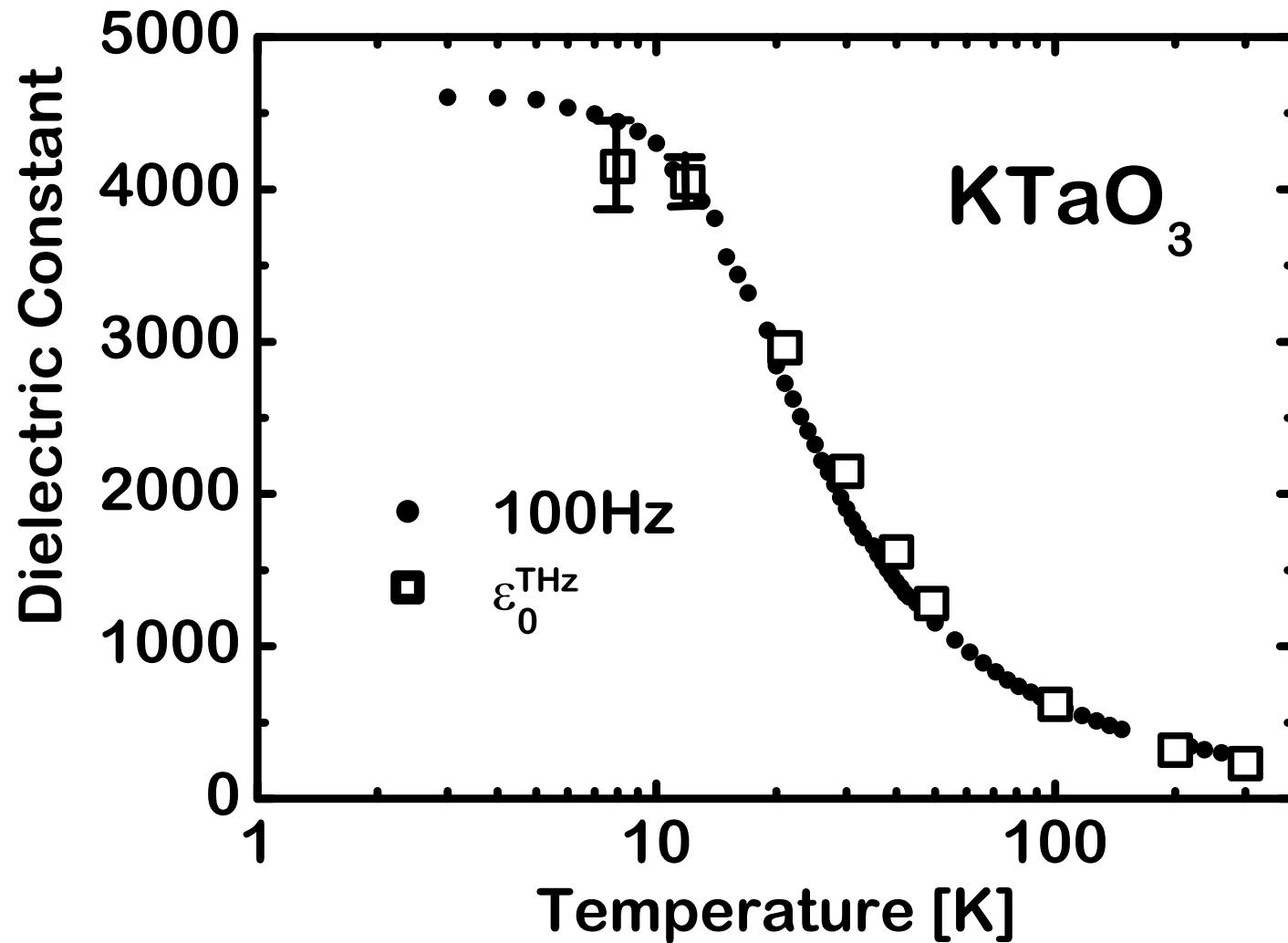
Ω_0 : Mode frequency

γ_0 : Damping

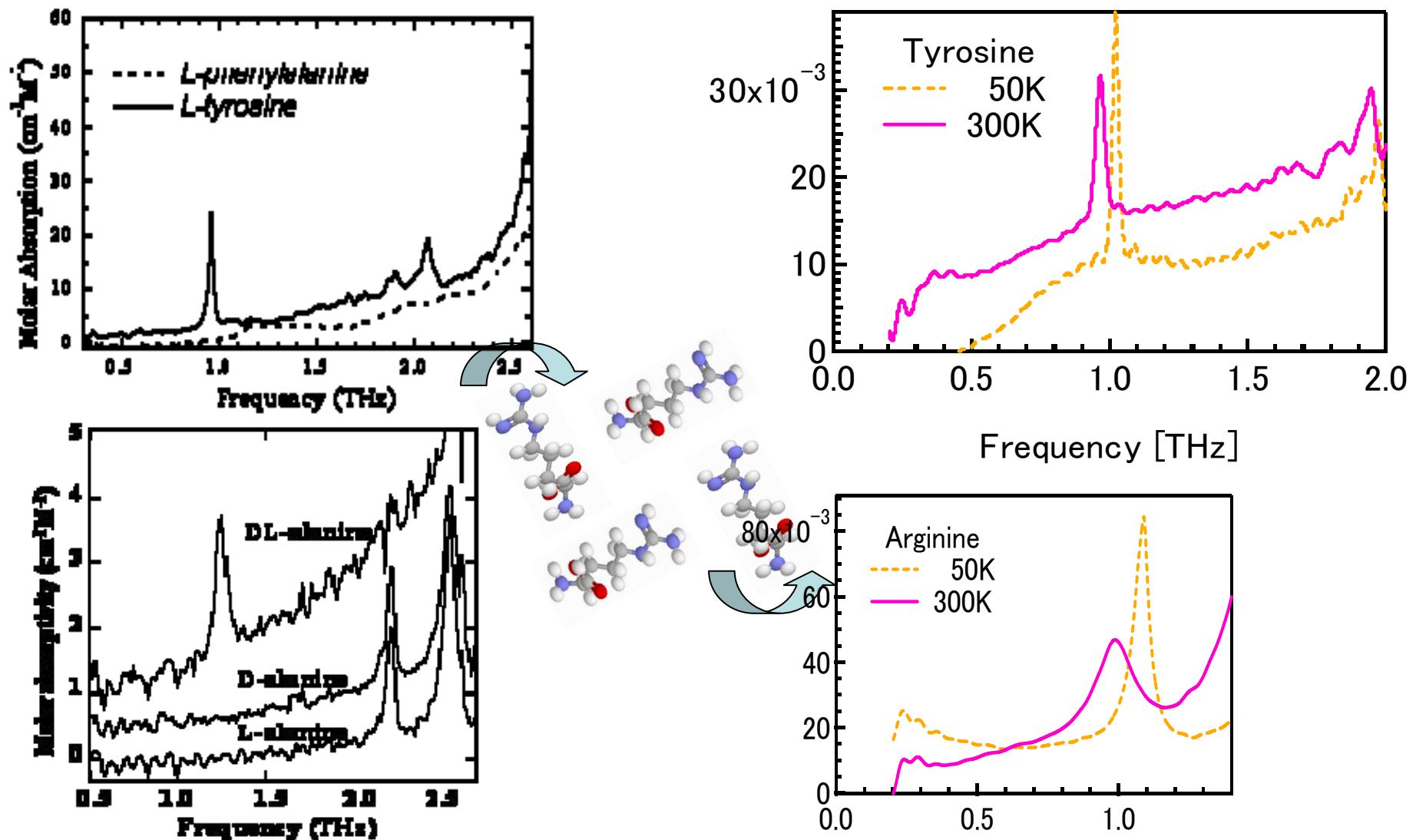
ε_0 : Static dielectric constant

Static dielectric constant

Ichikawa et al.: Physical Review B 71(2005) 086509.



アミノ酸結晶の吸収



谷他、「アミノ酸分子結晶のテラヘルツ時間領域分光」より引用

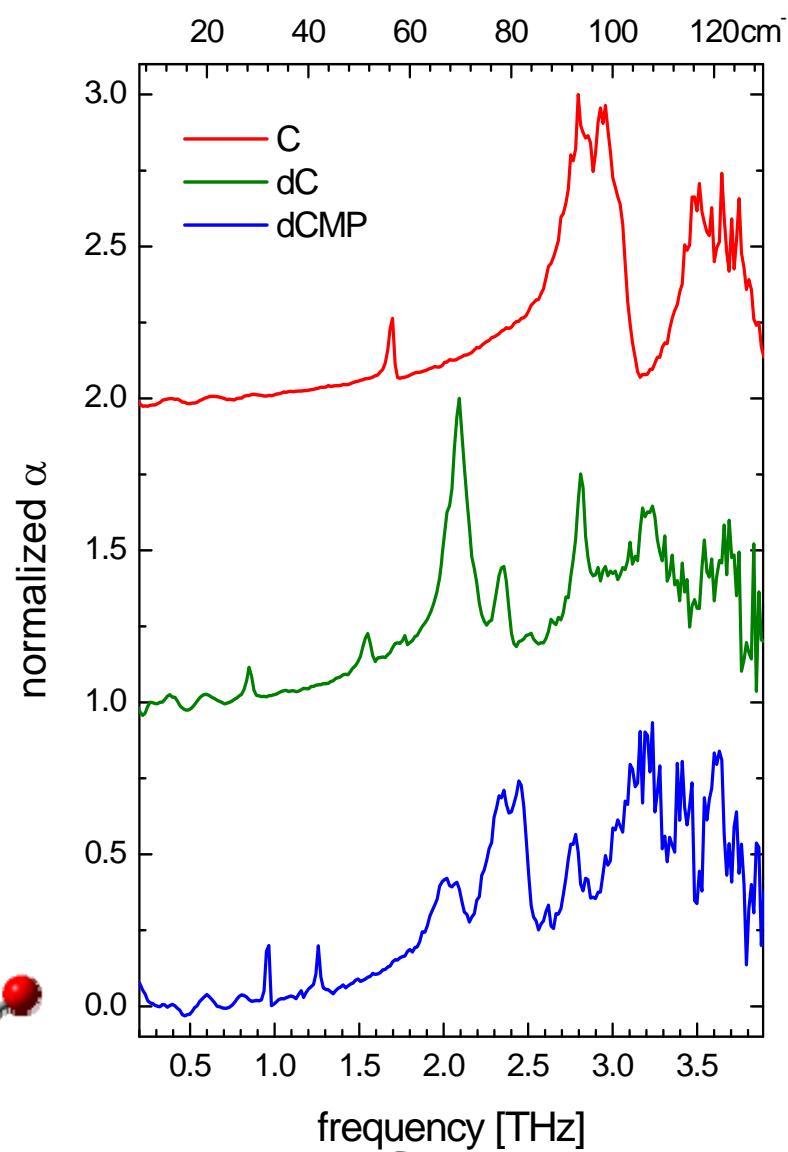
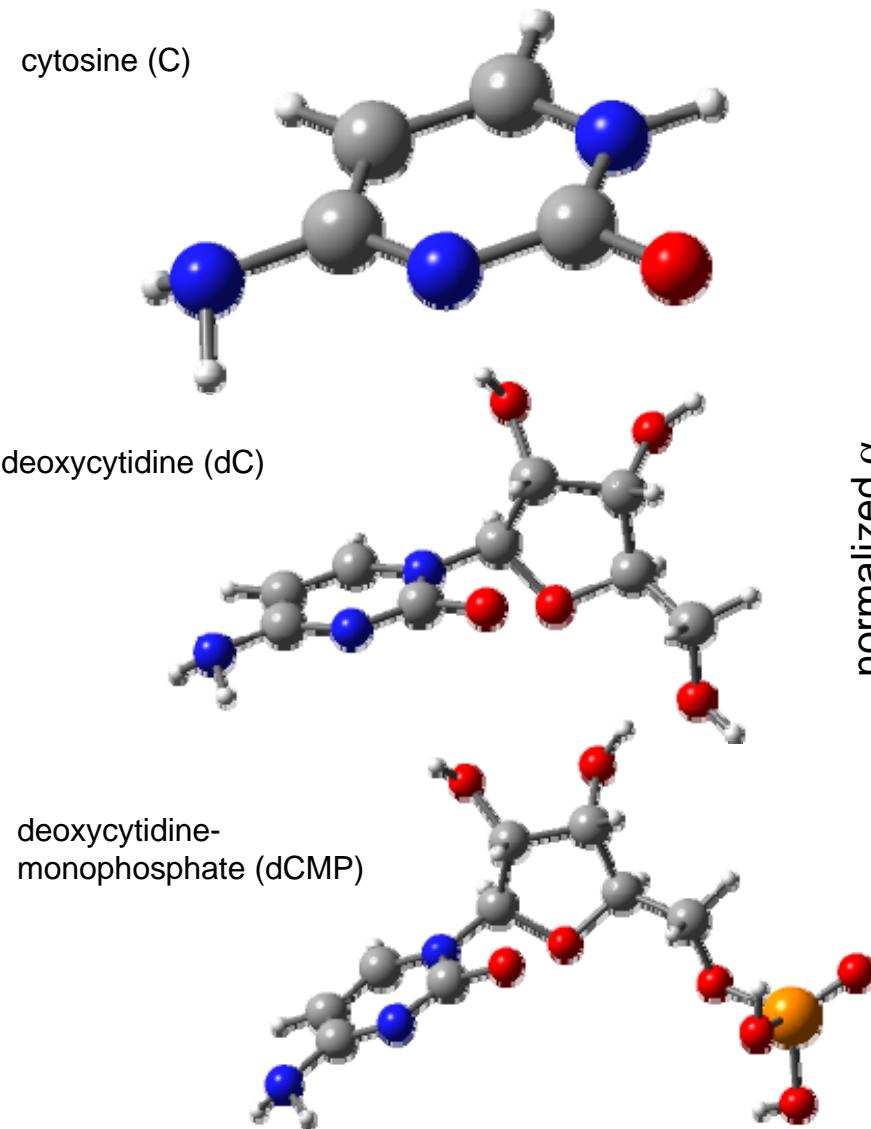
http://www.technova.co.jp/teratech/saizensen/tera_saizensen1.html

Frequency [THz]

Kyoto University 2006



isolated bases → DNA: first step



Evanescence wave in ATR

Internal total reflection condition

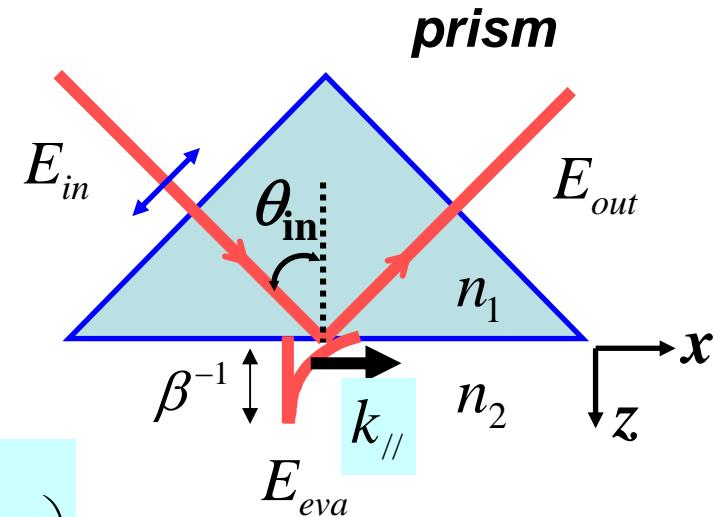
Evanescence wave
for *p*-polarization

$$\theta_{in} > \theta_c = \sin^{-1}\left(\frac{n_2}{n_1}\right)$$

$$\begin{pmatrix} H_x \\ H_y \\ H_z \end{pmatrix} = \begin{pmatrix} 0 \\ B \\ 0 \end{pmatrix} \exp(-\beta z) \exp i\left(\frac{n_1 k \sin \theta_{in}}{n_2} x - \omega t\right)$$

$$E_{eva} \begin{pmatrix} E_x \\ E_y \\ E_z \end{pmatrix} = \begin{pmatrix} A \\ 0 \\ C \end{pmatrix} \exp(-\beta z) \exp i\left(\frac{n_1 k \sin \theta_{in}}{n_2} x - \omega t\right)$$

exponentially decaying
away from the interface waves

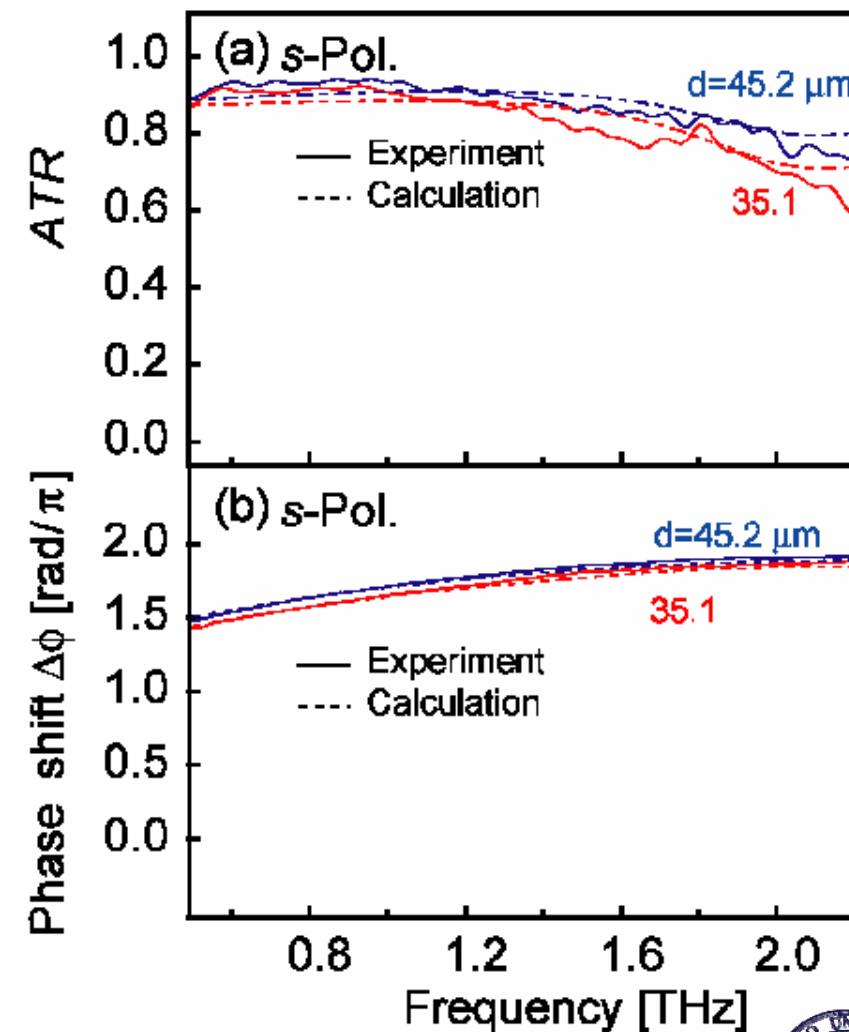
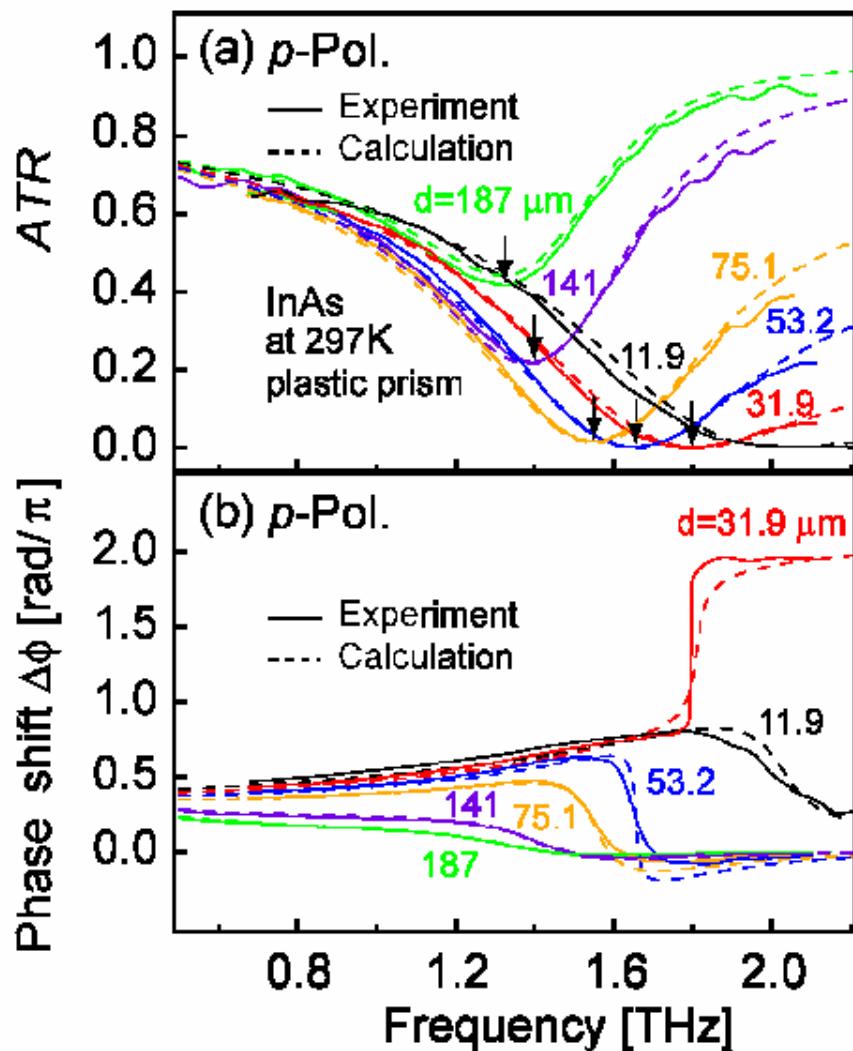


TM mode with large k_{\parallel}

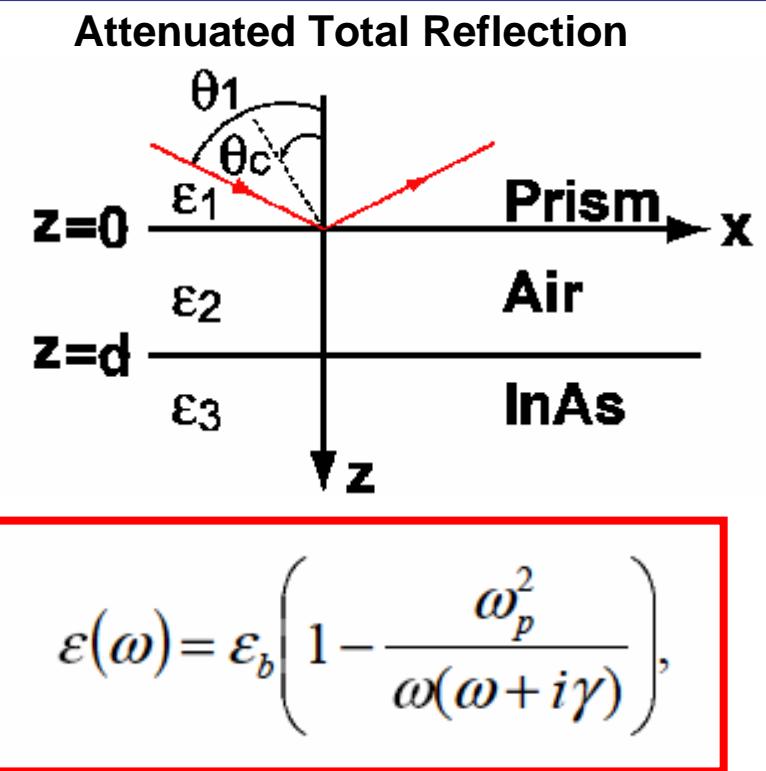
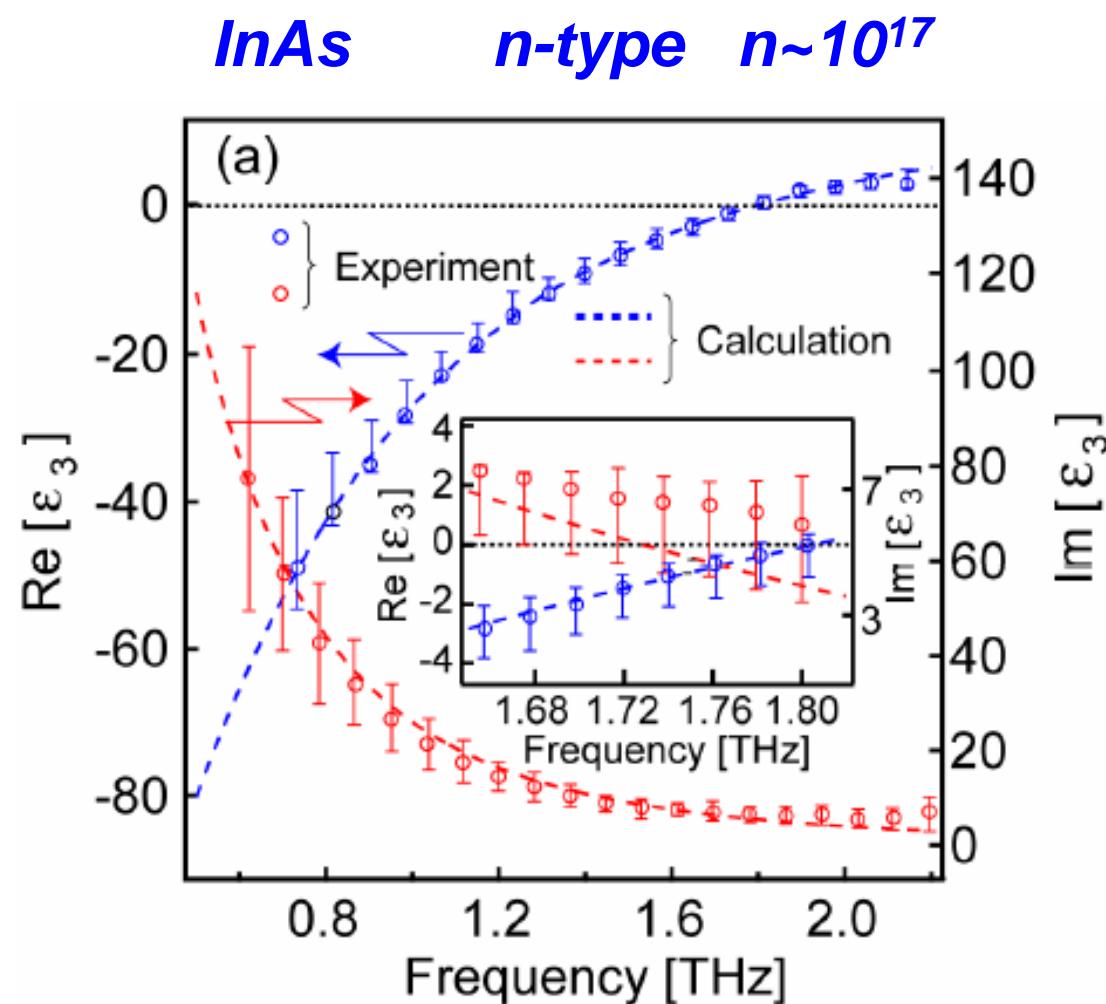
Enable to excite
surface modes.

Surface plasmon in semiconductor with THz-ATR

H. Hirori, M. Nagai, and K. Tanaka, Optics Express, 13, (26), 10801-10814 (2005).



Dielectric constant in doped semiconductor



$$\omega \ll \gamma$$

$$\epsilon(\omega) = i \frac{\epsilon_b \omega_p^2}{\gamma} \frac{1}{\omega} = i \frac{\sigma}{\omega}$$

H. Hirori, M. Nagai, and K. Tanaka, Optics Express, 13, (26), 10801-10814 (2005).

Kyoto University 2006

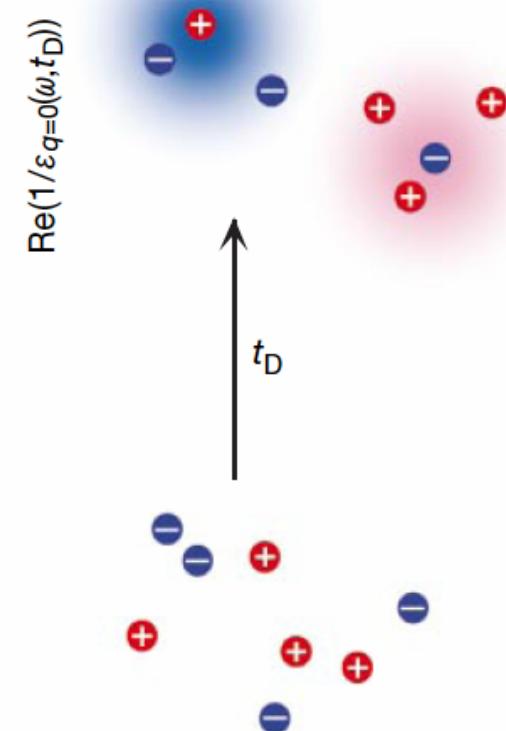
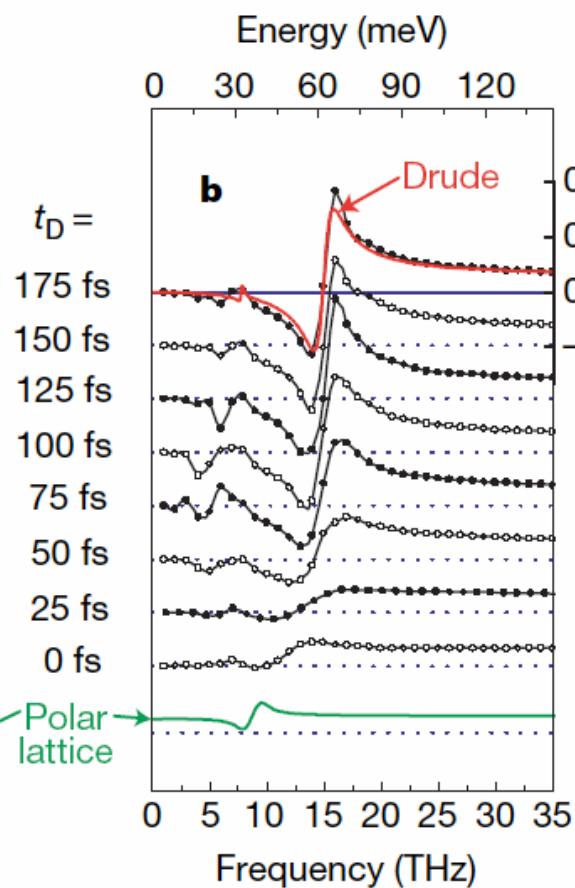
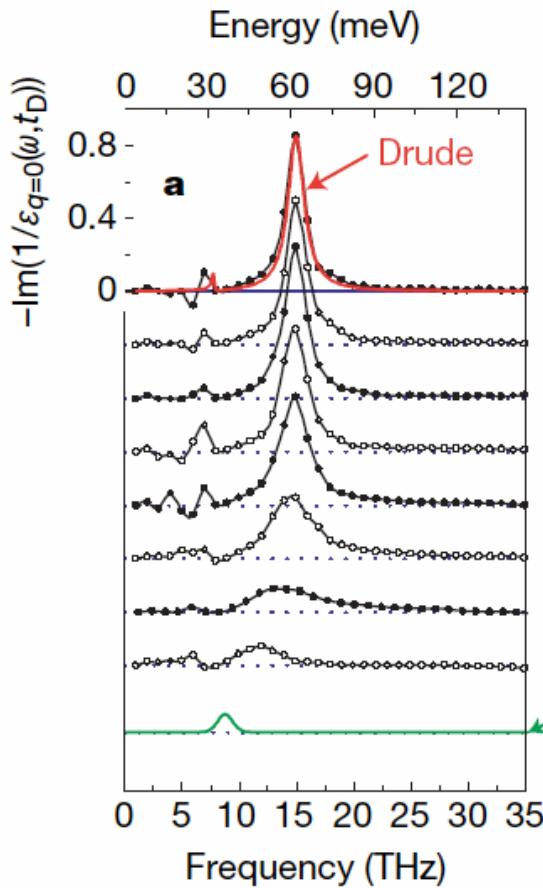


Pump and probe spectroscopy

“How many-particle interactions develop after ultrafast excitation of electron-hole plasma ”

R. Huber, F. Tauser, A. Brodschelm, M. Bichler, G. Abstreiter and A. Leitenstorfer

Nature. Vol.414 (2001) 286

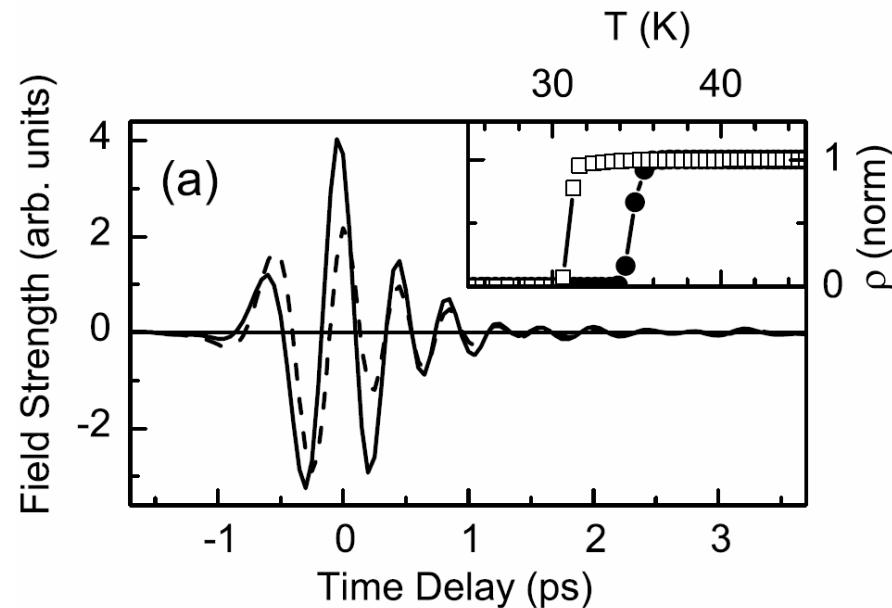


Superconducting gap

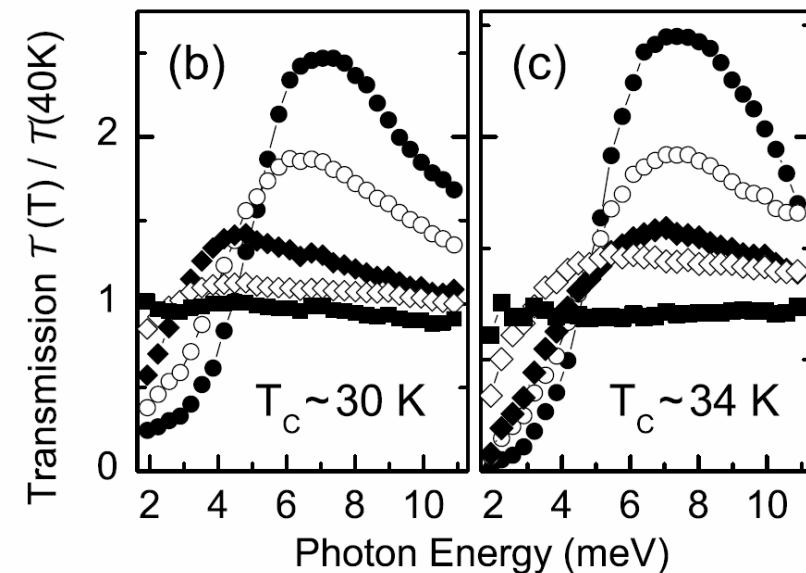
“Far-infrared optical conductivity gap in superconducting MgB₂ films”

R. A. Kaindl, M. A. Carnahan, J. Orestein, D. S. Chemla, H. M. Christen, H. Y. Zhai, M. Paranthaman and D. H. Lowndes

[Phys. Rev. Lett. 88 \(2002\) 027003](#)



100nm MgB₂薄膜における透過波形
6K(実線)と40K(破線)
Insetは100nm MgB₂薄膜(□)と200nm
MgB₂薄膜(●)の抵抗の温度変化

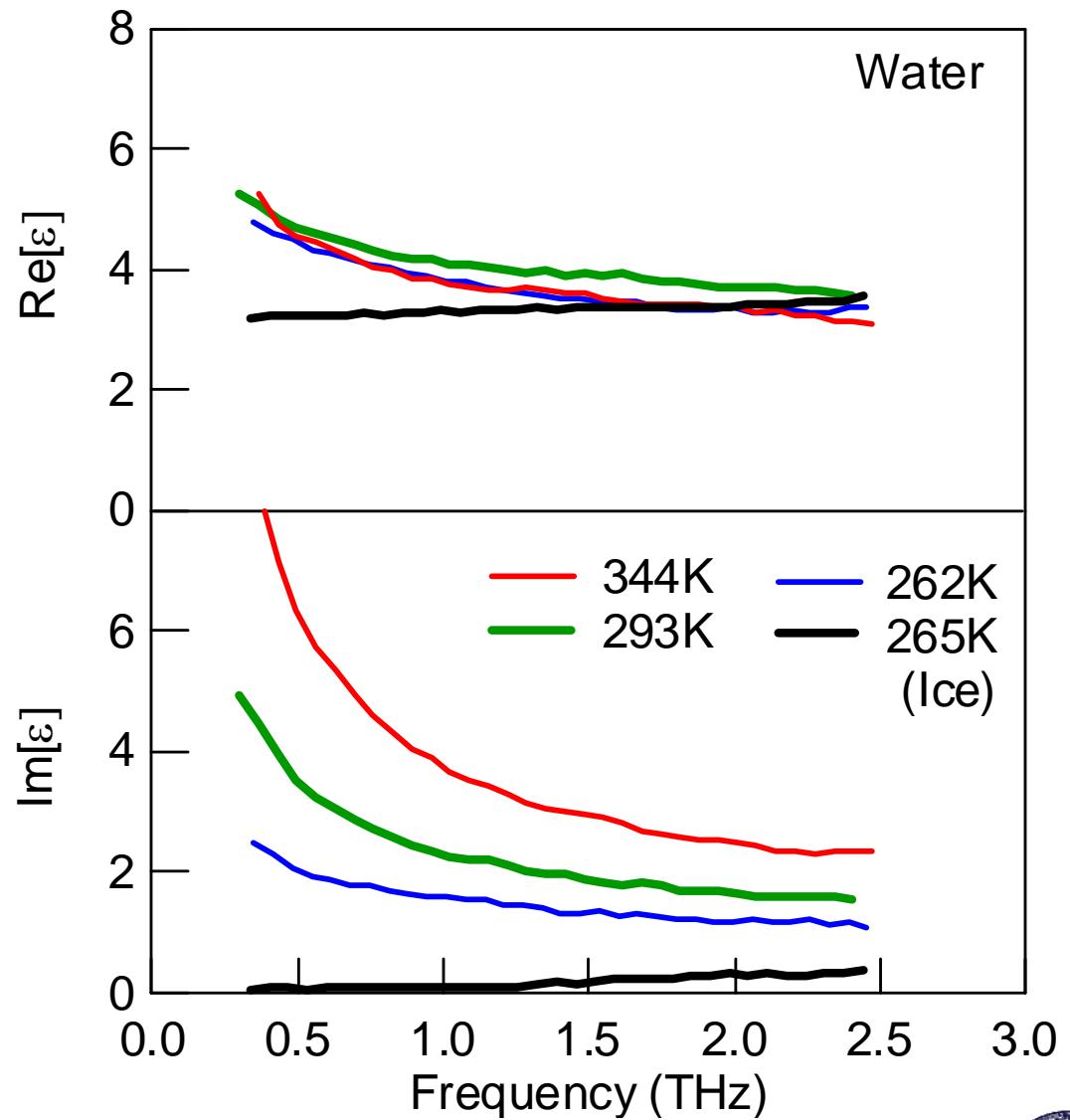
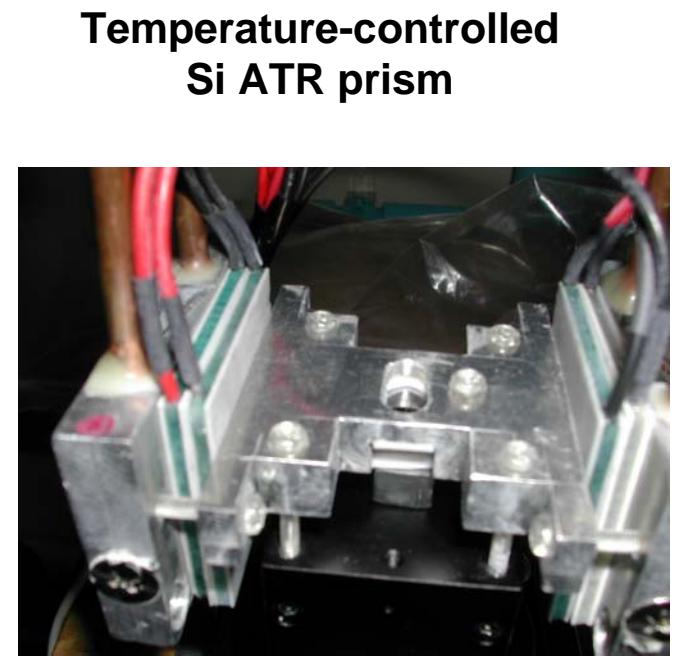


40Kの値で規格化した透過率

- (b)膜厚100nm MgB₂
6K(●), 20K(○) 27K(◆), 30K(◇), 33K(■)
(c)膜厚200nm MgB₂(6K, 20K 25K, 30K, 36K)
6K(●), 20K(○) 25K(◆), 30K(◇), 36K(■)



TD-ATR spectroscopy in water

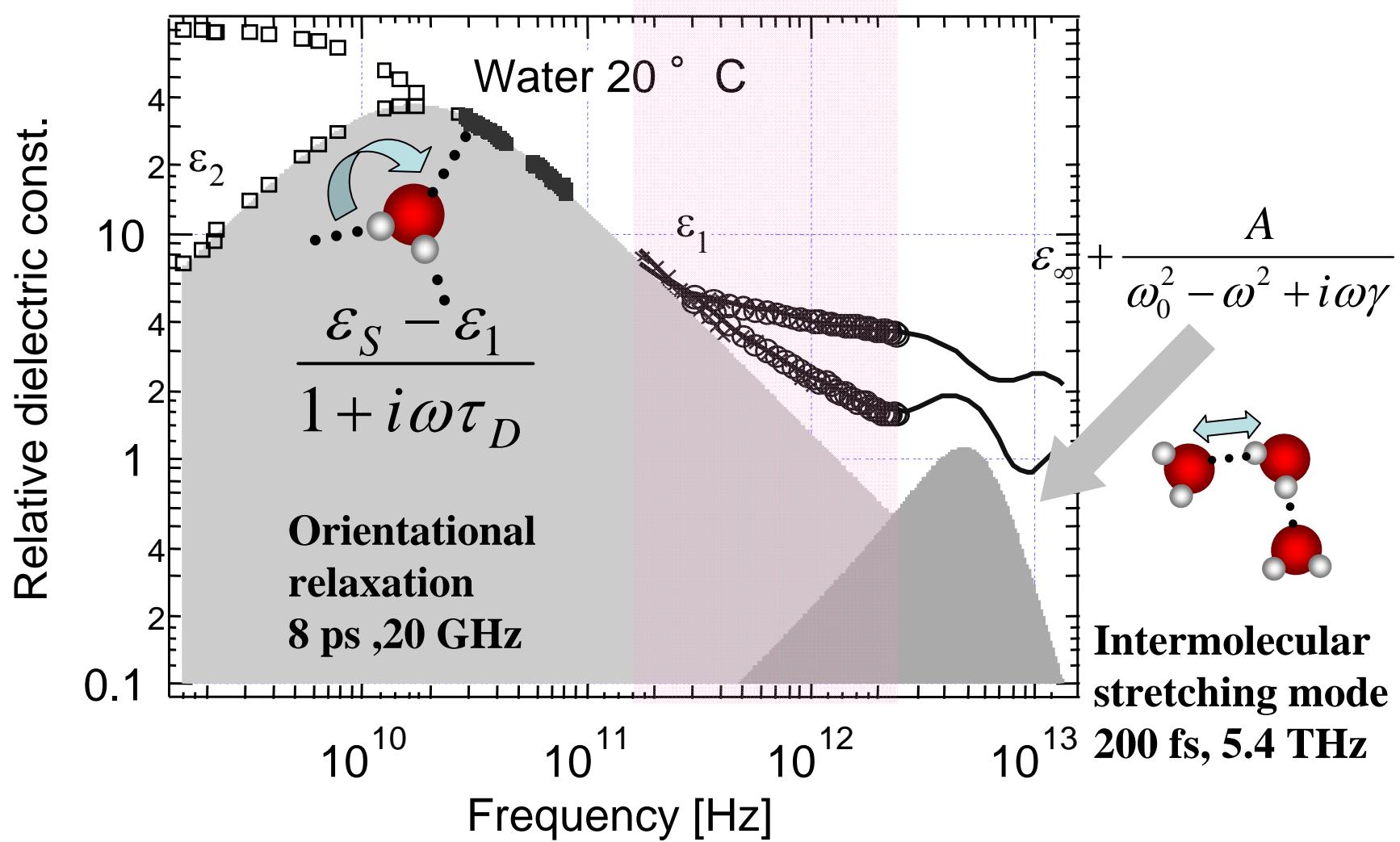


Nagai, et al. , Int. J. IRMMW, 27, 505 (2006).

Kyoto University 2006



Dielectric Constants in Water



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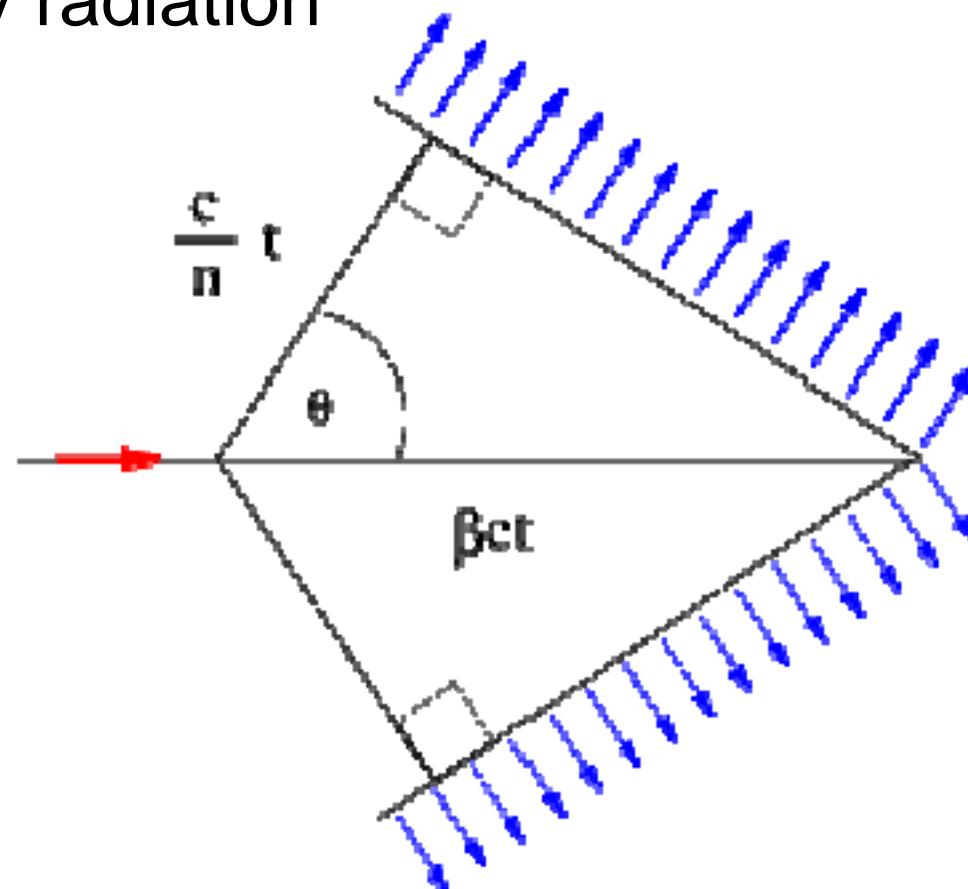
Non-linear THz spectroscopy

State of the art of the THz pulse power using fs laser

| | Pump power | | THz power | | |
|--|-------------------|----------------------|------------------|--------------------|-------------------------|
| | | /pulse | average | /pulse | Electric field (V/cm) |
| Ti: Al ₂ O ₃ laser | | | | | |
| Amplifier@1KHz | 650 mW | 650 µJ | 0.1µW | 125 pJ | 1 KV/cm |
| Oscillator with magnetic field (Sarukura 1.7T, @80MHz) | 1500 mW 650 mW | 18.75 nJ 8.125 nJ | 650 µW 110 µW | 8.125pJ 1.375pJ | 0.25 KV/cm 0.1 KV/cm |

Non-linear THz spectroscopy

- Cherenkov radiation



Cherenkov angle

$$\cos \theta = \frac{1}{n\beta}$$

Non-linear THz spectroscopy

- Basic idea of velocity matching by pulse front using Cherenkov effect with non-linear crystal

A. G. Stepanov, J. Hebling and J. Kuhl *et al.*, Appl. Phys. Lett. 83, 3000 (2003).

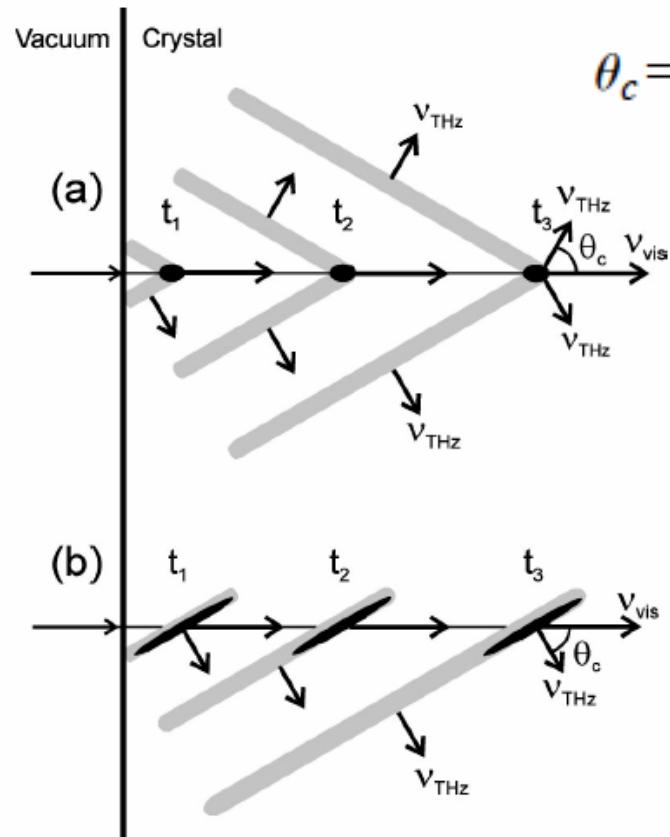


FIG. 1. Schematic illustration of the THz generation by tightly focused (a) and front tilted (b) femtosecond laser pulses propagating in an electro-optical crystal. Black ellipses and gray areas depict laser and generated THz pulses, respectively, at three different instants of time ($t_1 < t_2 < t_3$).

$$\theta_c = \arccos(v_{\text{THz}}/v_{\text{vis}}),$$

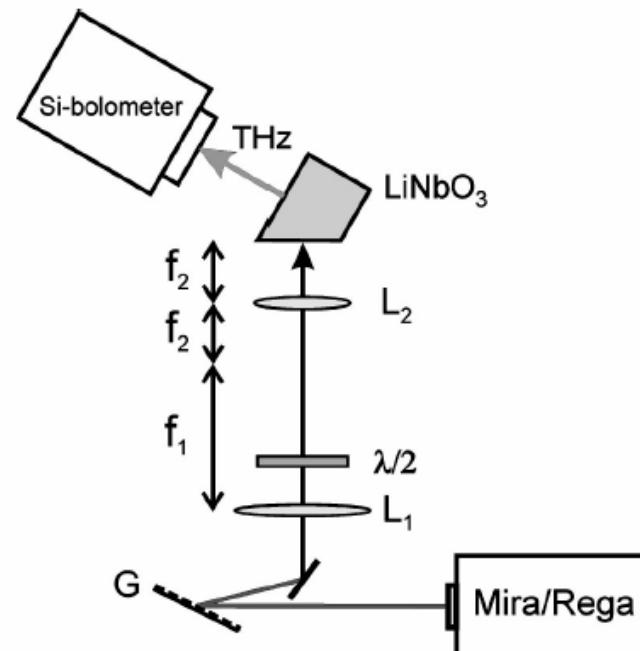
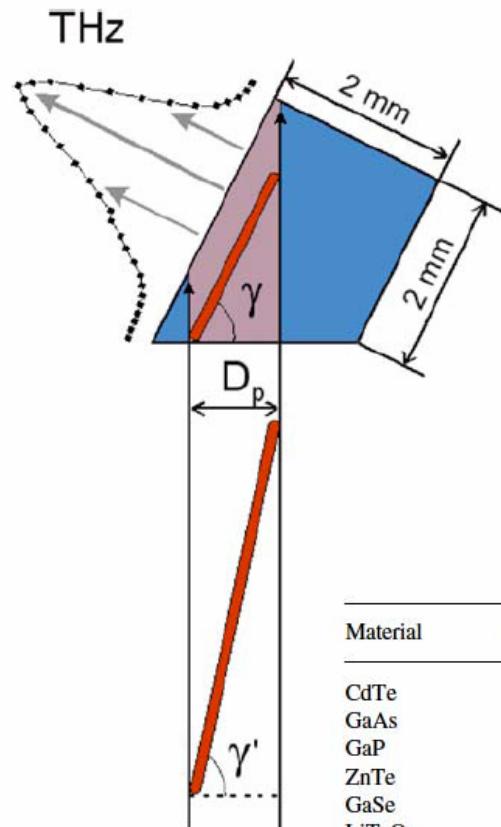


FIG. 2. Experimental setup used for the THz generation by femtosecond laser pulses with tilted pulse fronts.

Non-linear THz spectroscopy

- High power single cycle THz generation using tilted femtosecond light sources

J. Hebling and J. Kuhl et al., Appl. Phys. B 78, 593-599 (2004).



$$v_{vis}^{gr} \cdot \cos \gamma = v_{THz}^{ph}$$

$$v_{vis}^{gr} \gg v_{THz}^{ph}$$
$$n_v^{gr} \ll n_{THz}$$

| Material | r (pm/V) | d (pm/V) | n_v^* | n_v^{gr} | n_{THz} | α_{THz} (cm $^{-1}$) | FOM (pm 2 /V 2) | FOMA (pm 2 cm 2 /V 2) |
|--------------------|--------------------|------------|---------------------|------------|--------------------|------------------------------------|------------------------|---------------------------------|
| CdTe | 4.5 ²⁶ | 81.8 | 2.92 ²⁴ | 3.73 | 3.23 ²⁰ | 4.8 ²⁰ | 242 | 10.5 |
| GaAs | 1.43 ²⁶ | 65.6 | 3.68 ¹⁸ | 4.18 | 3.61 | 0.5 ¹⁸ | 87.9 | 352 |
| GaP | 0.97 ²² | 24.8 | 3.18 ¹⁸ | 3.57 | 3.34 ¹⁸ | 1.9 ¹⁸ | 18.2 | 5.0 |
| ZnTe | 4.04 ²⁶ | 68.5 | 2.87 ¹⁹ | 3.31 | 3.17 ²⁰ | 1.3 ²⁰ | 180 | 106 |
| GaSe | 1.7 ¹⁵ | 28.0 | 2.85 ¹⁷ | 3.13 | 3.72 ¹⁶ | 0.07 ¹⁸ | 25.9 | 5300 |
| LiTaO ₃ | 30.5 ²⁶ | 161 | 2.145 ²⁵ | 2.22 | 6.42 ²¹ | 46 ²¹ | 882 | 0.4 |
| LiNbO ₃ | 30.9 ²⁶ | 168 | 2.159 ²³ | 2.23 | 5.16 ²¹ | 10 ²¹ <small>LW</small> | | |

Below the table is a series of grayscale images showing THz pulses. The first four images are small, while the last three are larger and show more detail. A blue rectangle highlights the last three images.

Non-linear THz spectroscopy

- High power single cycle THz generation using tilted femtosecond light sources : Recent status

A. G. Stepanov and J. Heebling *et al.*, Optics Express 13 5762 (2005).

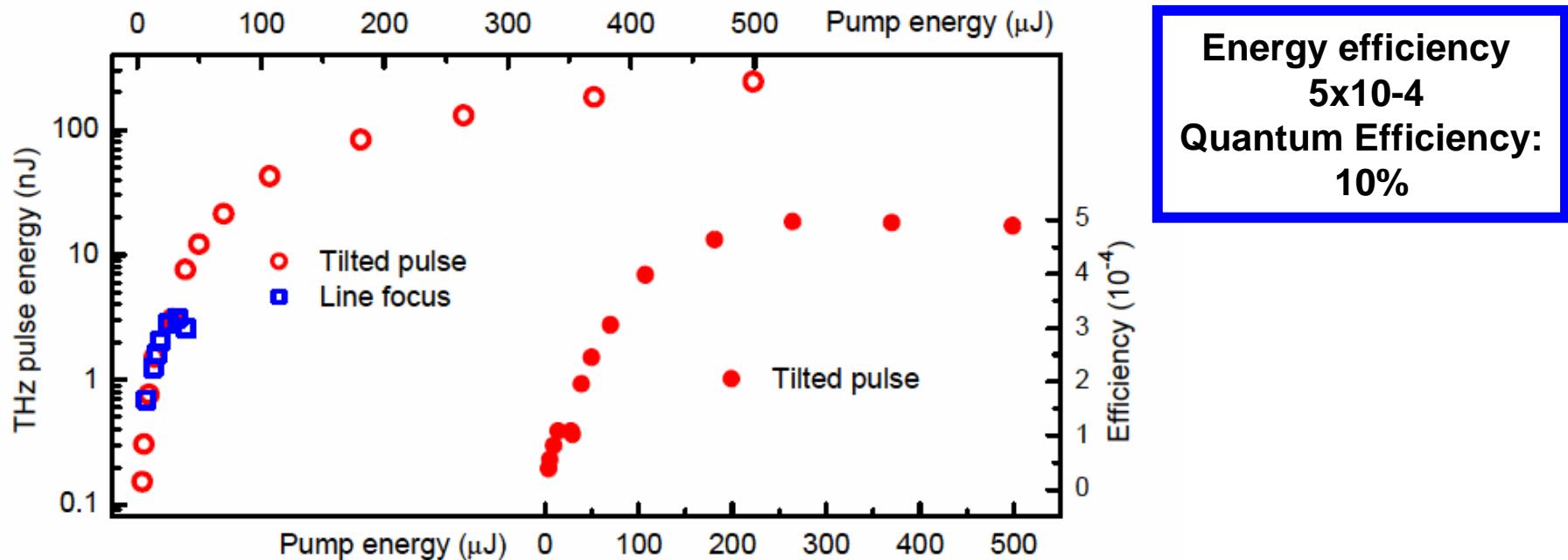


Fig. 1. Measured energy of THz pulses generated by the tilted pulse front (red circles) and line focusing (open blue squares) set-ups versus the energy of the 780 nm pump laser pulses (upper left part). Energy conversion efficiency versus the pump energy for the tilted pulse front set-up (lower right part).

$$P_{\text{THz}} = 100 \text{ nJ/pulse}, 10 \text{ nJ/cm}^2, 10 \text{ MW/cm}^2, |E_{\text{THz}}| \sim 30 \text{kV/cm} = 3 \text{ MV/m}$$