

## Polarization-dependent Ti *K* x-ray absorption and emission studies of Ti<sub>2</sub>O<sub>3</sub> single crystal

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### Introduction

Ti<sub>2</sub>O<sub>3</sub> exhibits a metal (high-temperature phase) - insulator (low-temperature phase) transition (MIT) around 450 K. The Ti *3d* levels in Ti<sub>2</sub>O<sub>3</sub> split into the *t<sub>2g</sub>* and *e<sub>g</sub>* levels. Due to the trigonally distorted crystal field, the *t<sub>2g</sub>* level further splits into the *e<sub>g</sub><sup>π</sup>* and *a<sub>1g</sub>* levels. Tanaka has proposed a new MIT model taking into account the many-body effect and predicted that the Ti *3d* configuration of the Ti ion pair along *c*-axis changes from (*a<sub>1g</sub>*↑, *a<sub>1g</sub>*↓) in the insulating phase to (*a<sub>1g</sub>*↑, *e<sub>g</sub><sup>π</sup>*↑) in the metallic phase [1]. Recently, we have detected the *3d* configuration change due to MIT by means of polarization-dependent Ti *2p-3d* absorption spectroscopy [2]. In this report, we present the polarization-dependent Ti *K* x-ray absorption and emission (XAS and XES) spectra of Ti<sub>2</sub>O<sub>3</sub> single crystal.

### Experiment

The Ti *K* XAS and XES experiments were carried out at BL-7C and BL-15B1 with polarized and depolarized configurations for XES, respectively. ESCARGOT with a Ge(400) crystal and a PSPC detector were used to measure the XES spectra. Both spectra were measured with *E*//*c* and *E*⊥*c* conditions, where *E* denotes a polarization vector of incidence light. Note that the Ti ion pair is parallel to the *c*-axis. Experimental data presented here is taken at room temperature.

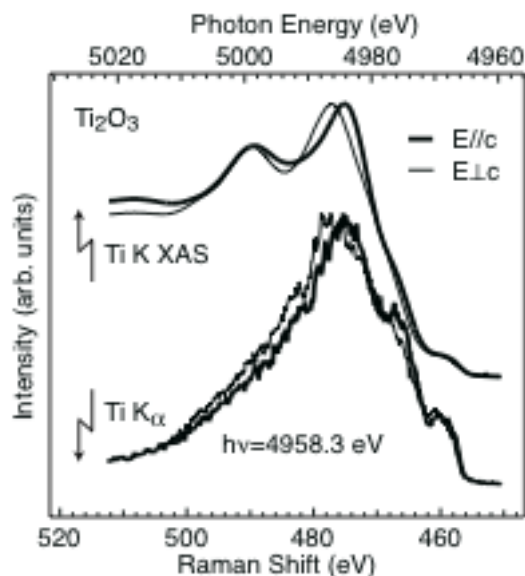
### Results and discussion

An upper side in Fig. 1 shows Ti *K* XAS spectra of Ti<sub>2</sub>O<sub>3</sub>. Thick and thin curves represent spectra measured with *E*//*c* and *E*⊥*c* conditions. One notices that a feature of the XAS spectra is different between two conditions. Peak energy of a white line of the *E*//*c*-spectra is lower than *E*⊥*c*-spectra by about 2 eV and a clear shoulder is found just below the white line. This difference shows a remarkable anisotropy in the unoccupied Ti *4p* states of Ti<sub>2</sub>O<sub>3</sub>. The lower energy of the *4p*<sub>∥</sub> bands than the *4p*<sub>⊥</sub> bands contradicts our intuition, considering the *a<sub>1g</sub>* states is mainly occupied.

A lower side in Fig. 1 shows the *E*//*c*- and *E*⊥*c*- XES spectra in the *K<sub>α</sub>* region measured at 4958.3 eV, far below

the Ti *K*-edge, with polarized configuration. Abscissa is denoted by the Raman shift. Interestingly, the polarization-dependent Raman spectra are very similar to the *K* XAS spectra with respect to the energy shift and the shoulder structure. This experimental result suggests that the structure of the Raman spectra originates from the *2p-4p* transition. This trend is also found for TiO<sub>2</sub> single crystal [3].

Features of the XES spectra strongly depend on the excitation energy below 4964 eV and show several fine structures depending on the *E*//*c* and *E*⊥*c* conditions. Detailed analysis is in progress and experiments above 450 K is also planned.



**Fig.1.** Polarization-dependent Ti *K* XAS and Raman spectra of Ti<sub>2</sub>O<sub>3</sub>. The Raman spectra were measured at 4958.3 eV, far below the *K*-edge.

### References

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